

Hydrogen Atom

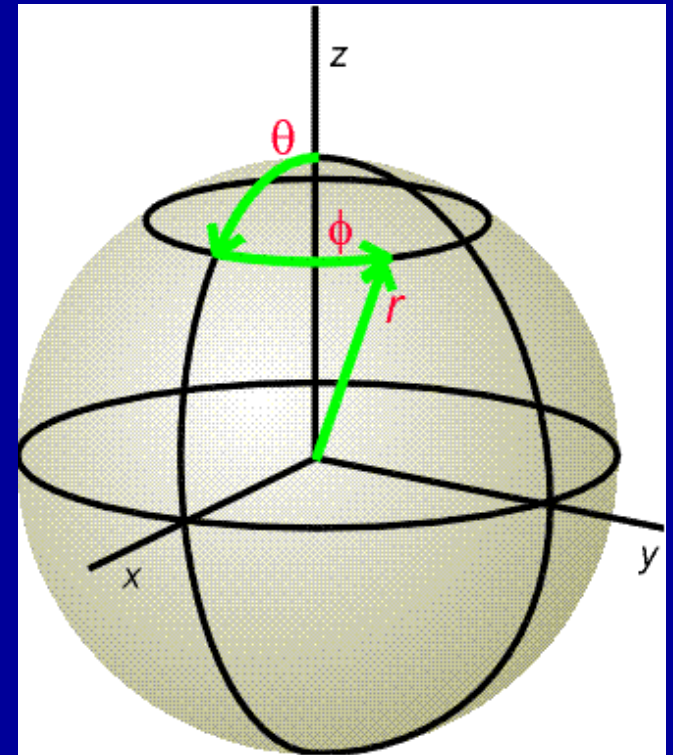
The simplest system: p + e

Schrodinger equation can be solved exactly

Spherical symmetry

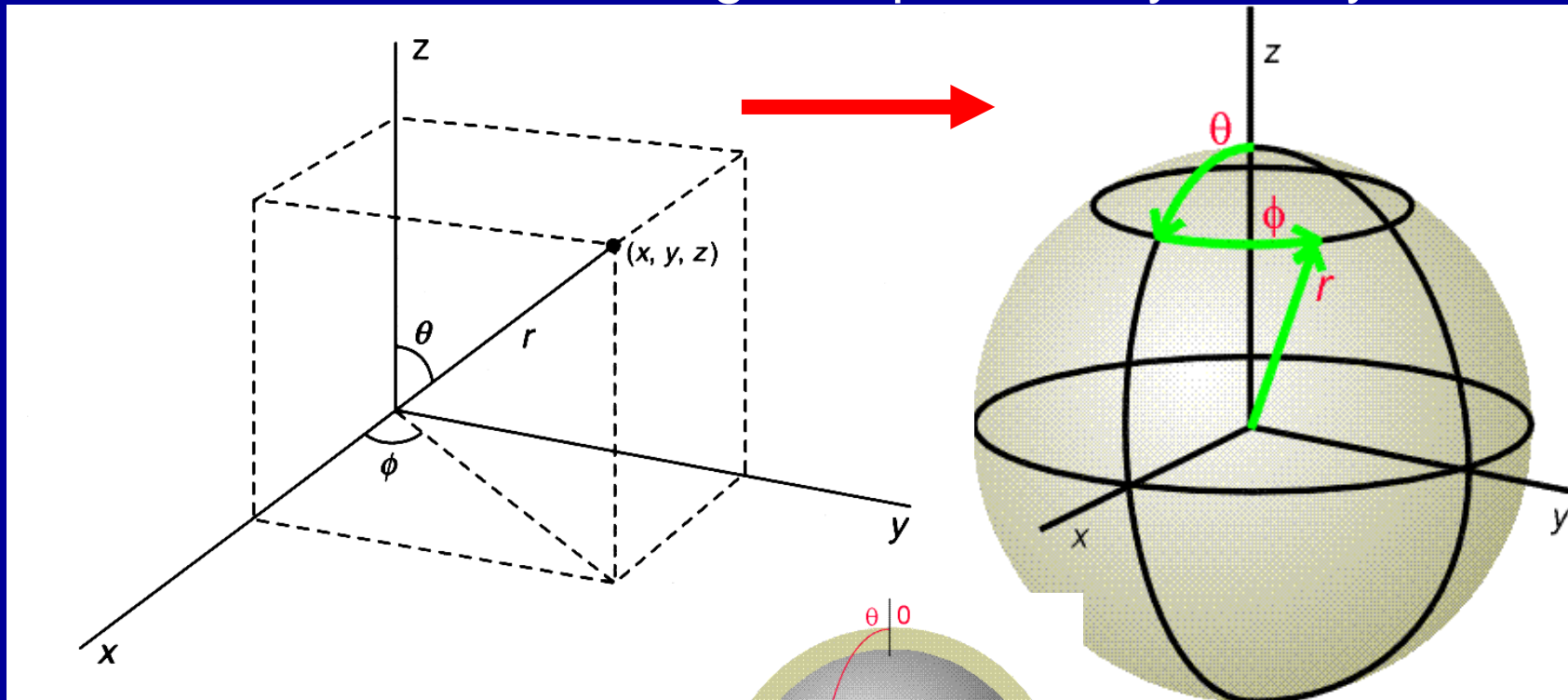
Potential energy between p + e

$$V = -\frac{e^2}{4\pi\epsilon_0 r}$$



Polar Coordinates

Take advantage of spherical symmetry



$$\Psi(x, y, z) \rightarrow \Psi(r, \theta, \phi)$$

$$x = ?$$

$$y = ?$$

$$z = r \cos \theta$$

Radial and Angular Part

$$\Psi_{n, l, m}(r, \theta, \phi) = N \times R_{n, l}(r) \times \chi_{l, m}(\theta, \phi)$$

Separation of variables

$R_{n, l}(r)$ = radial part of the wave function, depends only on distance from a nucleus - r

$\chi_{l, m}(\theta, \phi)$ = angular (angles) part of the wave function, depends only on direction - θ, ϕ

N = normalisation constant

In order to be $\int |\Psi|^2 dV = +1$

Normalisation condition - electron is definitely somewhere with probability = 1

Quantum Numbers

$R_{n, l}(r)$ depends on quantum numbers n and l

$\chi_{l, m}(\theta, \phi)$ depends on quantum numbers l and m_l

Principal quantum number n , (1 to ∞)

Angular momentum quantum number l , (0 to $n - 1$)

$l = 0$ (s), 1 (p), 2 (d), 3 (f), 4 (g), 5 (h),

Magnetic quantum number m_l , (+ l ,0, $-l$)

For each l there is $(2l + 1)$ values of m_l

Magnetic spin quantum number m_s ($\pm 1/2$)

Wave Functions of H atom

$$\varphi_{1,0,0} = \frac{1}{\sqrt{\pi a_0^3}} e^{-r/a_0}$$

$$\varphi_{2,0,0} = \frac{1}{\sqrt{8\pi a_0^3}} \left(1 - \frac{r}{2a_0}\right) e^{-r/2a_0}$$

$$\varphi_{2,1,0} = \frac{1}{4\sqrt{2\pi a_0^3}} \frac{r}{a_0} e^{-r/2a_0} \cos \theta$$

$$\varphi_{2,1,1} = -\frac{1}{8\sqrt{\pi a_0^3}} \frac{r}{a_0} e^{-r/2a_0} \sin \theta e^{i\phi}$$

$$\varphi_{3,0,0} = \frac{1}{3\sqrt{3\pi a_0^3}} \left[1 - \frac{2}{3} \frac{r}{a_0} + \frac{2}{27} \left(\frac{r}{a_0}\right)^2\right] e^{-r/3a_0}$$

$$\varphi_{3,1,m} = \frac{4}{27\sqrt{3\pi a_0^3}} \frac{r}{a_0} \left(1 - \frac{1}{6} \frac{r}{a_0}\right) e^{-r/3a_0} Y_{1,m}(\theta, \phi)$$

$$\varphi_{3,2,m} = \frac{2}{81\sqrt{15\pi a_0^3}} \left(\frac{r}{a_0}\right)^2 e^{-r/3a_0} Y_{2,m}(\theta, \phi)$$

- solution of Schrödinger equation
- complex function of coordinates x, y, z or better r, ϕ , θ
- no physical meaning
- positive and negative values
- $|\Psi|^2$ probability density of finding electron e

Radial Part of the Wave Function of H Atom

n	l	m_l	$R_{n,l}(r)$
1 (K)	0 (s)	0	$2 (Z/a_0)^{3/2} \exp(-Zr/a_0)$
2 (L)	1 (p)	0	$2 (Z/2a_0)^{3/2} (1 - Zr/2a_0) \exp(-Zr/2a_0)$
2 (L)	1 (p)	± 1	$2/\sqrt{3} (Z/2a_0)^{3/2} (Zr/2a_0) \exp(-Zr/2a_0)$

Electron Energy in H-type Atoms

$$E_n = -\frac{N_A \mu e^4 Z^2}{8 \epsilon_0^2 h^2 n^2}$$

μ = reduced mass of nucleus-electron

e = elementary charge, ϵ_0 = permittivity of vacuum

Z – the higher a nucleus charge the stronger is an electron bound, the lower energy has, one-electron ions (He⁺, Li²⁺,.....)

n – the higher a principal number the less stable e is

Corresponds to Bohr's eq.!!

$$E_n = -E_0 \frac{Z^2}{n^2}$$

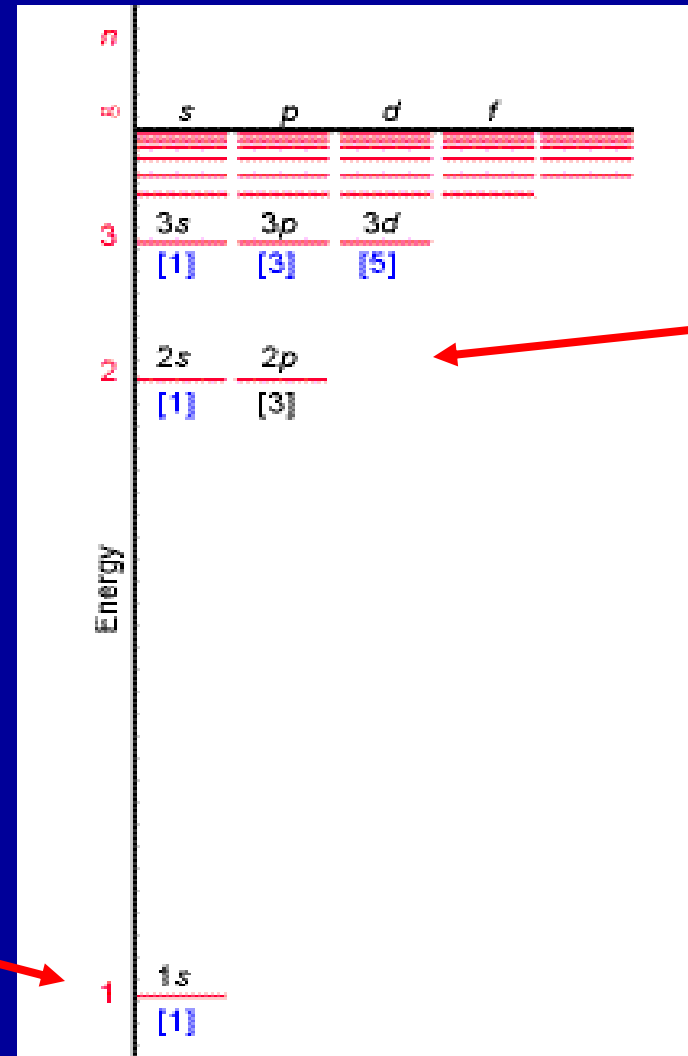
Electron Energy in H-type Atoms

$$E_n = -\frac{N_A \mu e^4 Z^2}{8 \epsilon_0^2 h^2 n^2}$$

Energy depend ONLY on n

$$E_1 = -13.6 \text{ eV}$$

$$(13.6 \text{ eV} = 1 \text{ Ry})$$



$E_2 = ?$

Principal Quantum Number n

Gives the levels energy
Higher n has higher energy -
less stable
 n same as in the Bohr's model

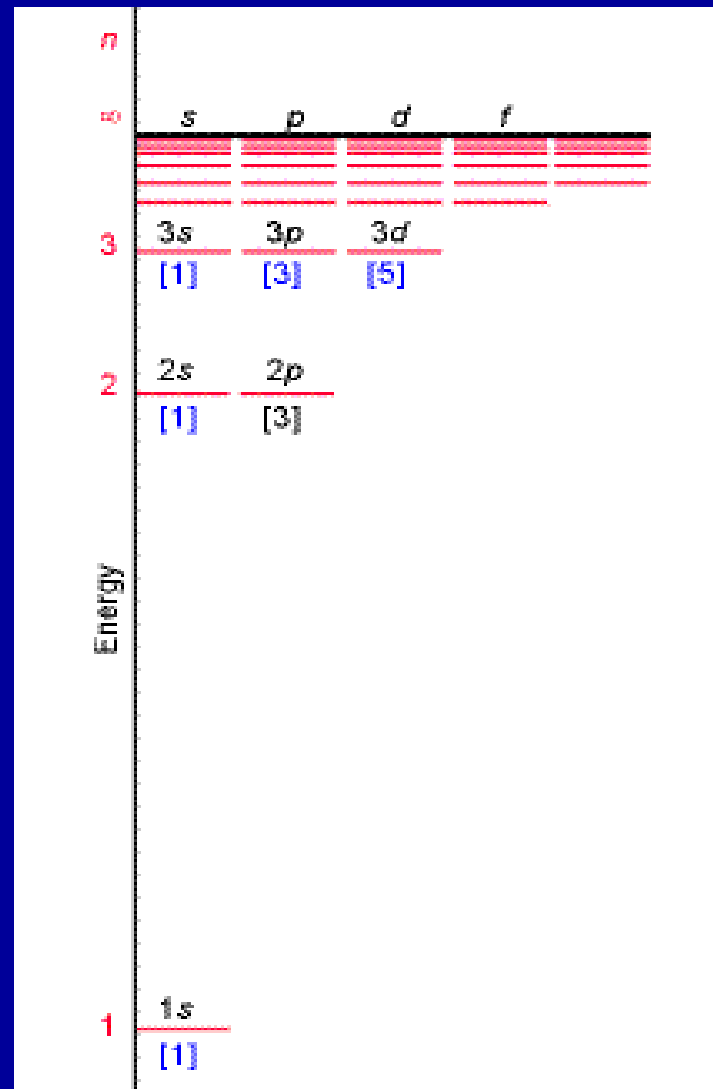
Attains values 1 to ∞

For each n there is n^2 of
degenerate levels

$$l = n - 1$$

$$\sum_{l=0}^{n-1} (2l + 1) = n^2$$

$$l = 0$$



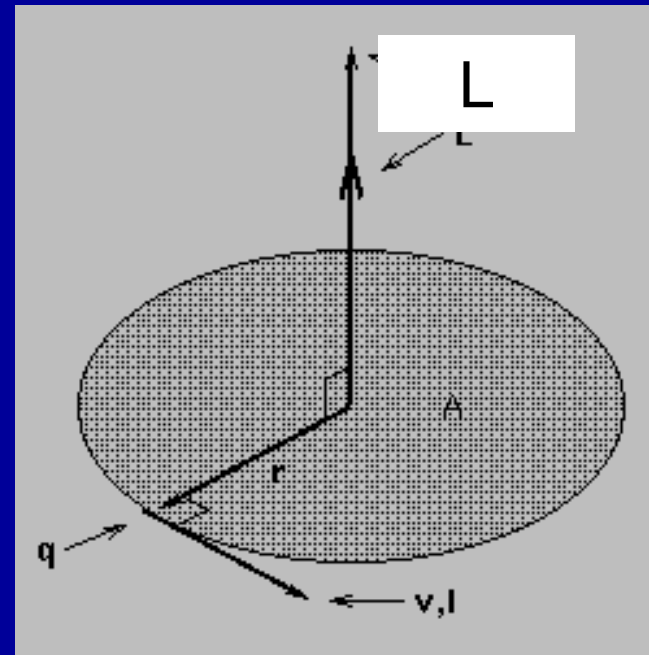
Orbital Angular Momentum

L = Orbital angular momentum (vector)

$$L = m \times v \times r = p \times r$$

$$L = \hbar \sqrt{l(l+1)}$$

Describes movements of electrons in orbitals



Angular Momentum Quantum Number l

Type of orbital, (0 to $n - 1$)

l	orbital
0	s
1	p
2	d
3	f
4	g
5	h
6	i
7	j
8	k

L = Orbital angular momentum
 $L = m \times v \times r$

$$L = \hbar \sqrt{l(l+1)}$$

these orbitals are not filled by electrons
in atoms in ground state

Magnetic Quantum Number m_l

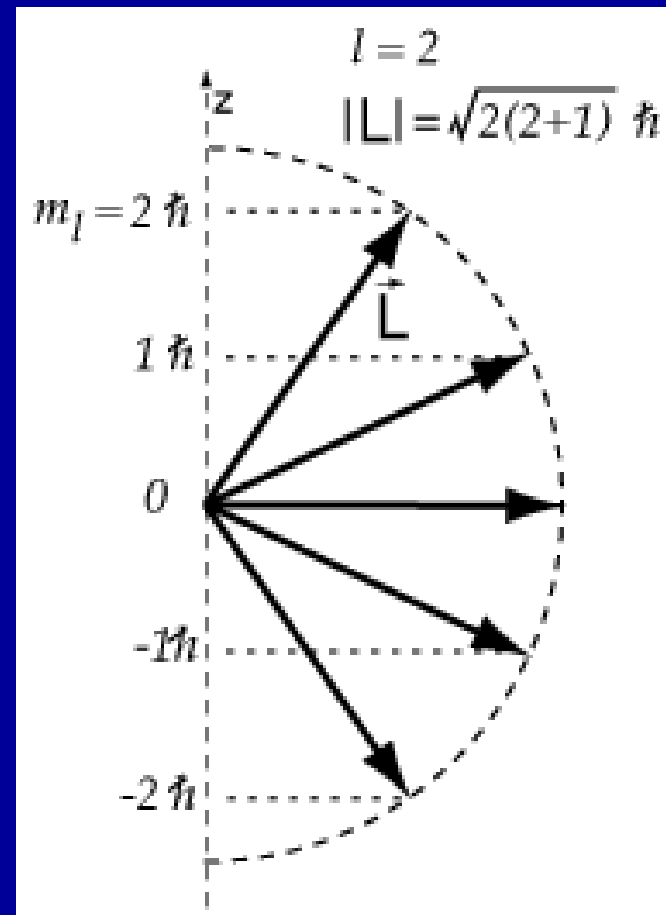
l	orbital	m_l
0	s	0
1	p	1, 0, -1
2	d	2, 1, 0, -1, -2
3	f	3, 2, 1, 0, -1, -2, -3
4	g	these orbitals are not filled by electrons in atoms in ground state
5	h	
6	i	

$$L_z = m_l \hbar = m_l \frac{h}{2\pi}$$

Quantization of Orbital Angular Momentum

$$L = \hbar \sqrt{l(l+1)}$$

$$L_z = m_l \hbar = m_l \frac{h}{2\pi}$$



For each n there is n^2 of degenerate levels

	s	p	d	f	g	h
$l =$	0	1	2	3	4	5
$n = 1$	1s					
$n = 2$	2s	2p				
$n = 3$	3s	3p	3d			
$n = 4$	4s	4p	4d	4f		
$n = 5$	5s	5p	5d	5f	5g	
$n = 6$	6s	6p	6d	6f	6g	6h

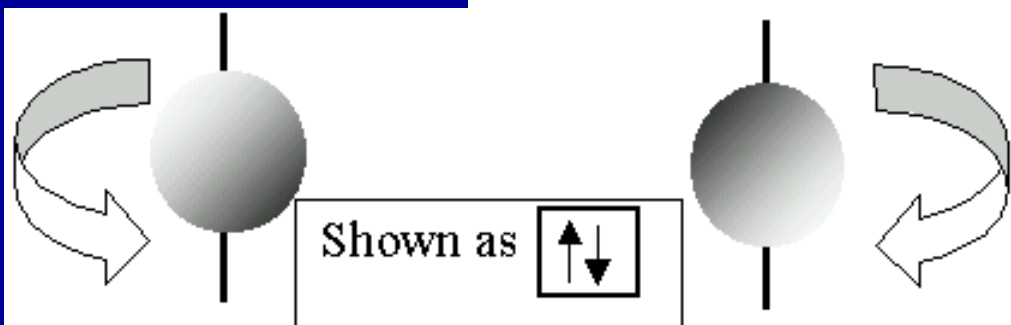
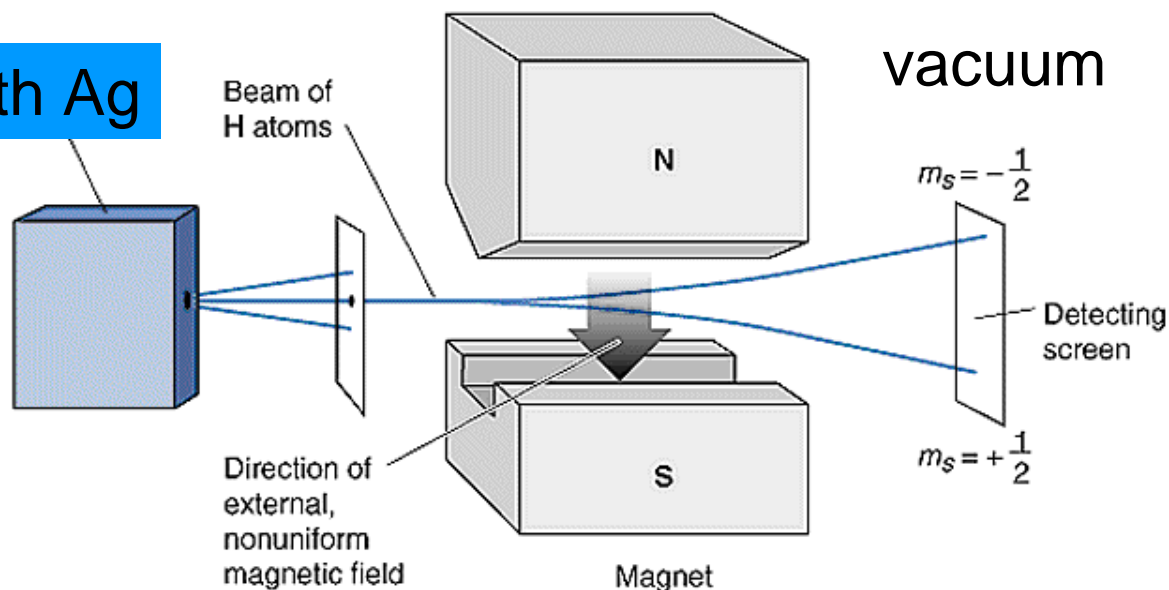
Magnetic Spin Quantum Number m_s

Stern-Gerlach experiment

Inhomogeneous magnetic field

Furnace with Ag

S = spin momentum

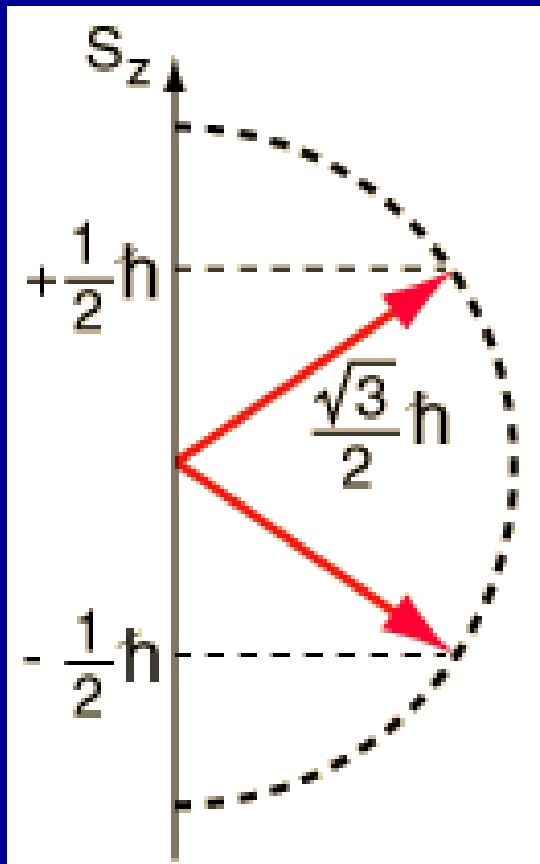


$$S = h/2\pi [s (s + 1)]^{1/2}$$

$$s = 1/2$$

$$S_z = m_s h/2\pi$$

Magnetic Spin Quantum Number m_s

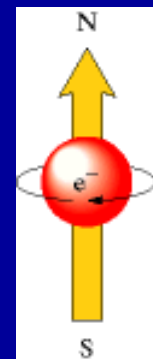
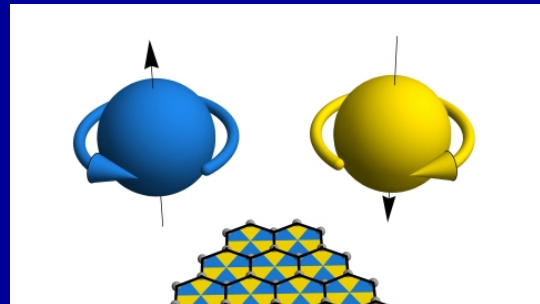


$$S = h/2\pi [s (s + 1)]^{1/2}$$

$$s = 1/2$$

$$S_z = m_s h/2\pi$$

$$m_s = \pm 1/2$$

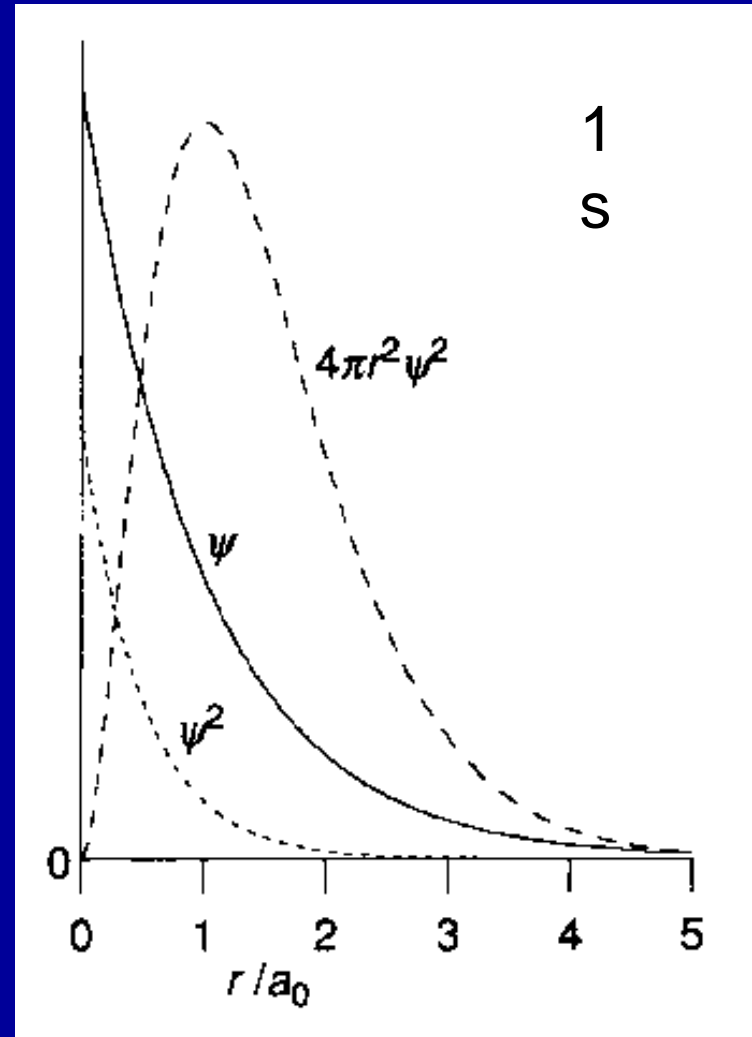


Ψ = Wave Functions

Ψ = solution of Schrödinger equation

$|\Psi|^2 =$ probability density of e

$|\Psi|^2 dV =$ probability density of finding electron e in volume $dV =$ distribution of electron density



Probability Density

Polar coordinates

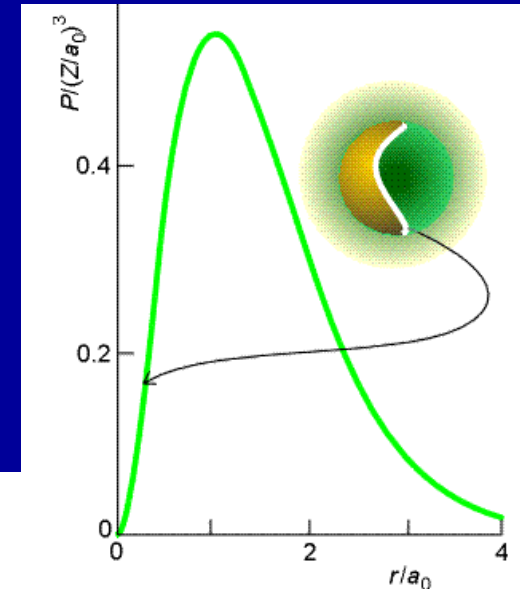
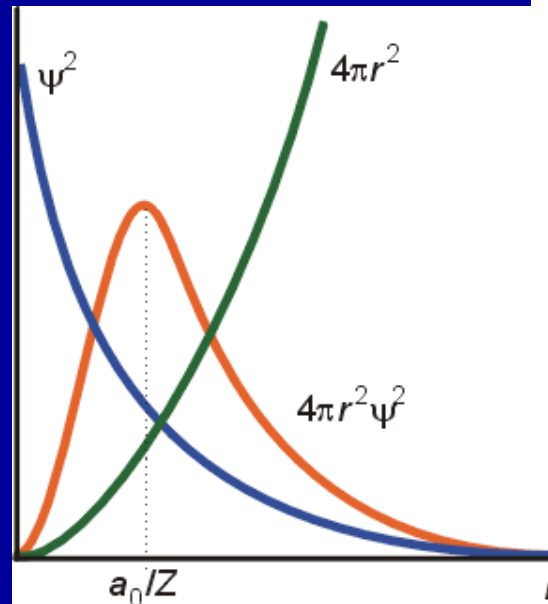
$R_{n,l}(r)$ radial function

$dV = 4\pi r^2 dr$ (spherical layer of thickness dr)

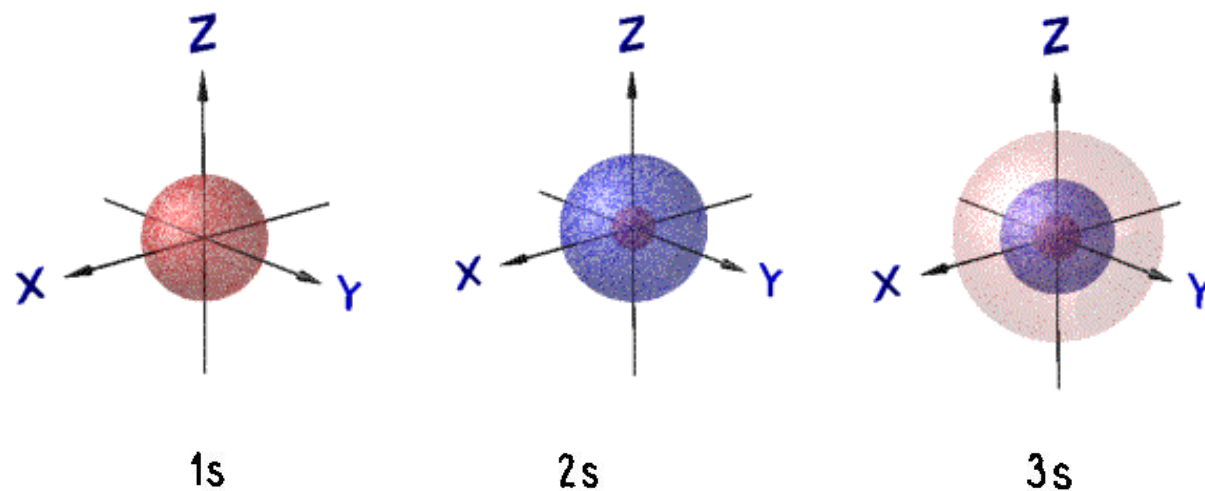
Radial distribution function

$$P = 4\pi r^2 |\Psi|^2 dr = 4\pi r^2 R_{n,l}^2(r) dr$$

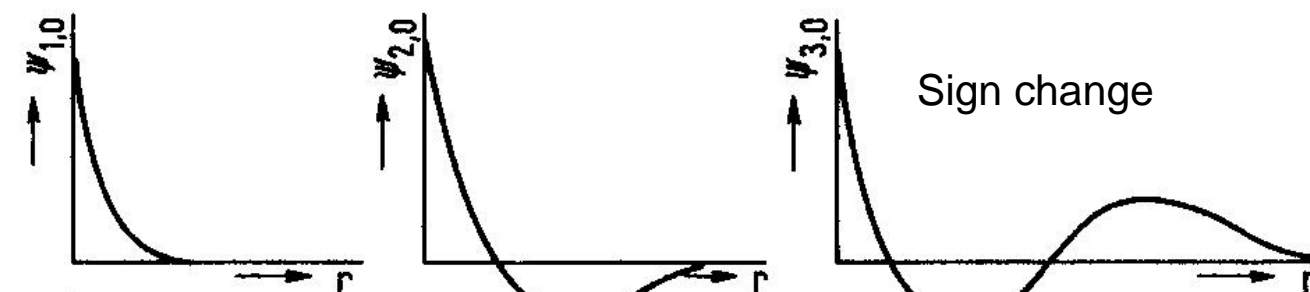
P = probability density
of finding electron e
in volume
of spherical layer
of thickness dr
In a distance r



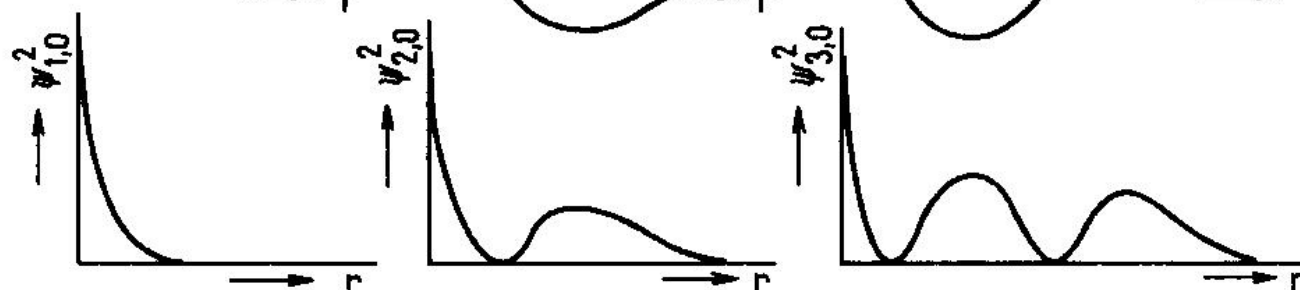
Orbital



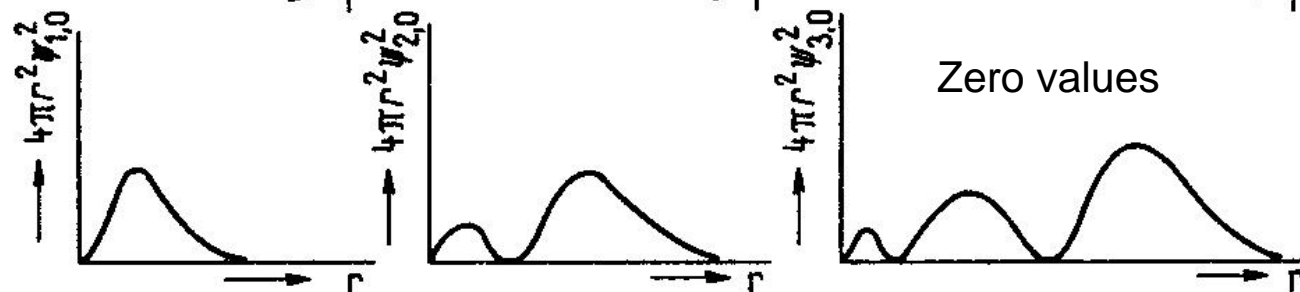
Wave Functions



Probability density



Radial distribution function

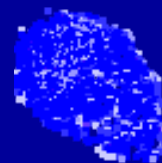
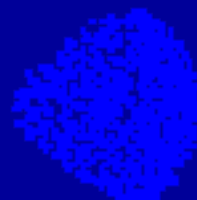
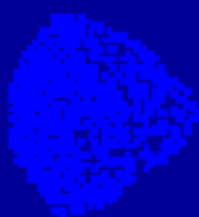
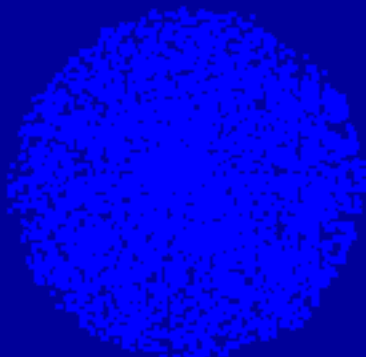


Orbitals

Position of electrons cannot be established – Heisenberg's principle – only probability

Radial function – probability of finding e in a direction away from nucleus (to $r = \infty$) and number of nodes = zero values of radial distribution function

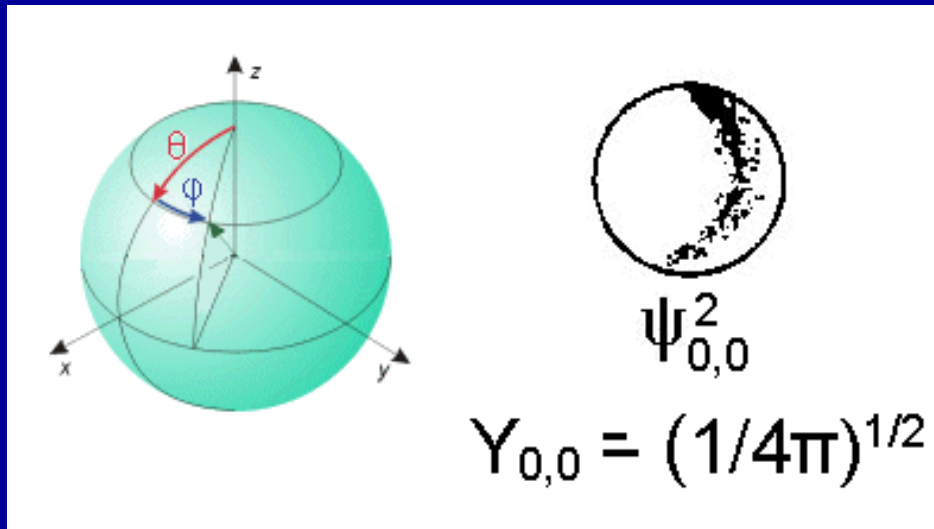
Angular function = shape of orbitals (number of nodal planes)



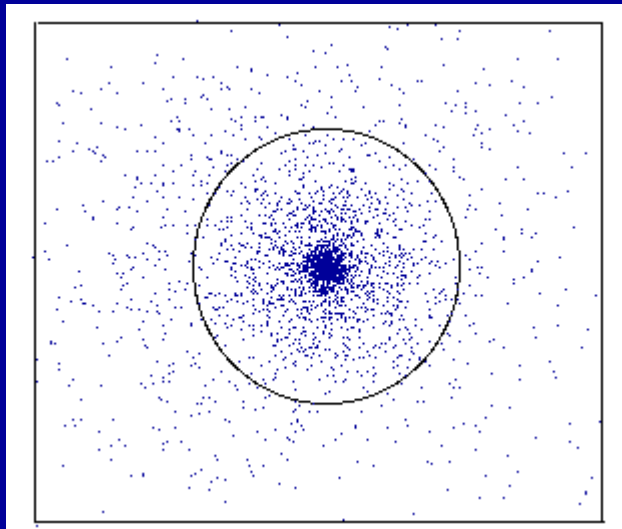
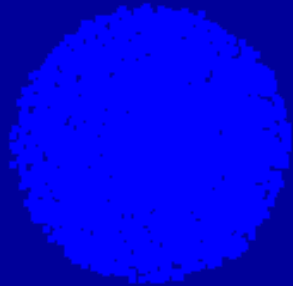
s - Orbitals

$R_{n,l}(r)$ = Radial function, depends on r only

$\chi_{l,m}(\theta, \phi)$ = angular function, is a **constant** for s-orbitals ($l = 0$) = SPHERICAL SHAPE

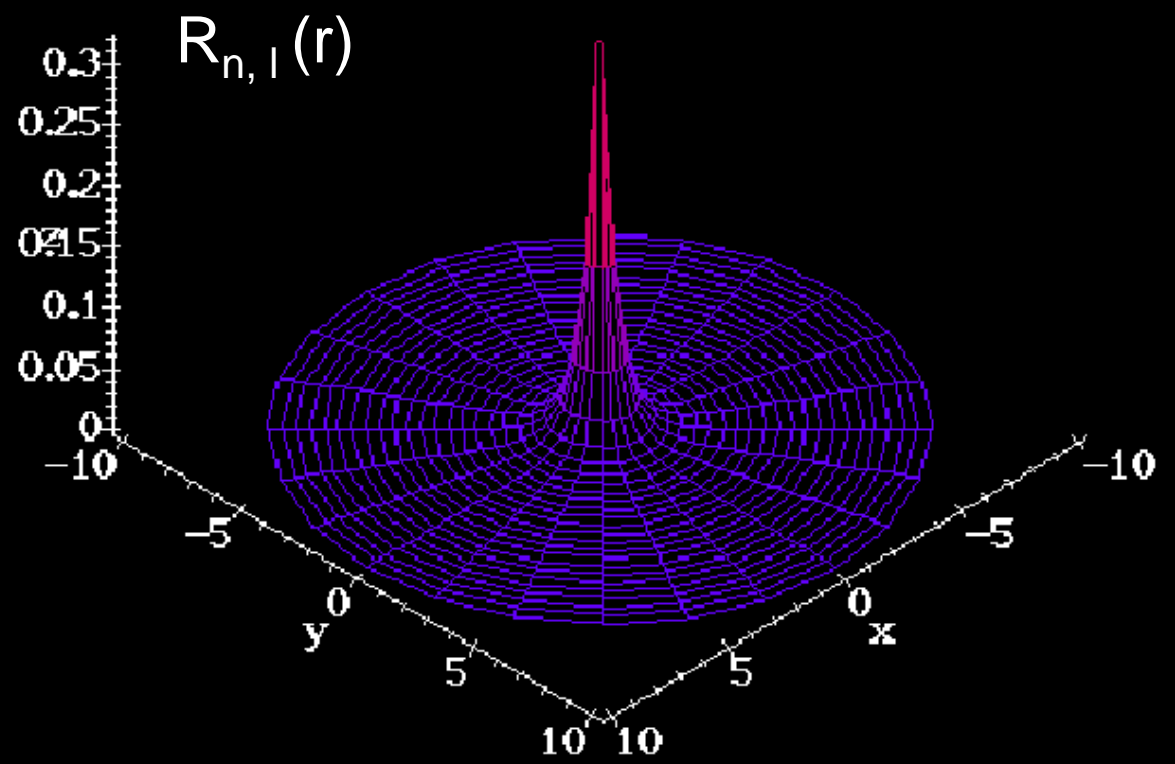


Atomic Orbital 1s



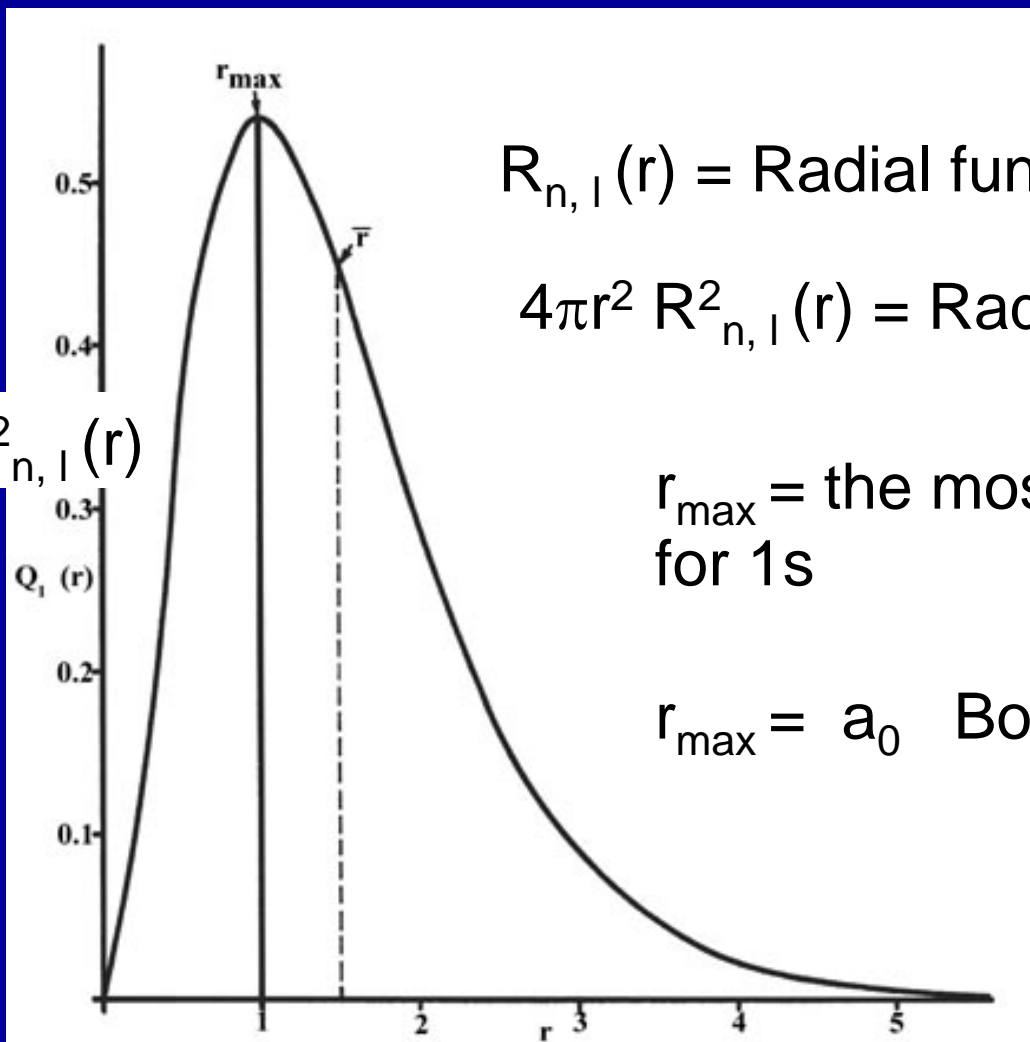
Wave Function 1s

$n = 1, l = 0$



Radial Distribution Function

$$4\pi r^2 R_{n,l}^2(r)$$

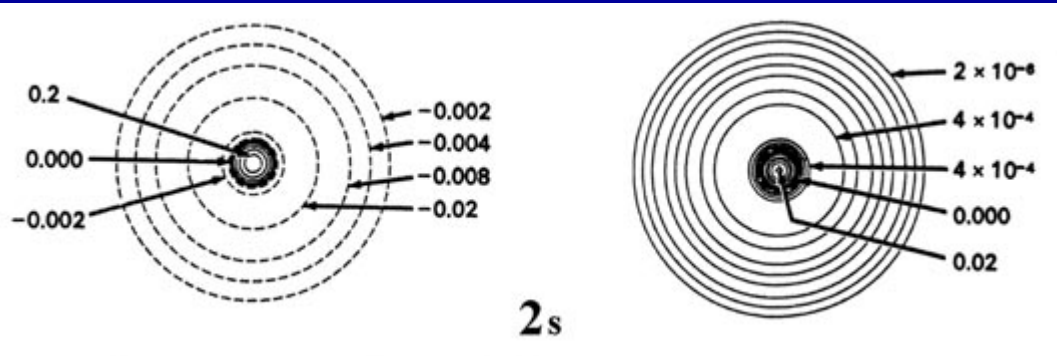


$R_{n,l}(r)$ = Radial function of H atom

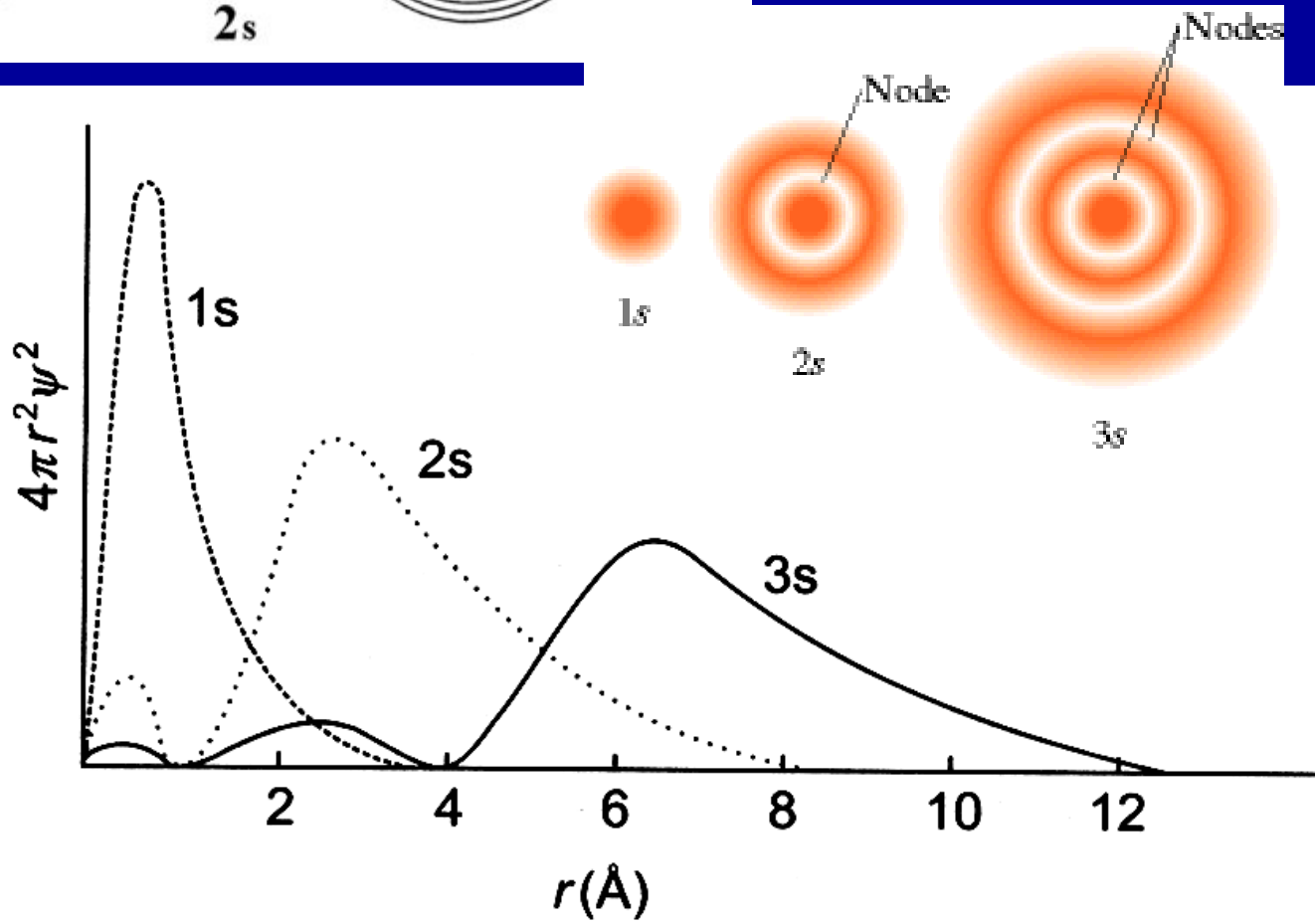
$4\pi r^2 R_{n,l}^2(r)$ = Radial distribution function

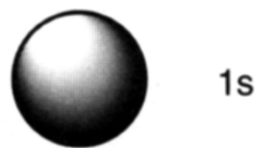
r_{\max} = the most probable radius
for 1s

$r_{\max} = a_0$ Bohr's radius

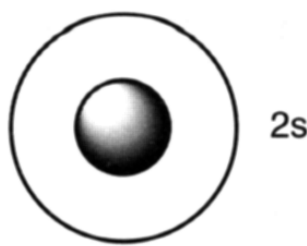


$4\pi r^2 R^2_{n,l}(r) = \text{Radial distribution function}$

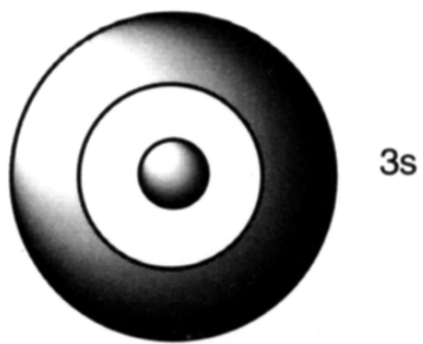




1s



2s



3s

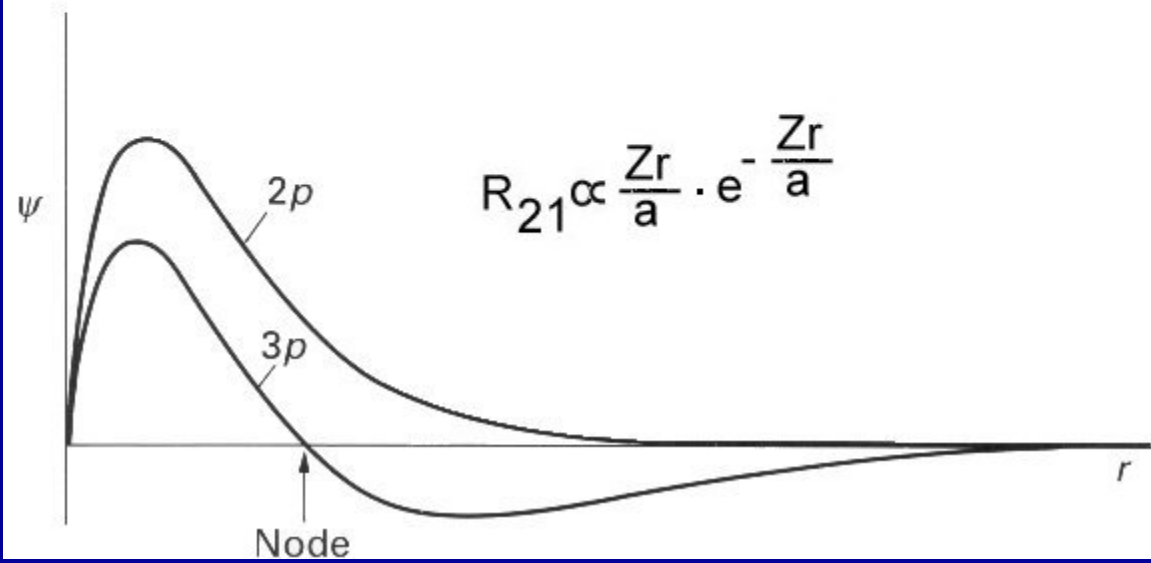
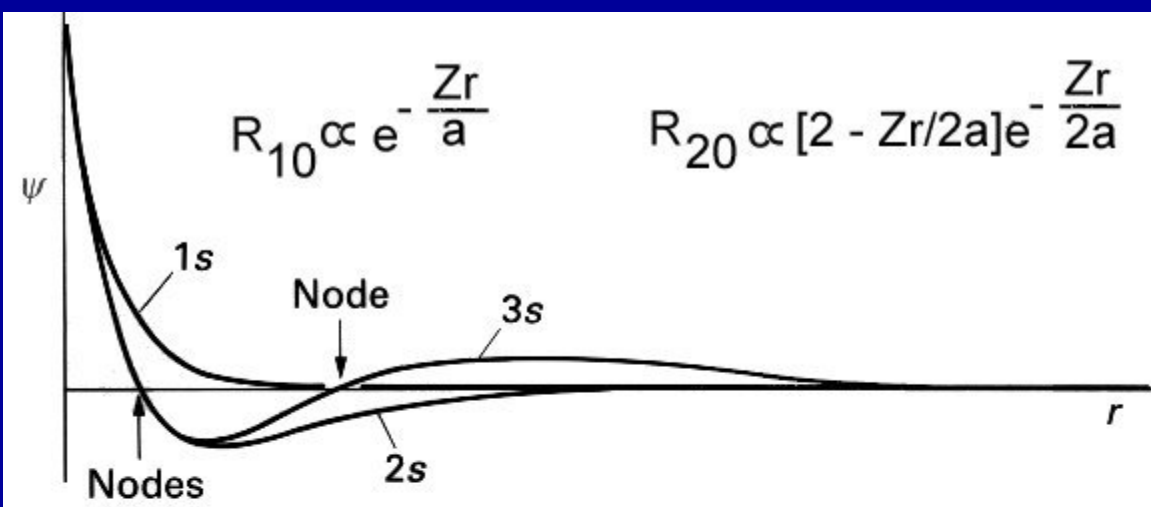


2p_x

nodal surface

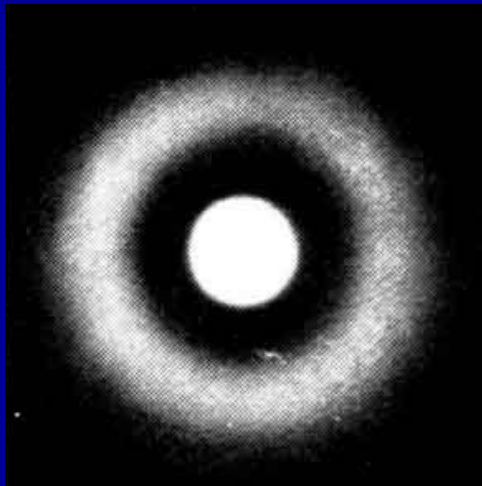


3p_x



Nodes

- Wavefunction changes sign
- Radial distribution function has zero value

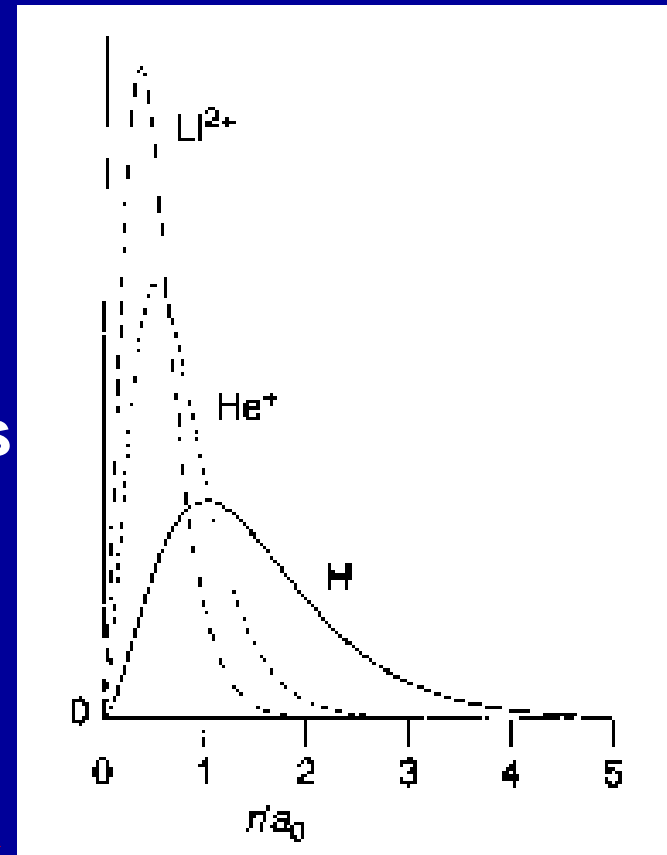


$$\text{Number of nodes} = n - l - 1$$

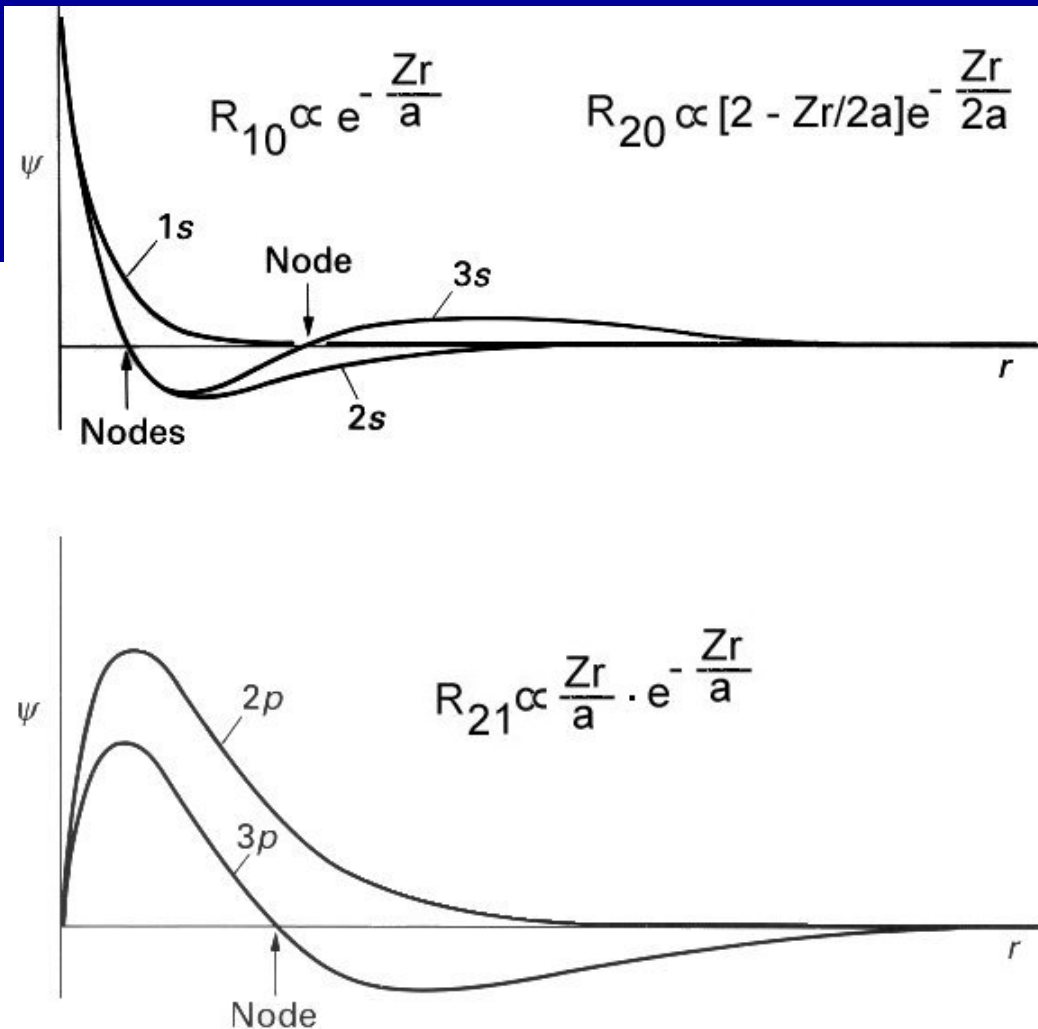
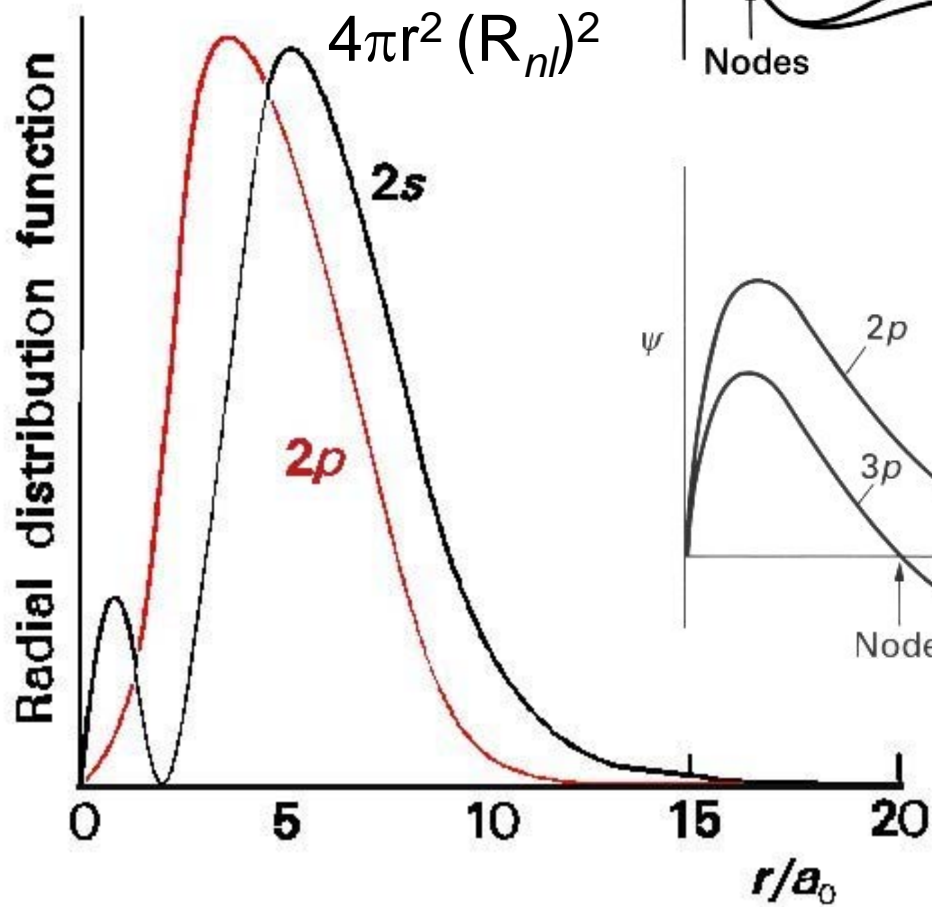
Effect of Z Radial Functions

$$R_{n,1}(r) = 2\sqrt{\left(\frac{Z}{a_0}\right)^3} \exp\left(-\frac{Zr}{a_0}\right)$$

Radial distribution function 1s



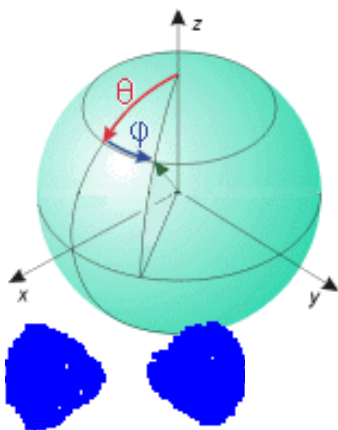
With increasing nucleus charge the maximum of radial distribution function approaches closer to the nucleus



Angular Wave Function

Angular wave function gives the shape of orbitals

The same for all values of n



$$\psi_{1,0}^2$$

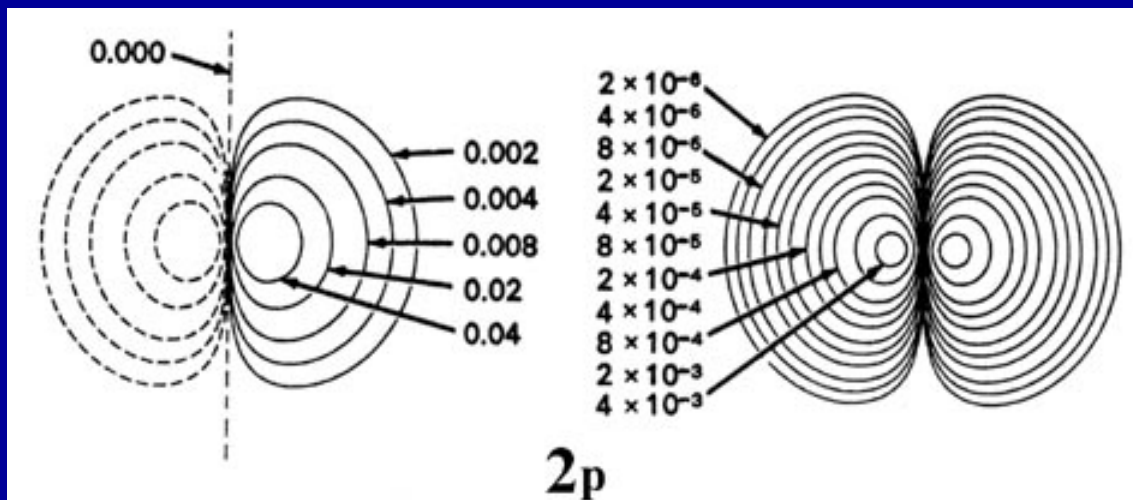
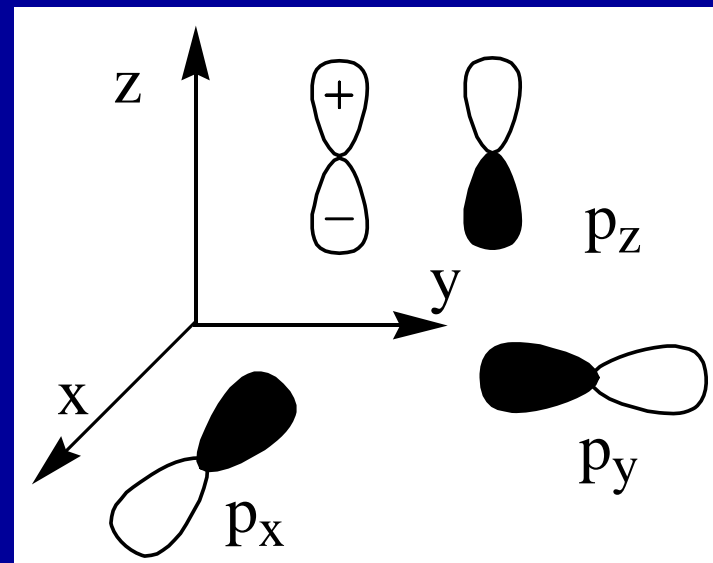
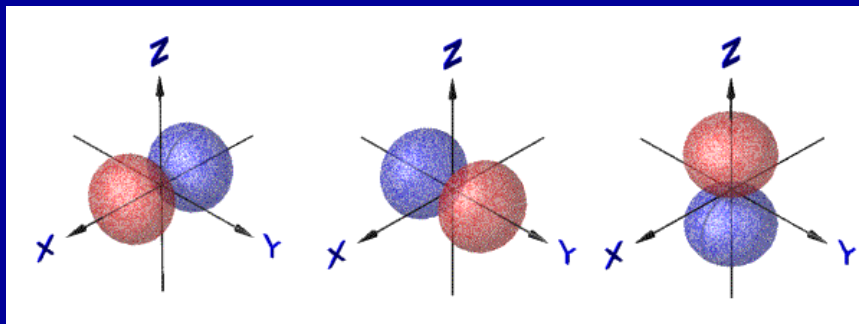


$$\psi_{1,\pm 1}^2$$

$$Y_{1,0} = (1/4\pi)^{1/2} 3^{1/2} \cos\theta$$

$$Y_{1,\pm 1} = (1/4\pi)^{1/2} (\mp 3/2)^{1/2} \sin\theta e^{\pm i\phi}$$

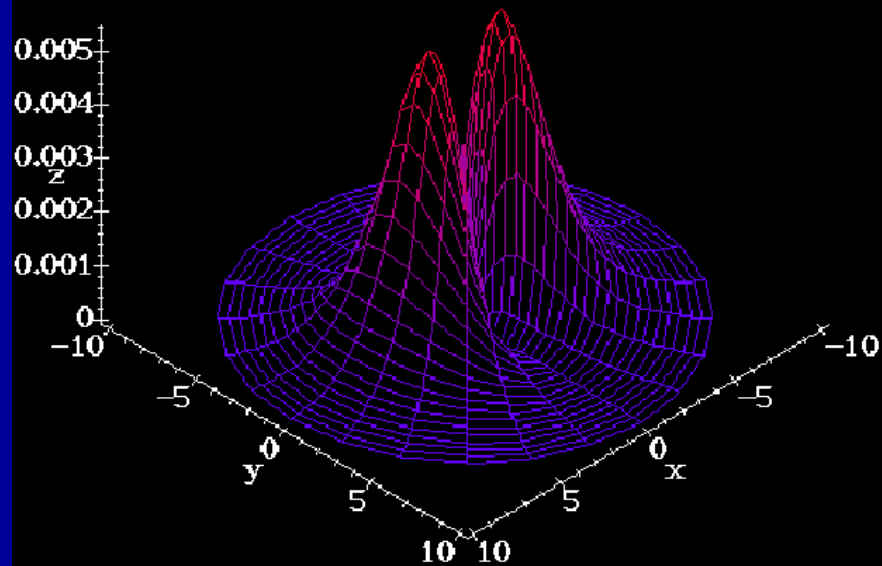
$n = 2, l = 1, m = 1, 0, -1$ p - Orbitals



2p - orbitals

2p orbital: $n=2, l=1, m=0$

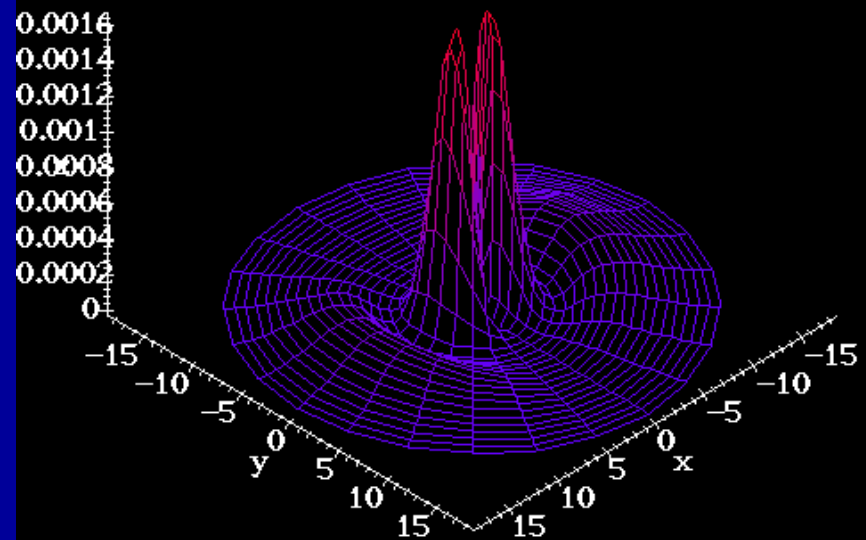
$$n = 2, l = 1, m = 0$$

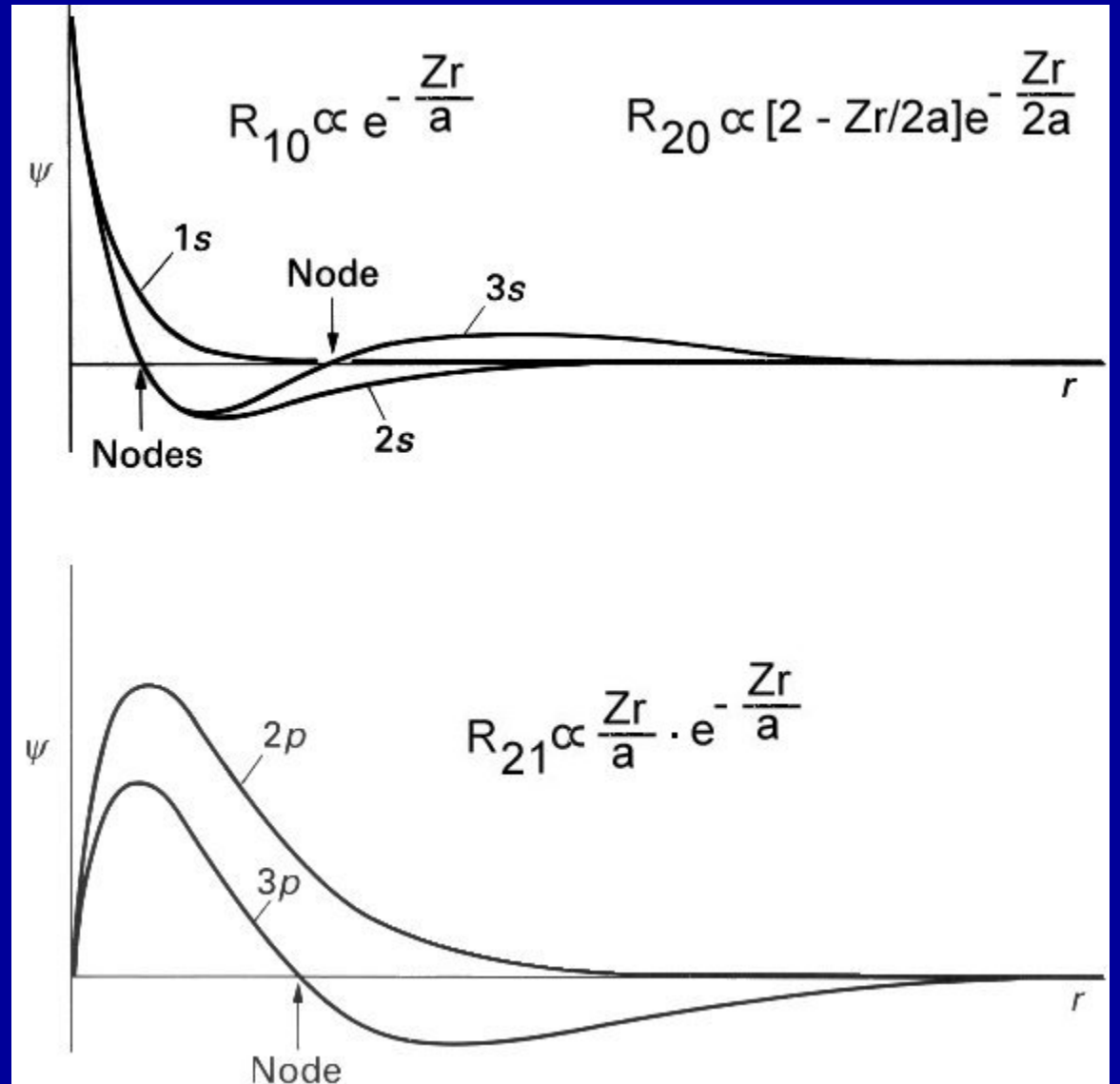
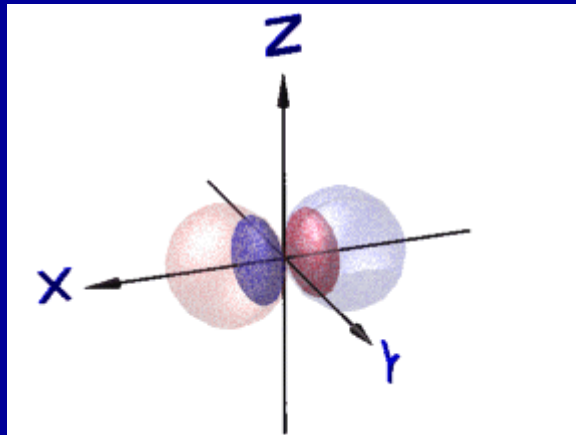


3p - orbitals

3p orbital: $n=3, l=1, m=0$

$$n = 3, l = 1, m = 0$$

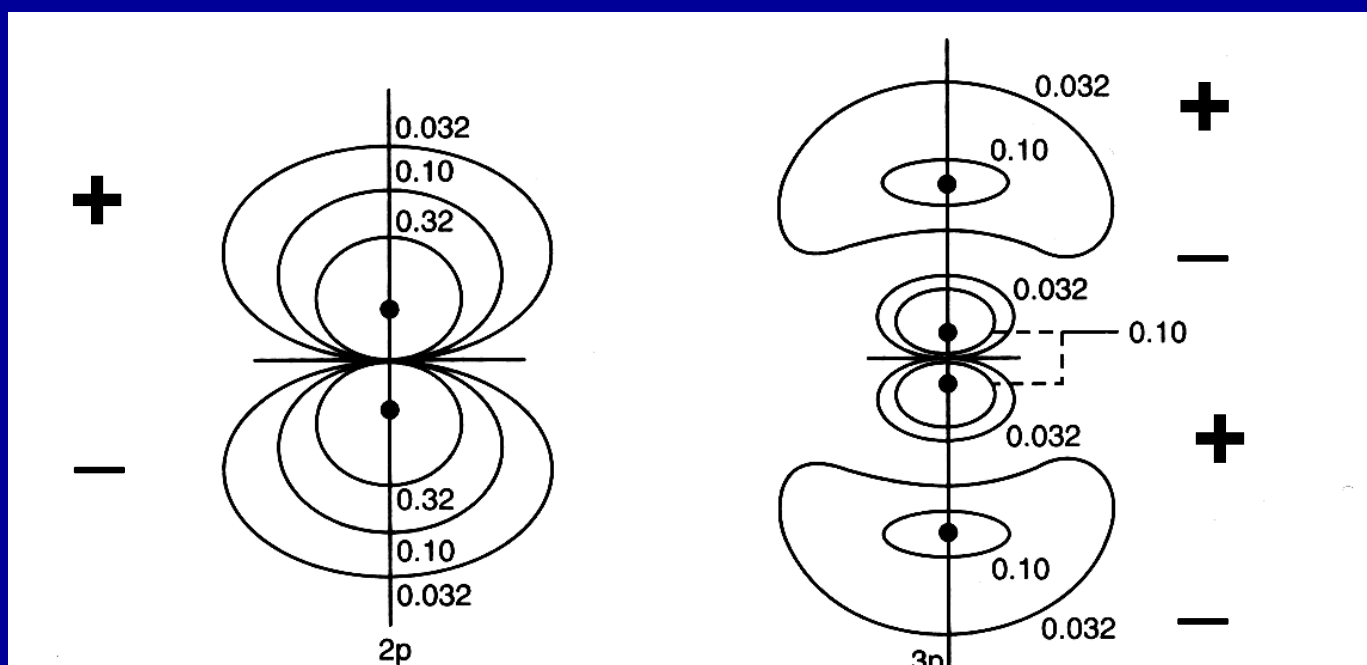




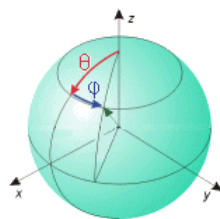
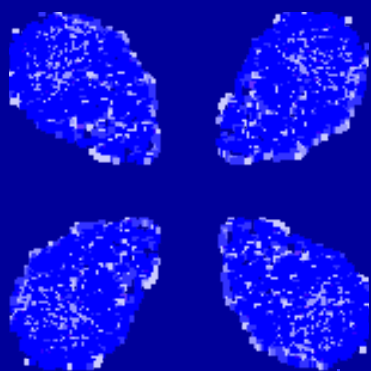
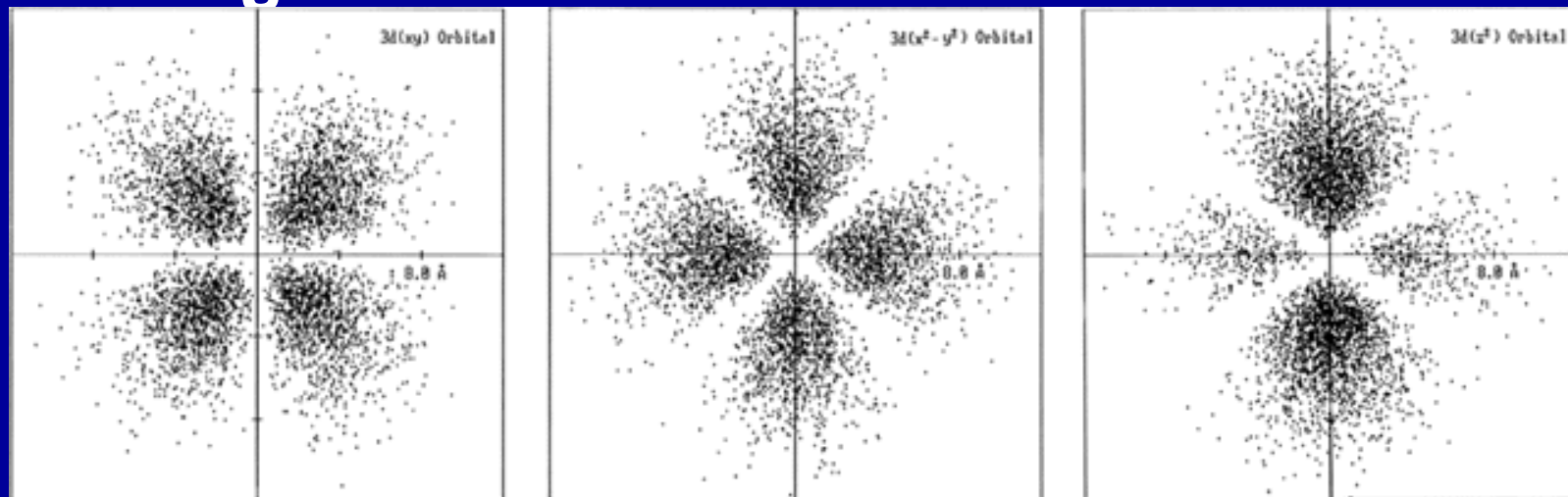
2p - orbitals

3p - orbitals

Wave Function = Radial \times Angular



Angular Wave Function of d-Orbitals



$$\psi_{2,0}^2$$



$$\psi_{2,\pm 1}^2$$



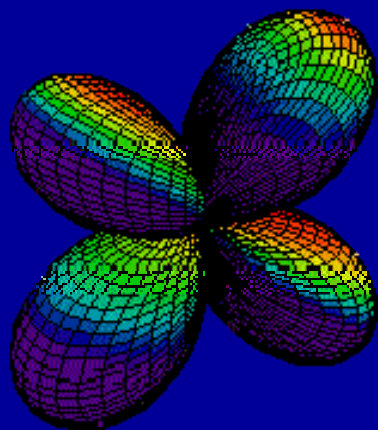
$$\psi_{2,\pm 2}^2$$

$$Y_{2,0} = (1/4\pi)^{1/2} (5/4)^{1/2} (3\cos^2\theta - 1)$$

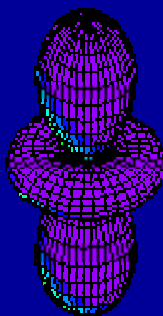
$$Y_{2,\pm 1} = (1/4\pi)^{1/2} (\mp 15/4)^{1/2} \cos\theta \sin\theta e^{\pm i\phi}$$

$$Y_{2,\pm 2} = (1/4\pi)^{1/2} (15/8)^{1/2} \sin^2\theta e^{\pm 2i\phi}$$

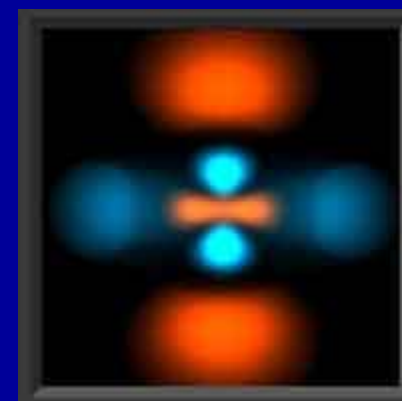
d - Orbitals



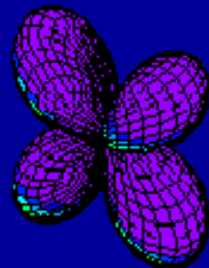
d_{z^2}



$d_{x^2-y^2}$



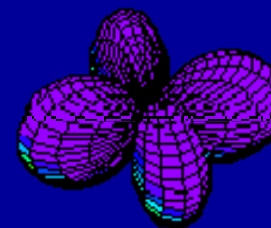
d_{yz}



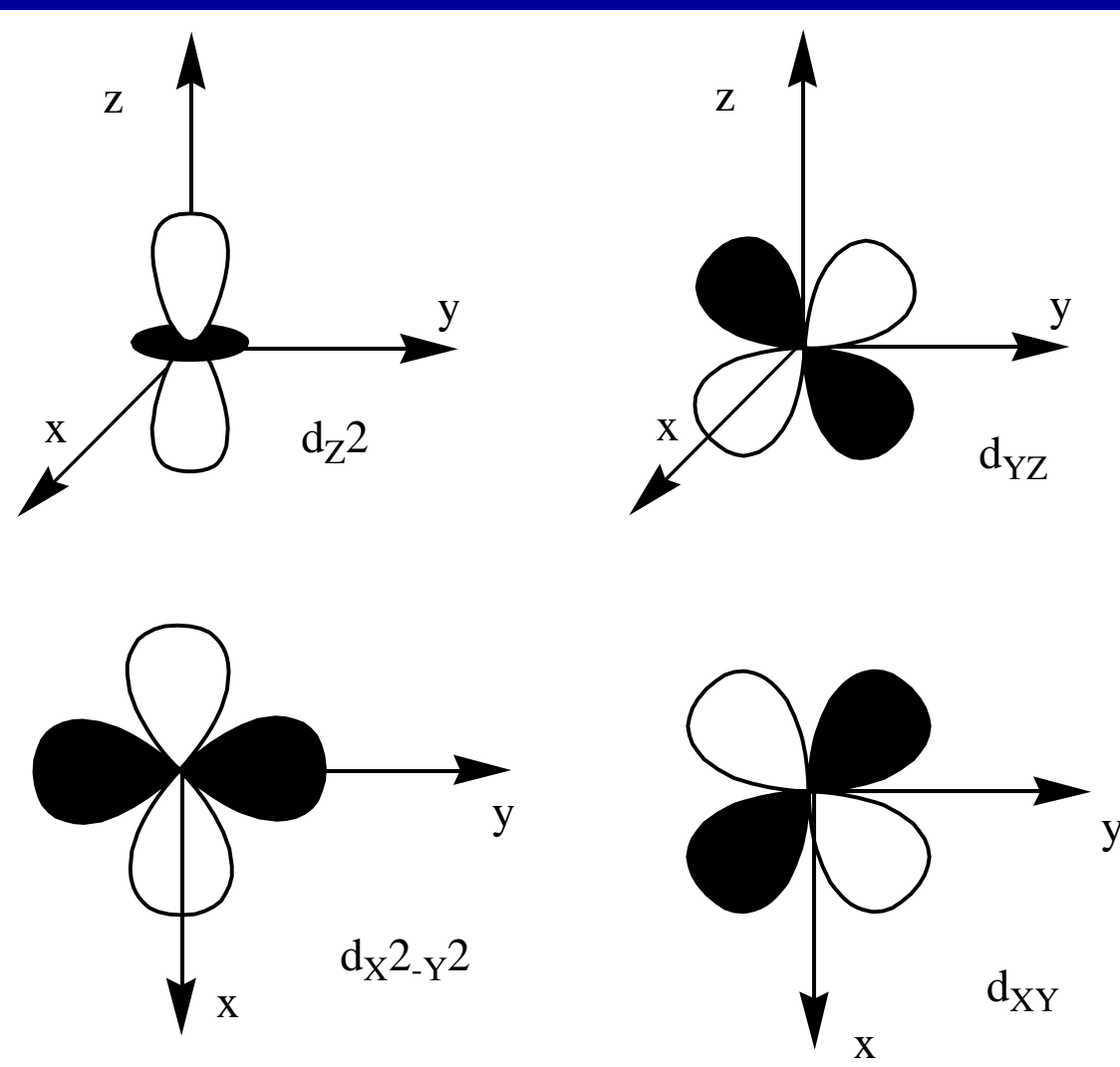
d_{xz}



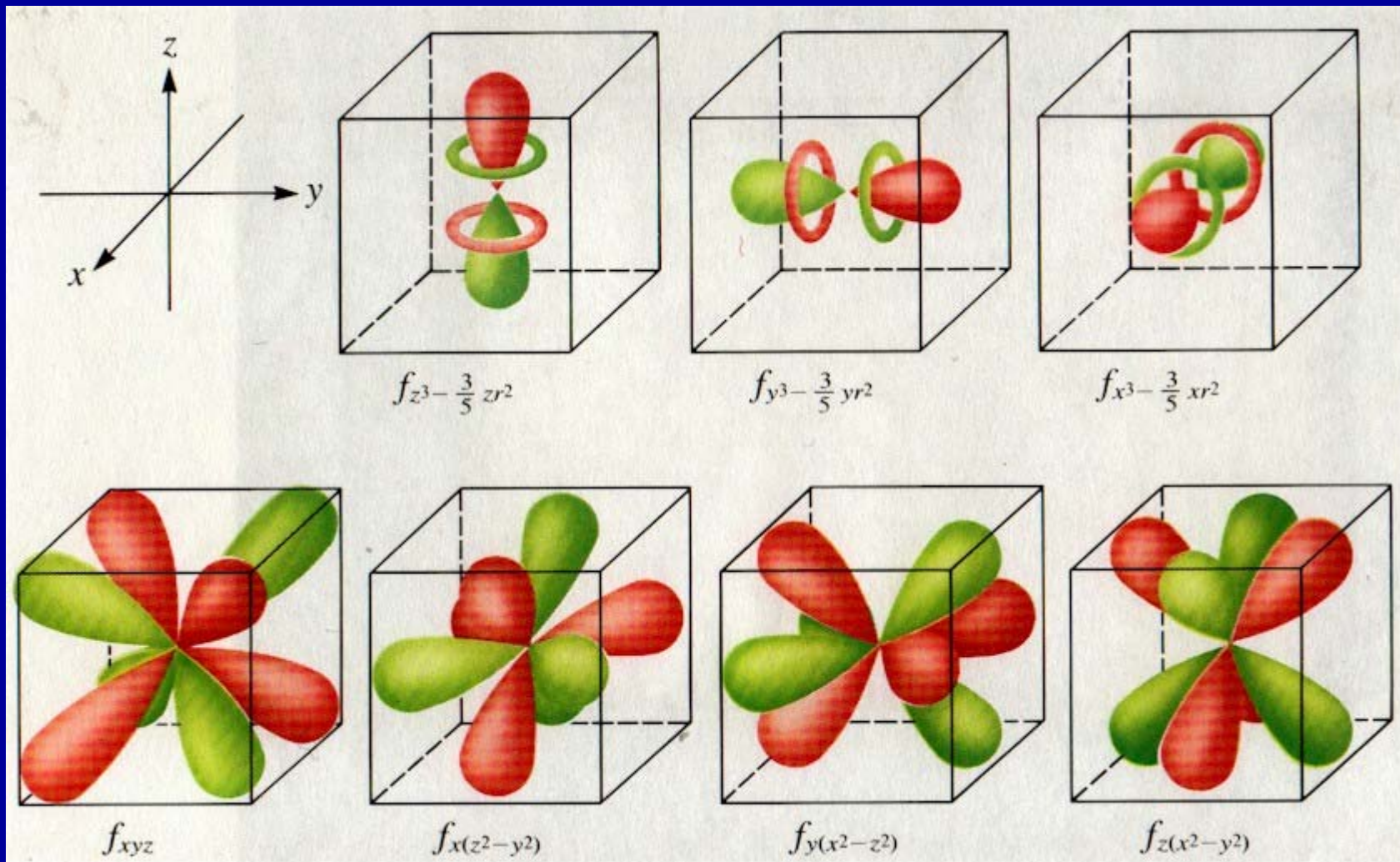
d_{xy}



d - Orbitals



f - Orbitals



Nodes

Spherical nodes = $n - l - 1$

for s, p, d, f,....

Radial wave function

Nodal planes

Angular wave function:

Orbital no.

s

0

p

1

d

2

f

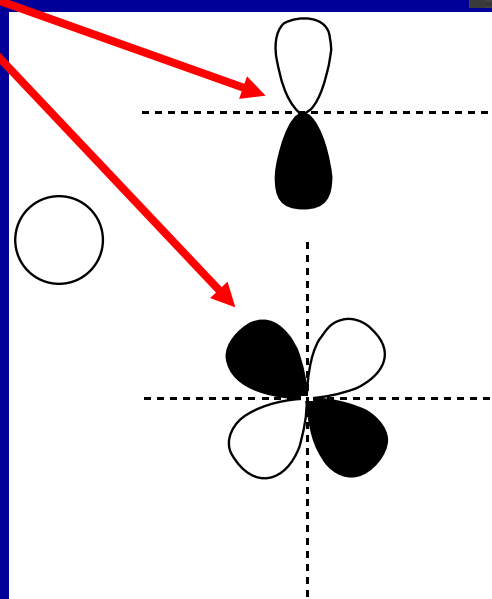
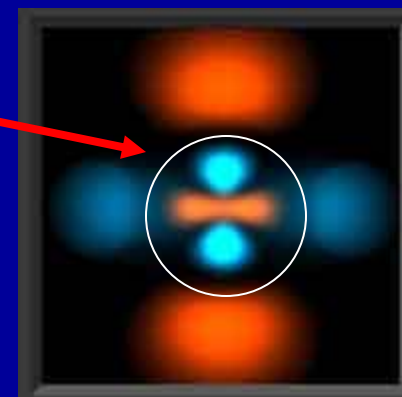
3

.

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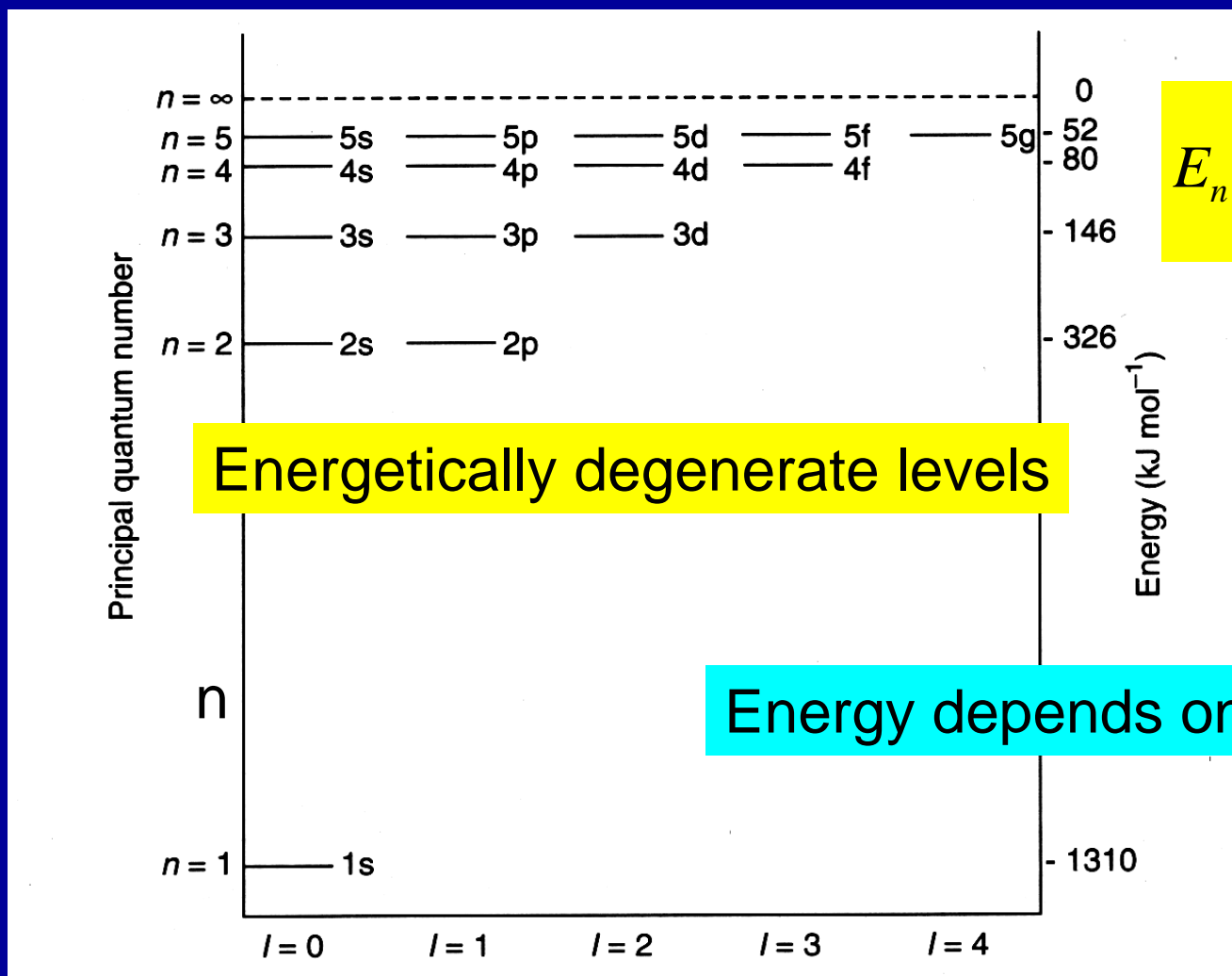
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Only s-orbitals have non-zero value of wave function at the nucleus

Energy of H-Atom Orbitals

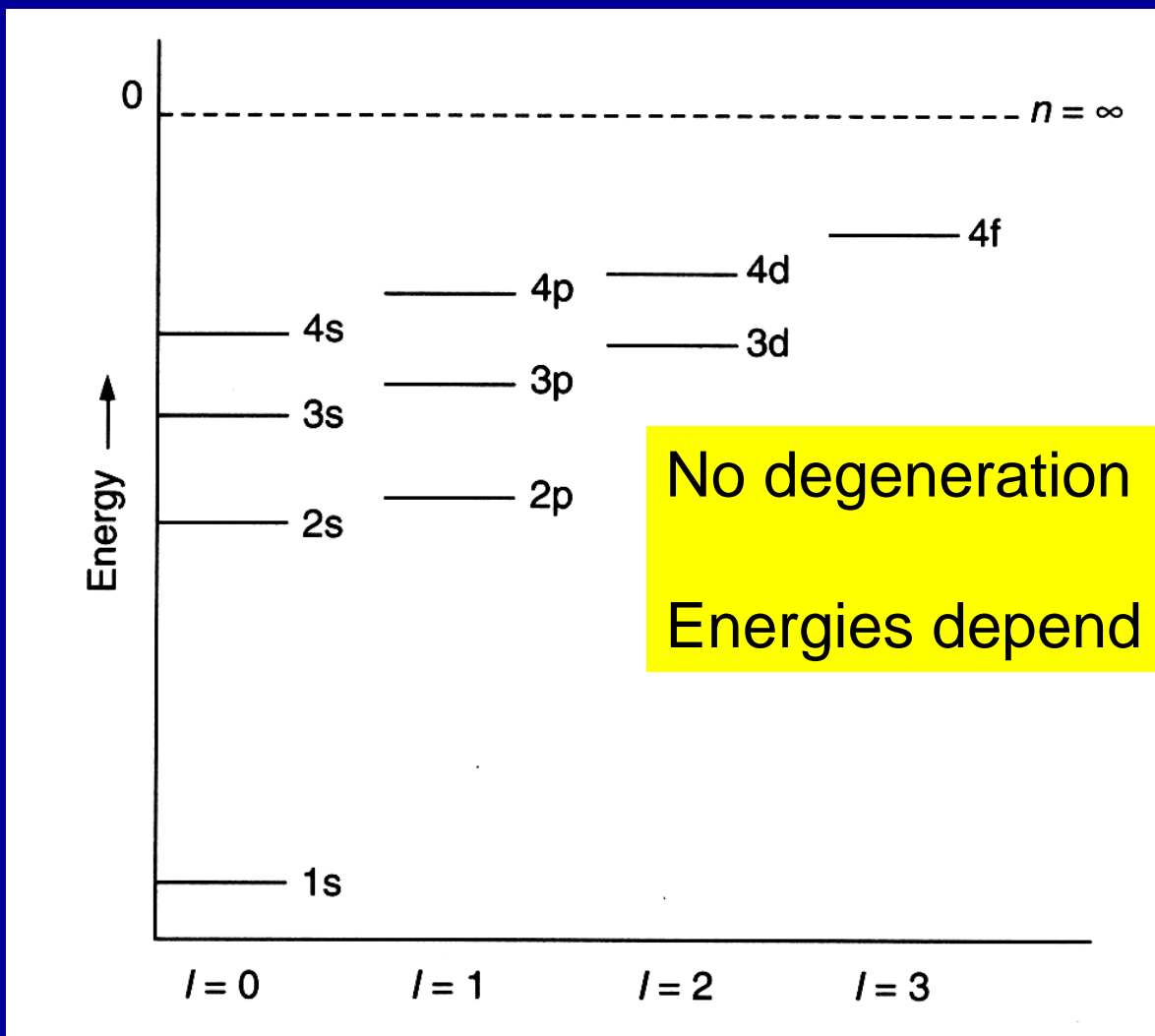


$$E_n = -\frac{N_A \mu e^4 Z^2}{8 \epsilon_0^2 h^2 n^2}$$

Energetically degenerate levels

Energy depends only on n

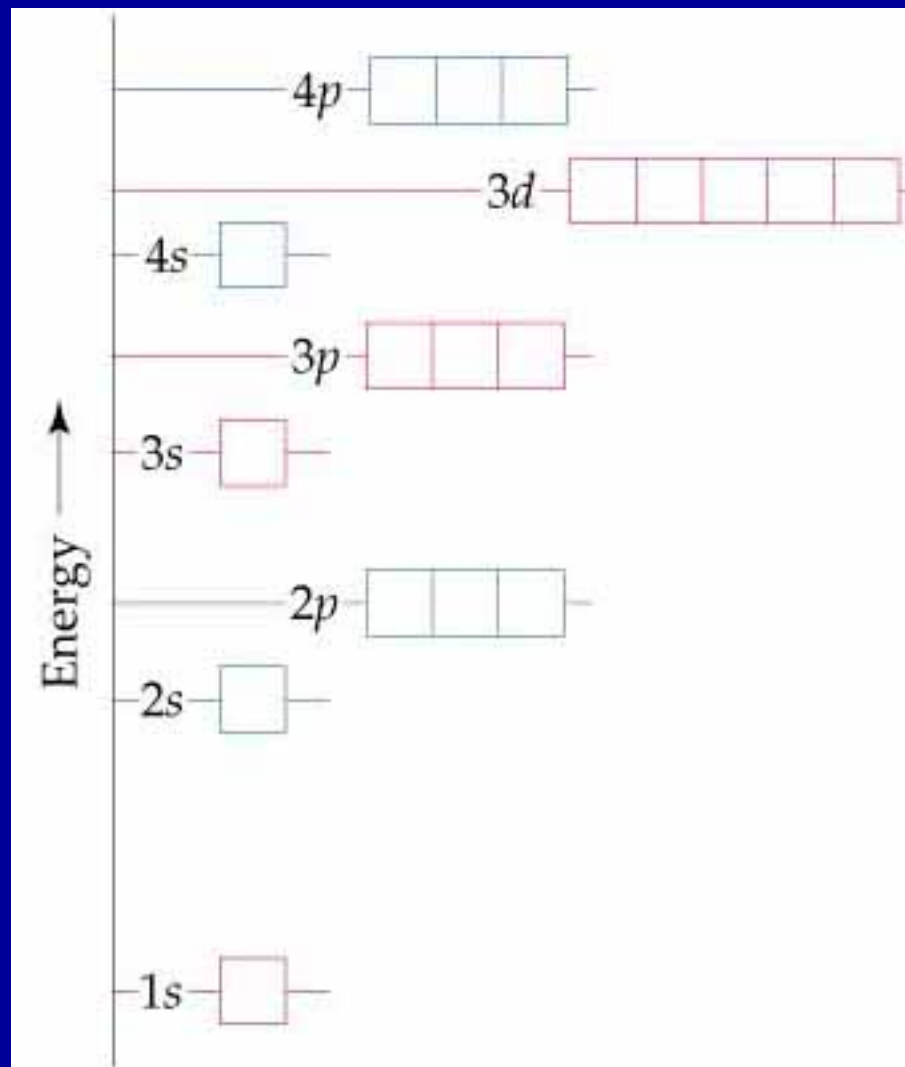
Energy Levels in Many-Electron Atoms



No degeneration

Energies depend on n and l

Energy Levels in Many-Electron Atoms



More stable orbital
has a lower energy

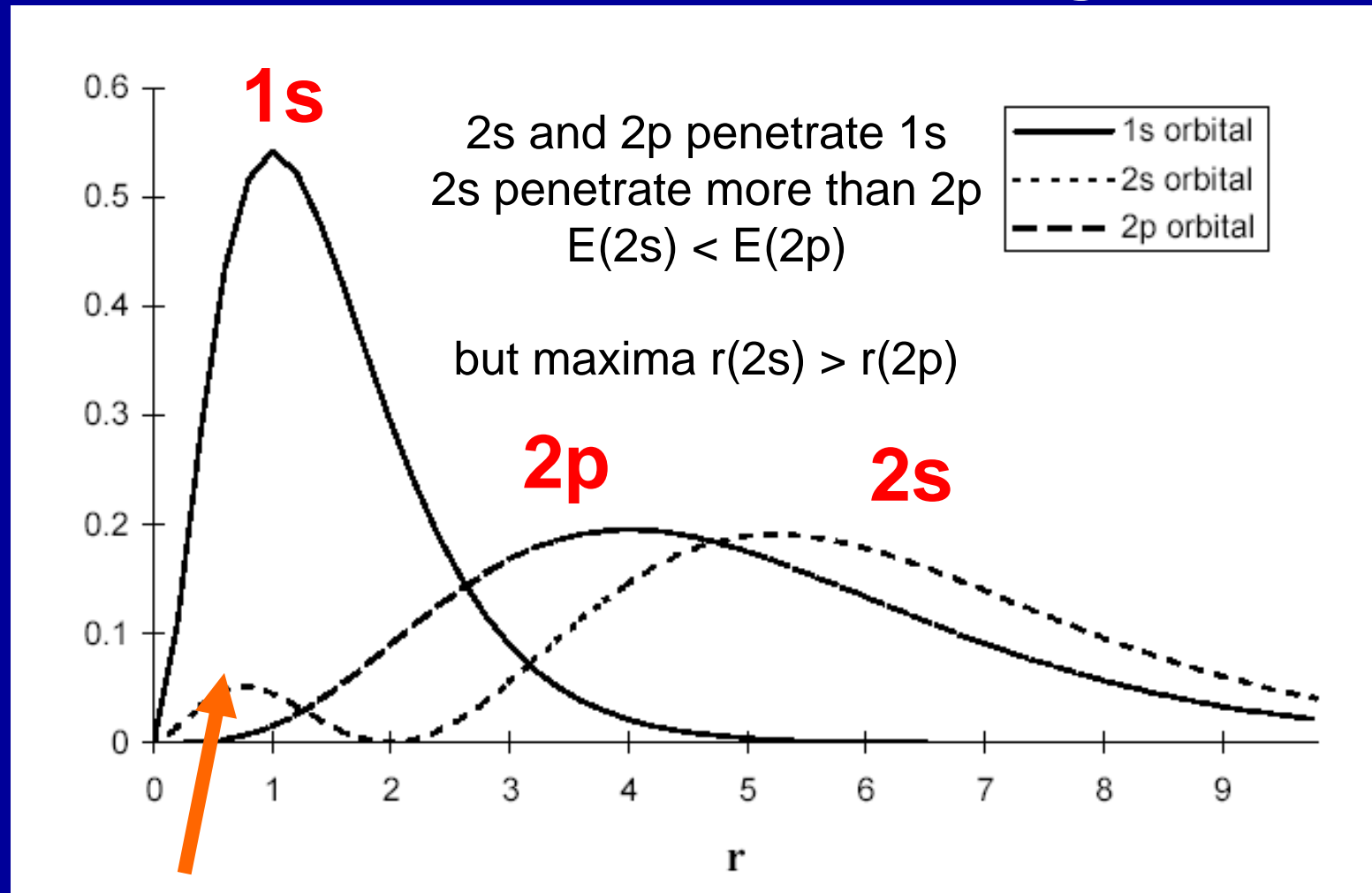
Madelung's Rule
(up to Ca)

1. Lower for $(n + l)$
2. When $n + l$ same
lower n

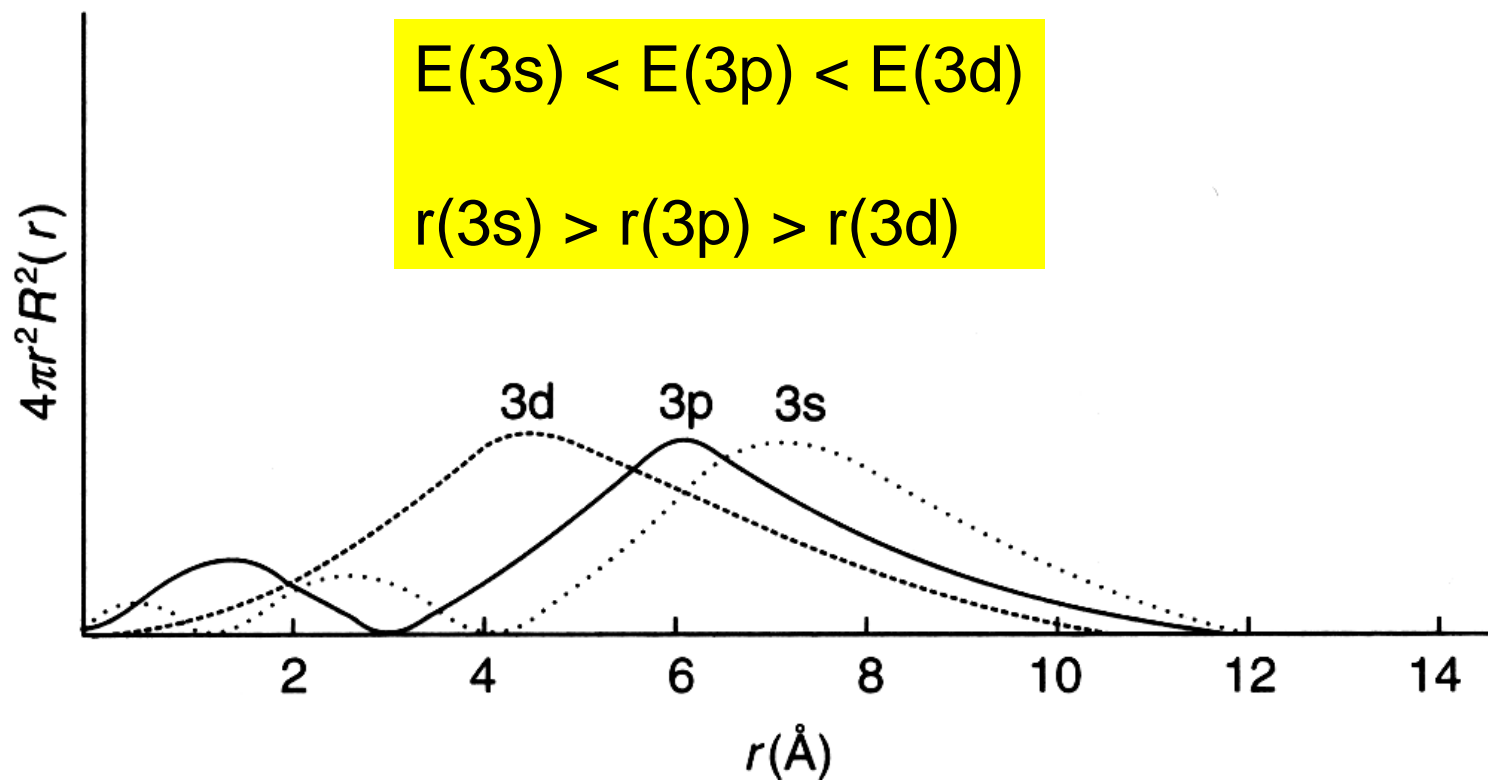
3p 4s

4p 3d

Many-Electron Atoms Penetration and Screening



Relative Energies of s, p, d Orbitals



Slater's Orbitals

Orbitals for many-electron atoms - approximate

- orbitals (wave functions) of hydrogen type
- angular part: same as for H
- radial part:

$$R(r) = N r^{n^*-1} \exp(-Z^* r/n^*)$$

Z^* = A charge acting on an electron

= Nucleus charge ($Z+$) – charge of other electrons

n^* = effective quant. number (for K, L, M = n)

$$E_i = -N (Z_i^*/n_i) \quad N = 1313 \text{ kJ mol}^{-1}$$

Effective Nucleus Charge

$$Z^* = Z - \sigma$$

σ = screening constant, sum for all electrons

(1s)(2s,2p)(3s,3p)(3d)(4s,4p)(4d)(4f)(5s,5p)(5d)(5f)...

Slater's rules

e on the right does not screen, no contribution to σ

Within a group screens 0.35 (1s only 0.30)

$n - 1$ (s,p) screens 0.85

$n - 2$ and lower screens 1.00

If an electron is in d or f, all on the left screens 1.00

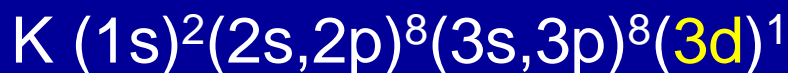
Effective Nucleus Charge

Z^* = Effective Nucleus Charge

$$Z^* = Z - \sigma$$

A charge acting on an electron

= Nucleus charge (Z^+) – charge of other electrons



$$\sigma(3d) = 0 \times (0.35) + 8 \times 1.00 + 10 \times 1.00 = 18$$

$$Z^* = 19 - 18 = 1$$



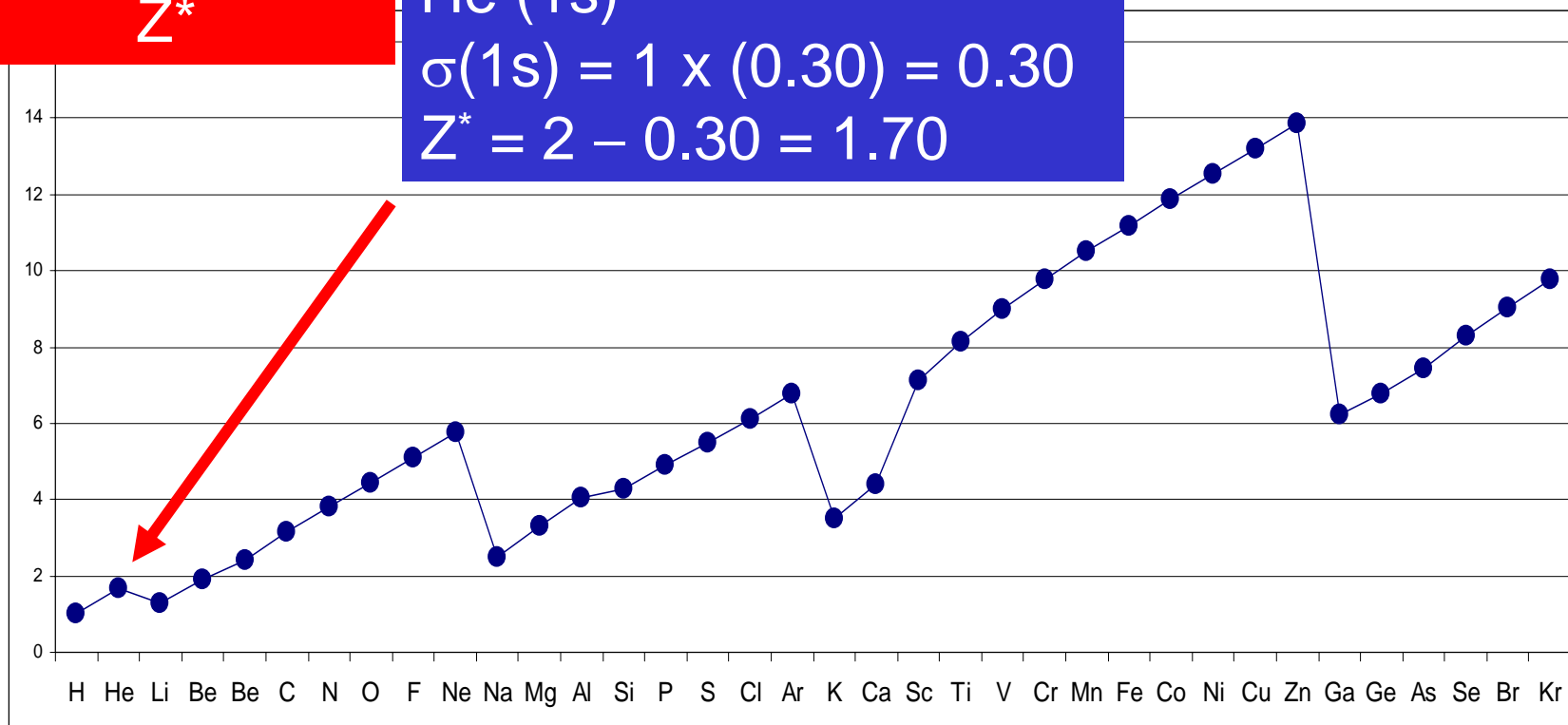
$$\sigma(4s) = 0 \times (0.35) + 8 \times 0.85 + 10 \times 1.00 = 16.8$$

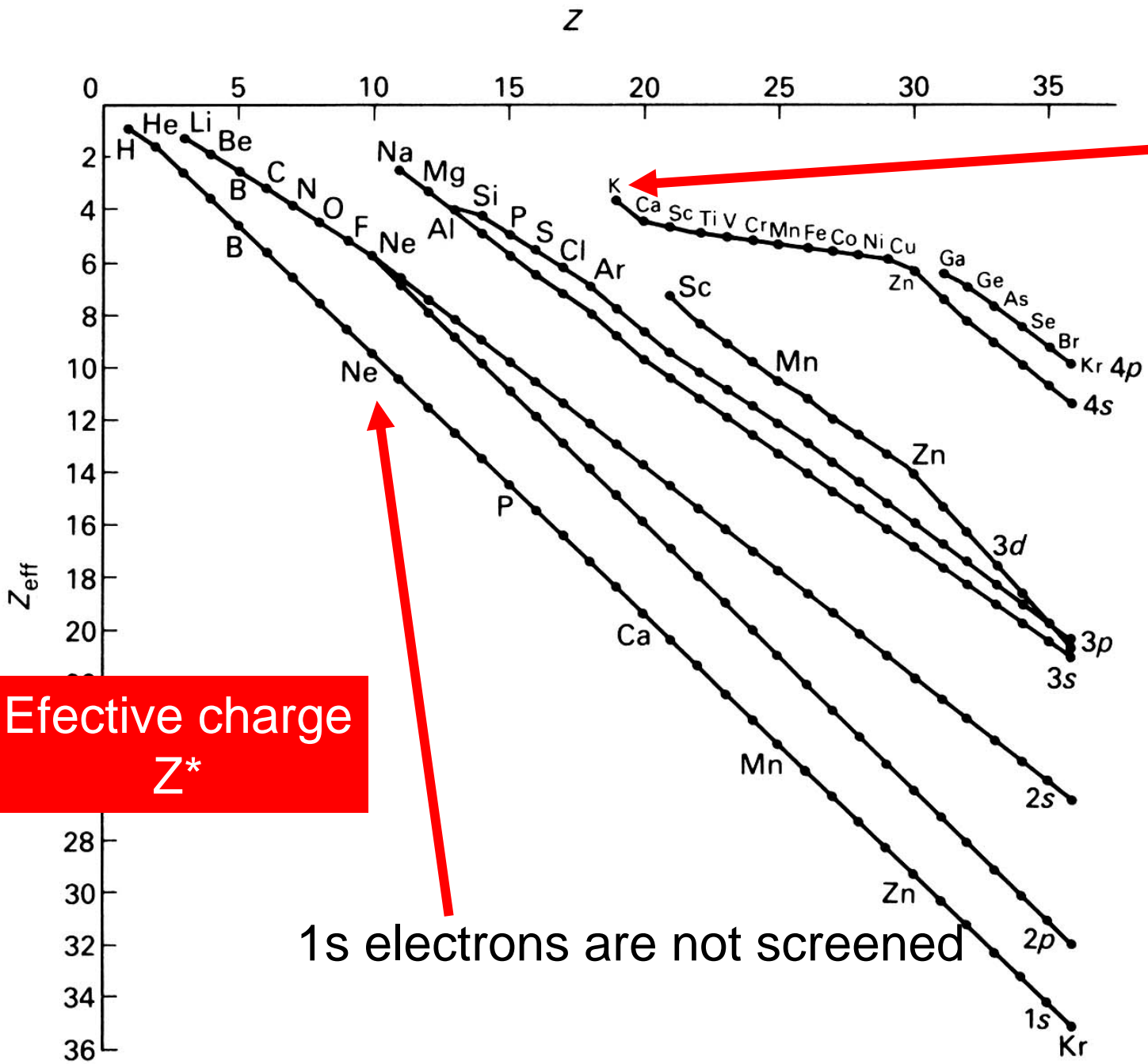
$$Z^* = 19 - 16.8 = 2.2$$

Effective Nucleus Charge

Effective charge
 Z^*

He $(1s)^2$
 $\sigma(1s) = 1 \times (0.30) = 0.30$
 $Z^* = 2 - 0.30 = 1.70$



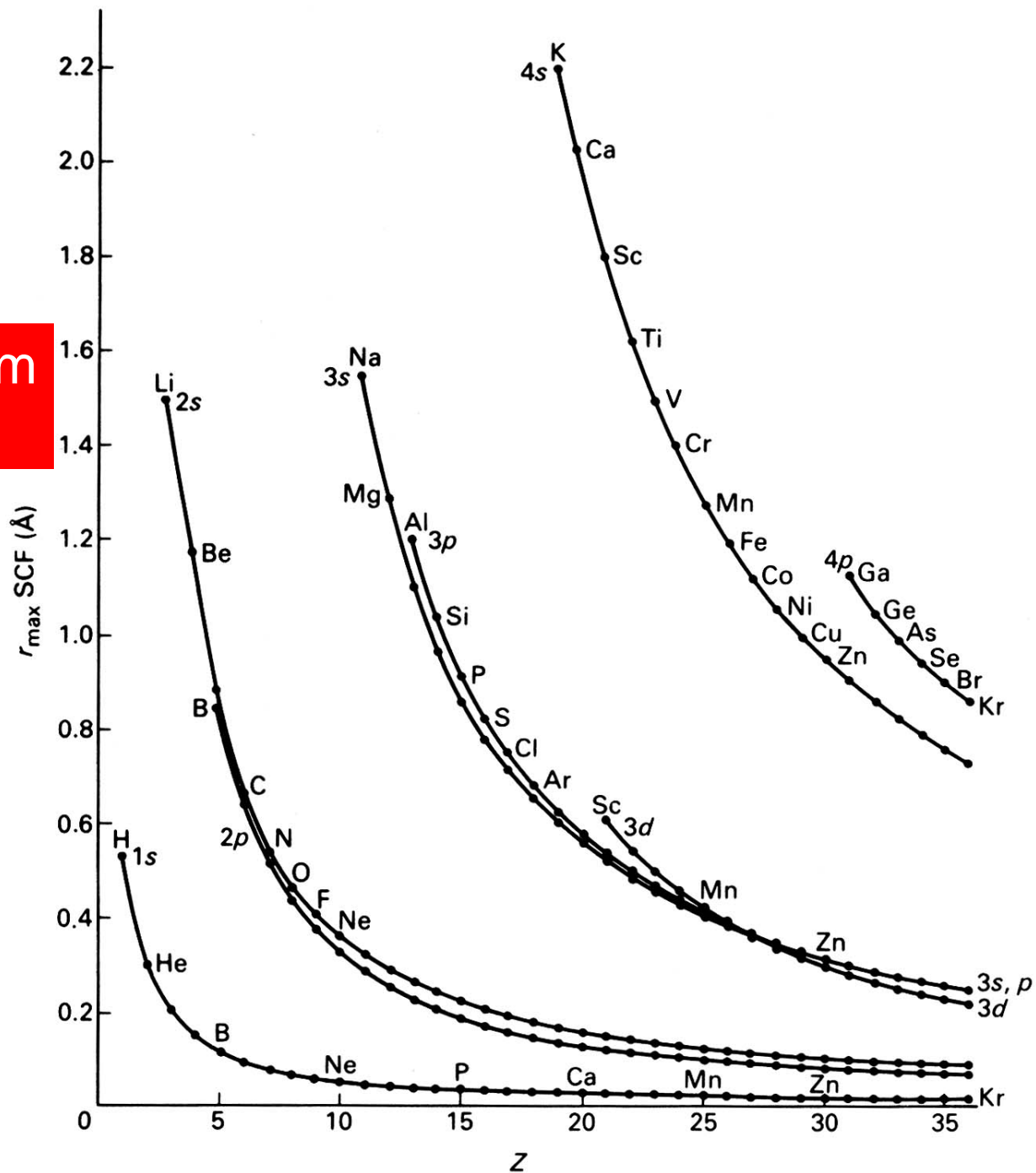


Other
electrons
are
screened

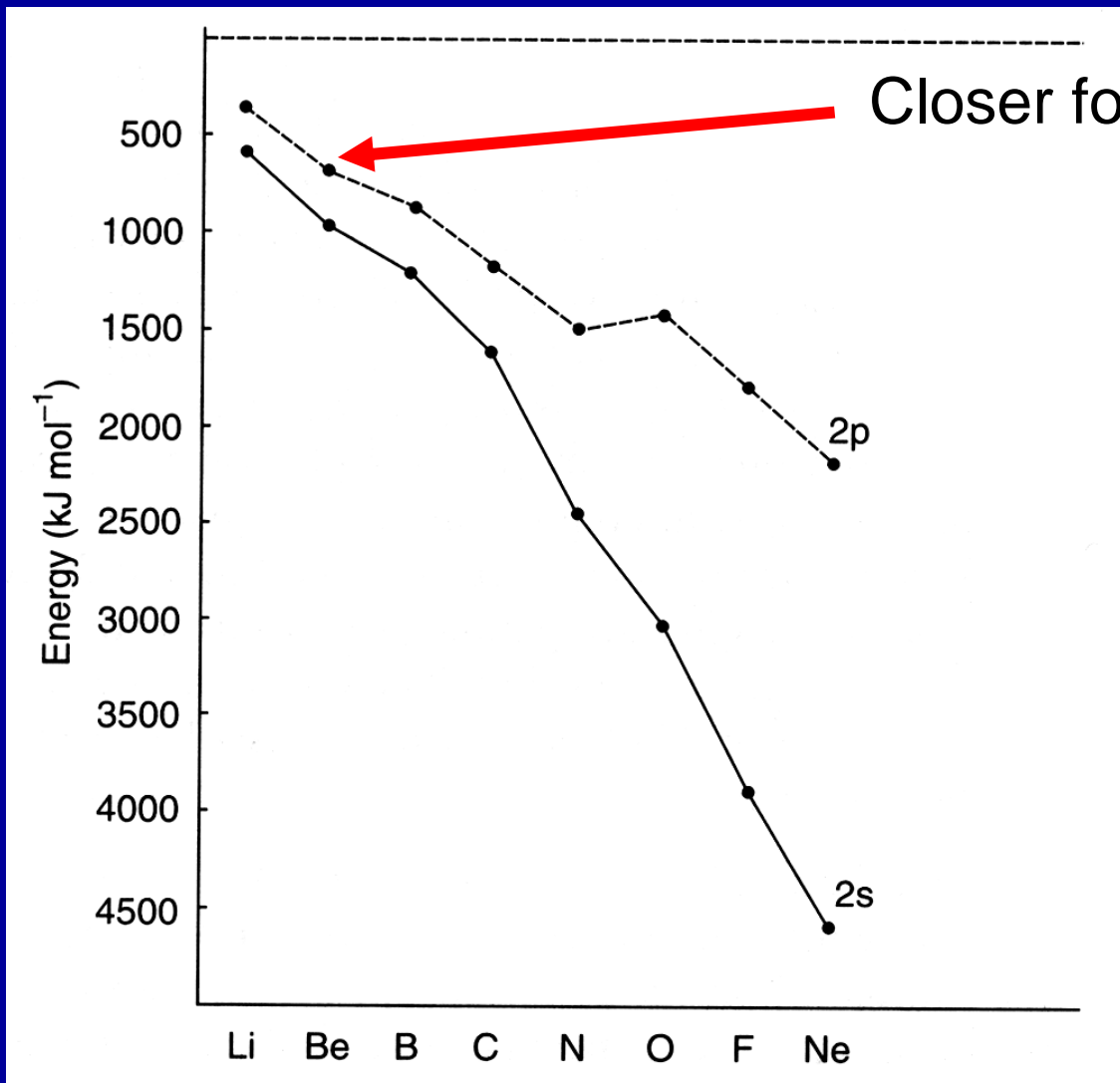
Radius of maximum electron density

$$r(2s) > r(2p)$$

$$r(3s) \sim r(3p)$$



Energies of 2s and 2p Orbitals



Electron Configurations of Ground State Atoms

Aufbau Principle:

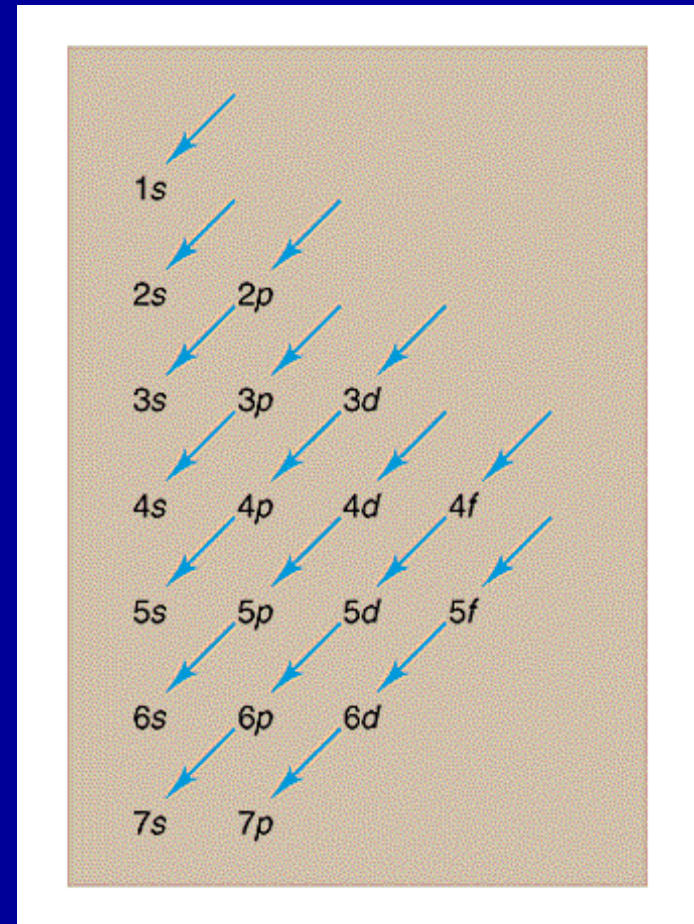
Electron levels are filled by electrons in the order of increasing energy, to maintain the lowest atom energy

Pauli Principle:

Two electrons cannot have all 4 quantum numbers the same

Hund's Rule:

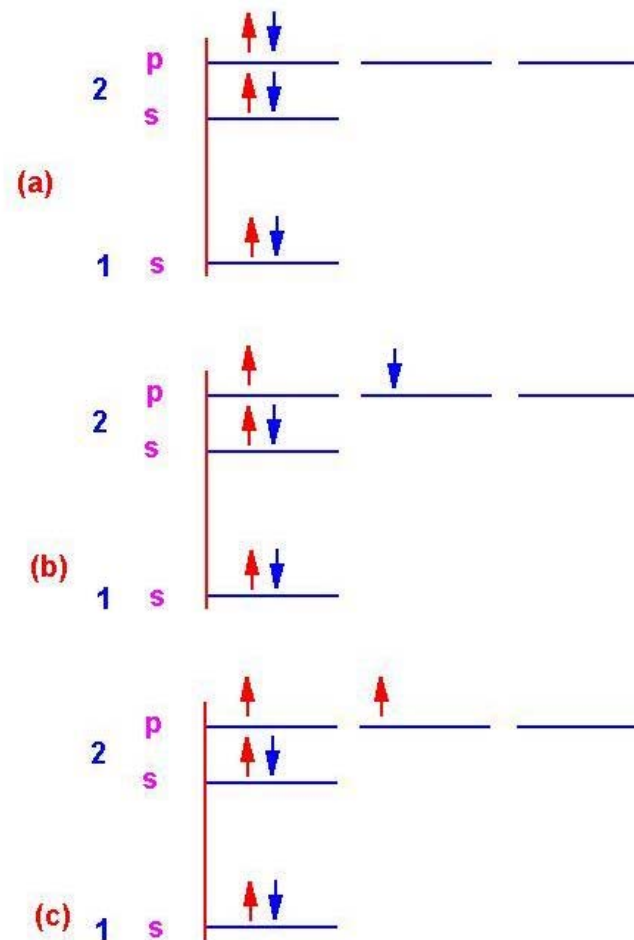
In degenerate orbitals, the state with maximum unpaired electrons is the most stable



1												18					
1 H 1.0079	2											13	14	15	16	17	2 He 4.0026
3 Li 6.941	4 Be 9.0122											5 B 10.811	6 C 12.011	7 N 14.007	8 O 15.999	9 F 18.998	10 Ne 20.180
11 Na 22.990	12 Mg 24.305	3	4	5	6	7	8	9	10	11	12	13 Al 26.982	14 Si 28.086	15 P 30.974	16 S 32.065	17 Cl 35.453	18 Ar 39.948
19 K 39.098	20 Ca 40.078	21 Sc 44.956	22 Ti 47.867	23 V 50.942	24 Cr 51.996	25 Mn 54.938	26 Fe 55.845	27 Co 58.933	28 Ni 58.693	29 Cu 63.546	30 Zn 65.38	31 Ga 69.723	32 Ge 72.64	33 As 74.922	34 Se 78.96	35 Br 79.904	36 Kr 83.798
37 Rb 85.468	38 Sr 87.62	39 Y 88.906	40 Zr 91.224	41 Nb 92.906	42 Mo 95.96	43 Tc (98)	44 Ru 101.07	45 Rh 102.91	46 Pd 106.42	47 Ag 107.87	48 Cd 112.41	49 In 114.82	50 Sn 118.71	51 Sb 121.76	52 Te 127.60	53 I 126.90	54 Xe 131.29
55 Cs 132.91	56 Ba 137.33	57-71 *	72 Hf 178.49	73 Ta 180.95	74 W 183.84	75 Re 186.21	76 Os 190.23	77 Ir 192.22	78 Pt 195.08	79 Au 196.97	80 Hg 200.59	81 Tl 204.38	82 Pb 207.2	83 Bi 208.98	84 Po (209)	85 At (210)	86 Rn (222)
87 Fr (223)	88 Ra (226)	89-103 #	104 Rf (261)	105 Db (262)	106 Sg (266)	107 Bh (264)	108 Hs (270)	109 Mt (268)	110 Ds (281)	111 Rg (272)	112 Uub (285)	113 Uut (284)	114 Uuq (289)	115 Uup (288)	116 Uuh (291)		118 Uuo (294)
* Lanthanide series			57 La 138.91	58 Ce 140.12	59 Pr 140.91	60 Nd 144.24	61 Pm (145)	62 Sm 150.36	63 Eu 151.96	64 Gd 157.25	65 Tb 158.93	66 Dy 162.50	67 Ho 164.93	68 Er 167.26	69 Tm 168.93	70 Yb 173.05	71 Lu 174.97
# Actinide series			89 Ac (227)	90 Th 232.04	91 Pa 231.04	92 U 238.03	93 Np (237)	94 Pu (244)	95 Am (243)	96 Cm (247)	97 Bk (247)	98 Cf (251)	99 Es (252)	100 Fm (257)	101 Md (258)	102 No (259)	103 Lr (262)

Electron Configurations of C

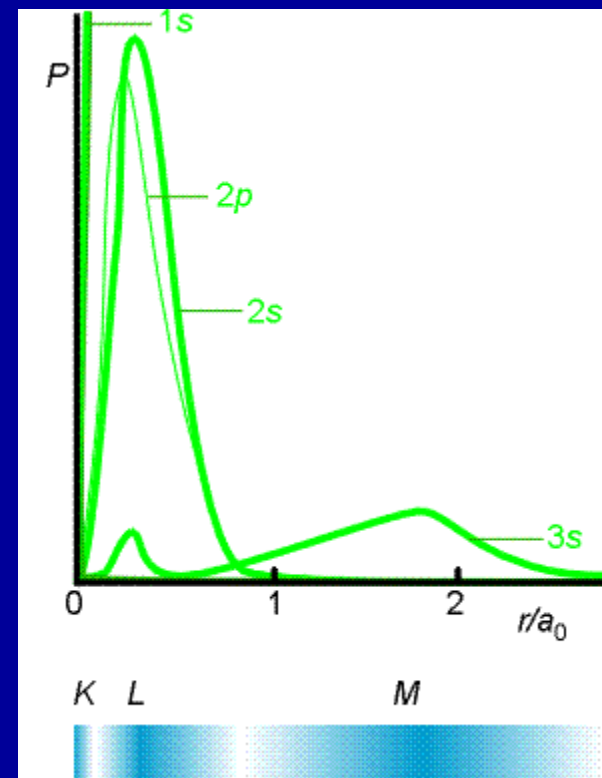
	1s	2s	2p		
a)	$\uparrow\downarrow$	$\uparrow\downarrow$	\uparrow	\downarrow	—
b)	$\uparrow\downarrow$	$\uparrow\downarrow$	$\uparrow\uparrow$	—	—
c)	$\uparrow\downarrow$	$\uparrow\downarrow$	\uparrow	\uparrow	—
d)	$\uparrow\downarrow$	$\uparrow\downarrow$	$\uparrow\downarrow$	—	—



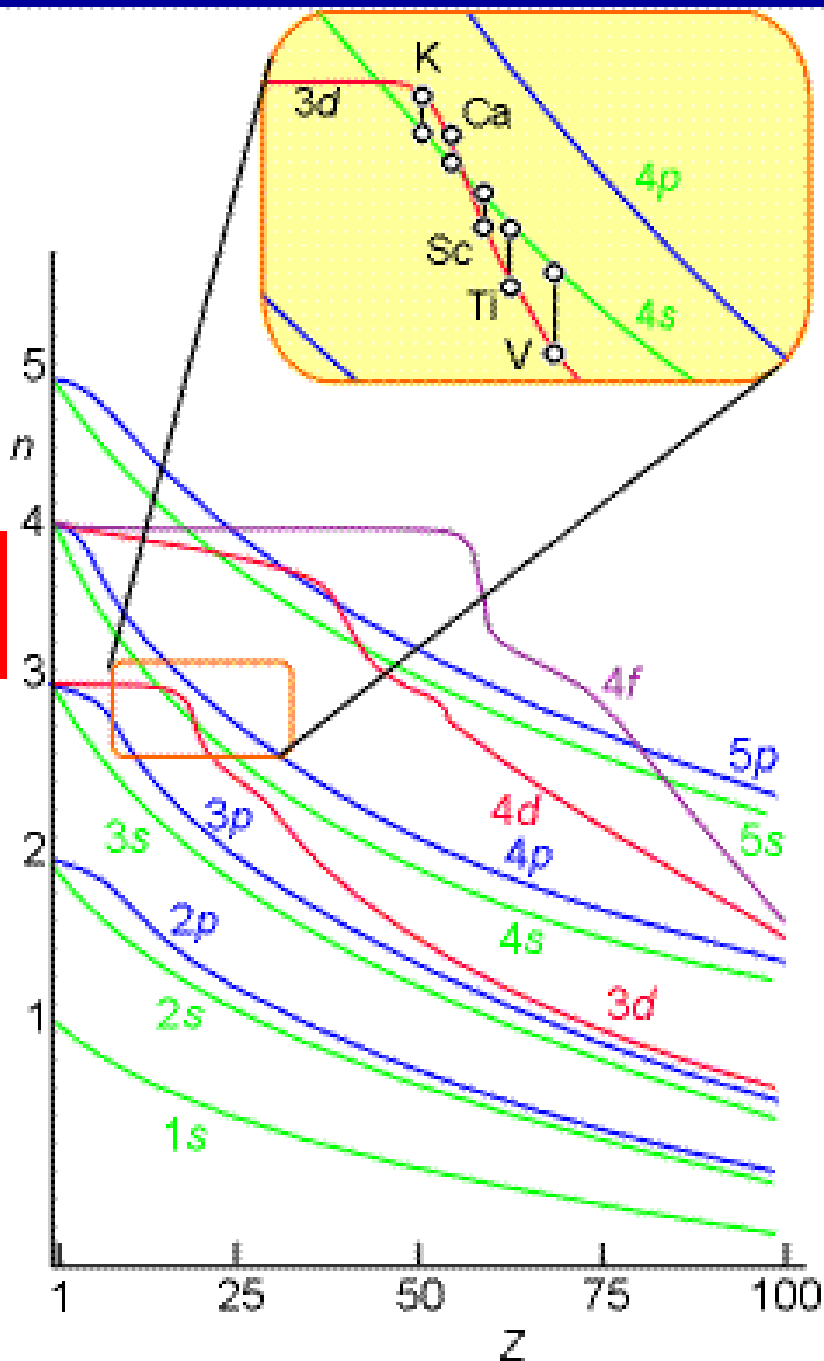
Electron Configurations of Valence Shell

(Ne)

	3s	3p
Na	\uparrow	
Mg	$\uparrow\downarrow$	
Al	$\uparrow\downarrow$	\uparrow
Si	$\uparrow\downarrow$	$\uparrow\uparrow$
P	$\uparrow\downarrow$	$\uparrow\uparrow\uparrow$
S	$\uparrow\downarrow$	$\uparrow\downarrow\uparrow\uparrow$
Cl	$\uparrow\downarrow$	$\uparrow\downarrow\uparrow\downarrow\uparrow$
Ar	$\uparrow\downarrow$	$\uparrow\downarrow\uparrow\downarrow\uparrow\downarrow$

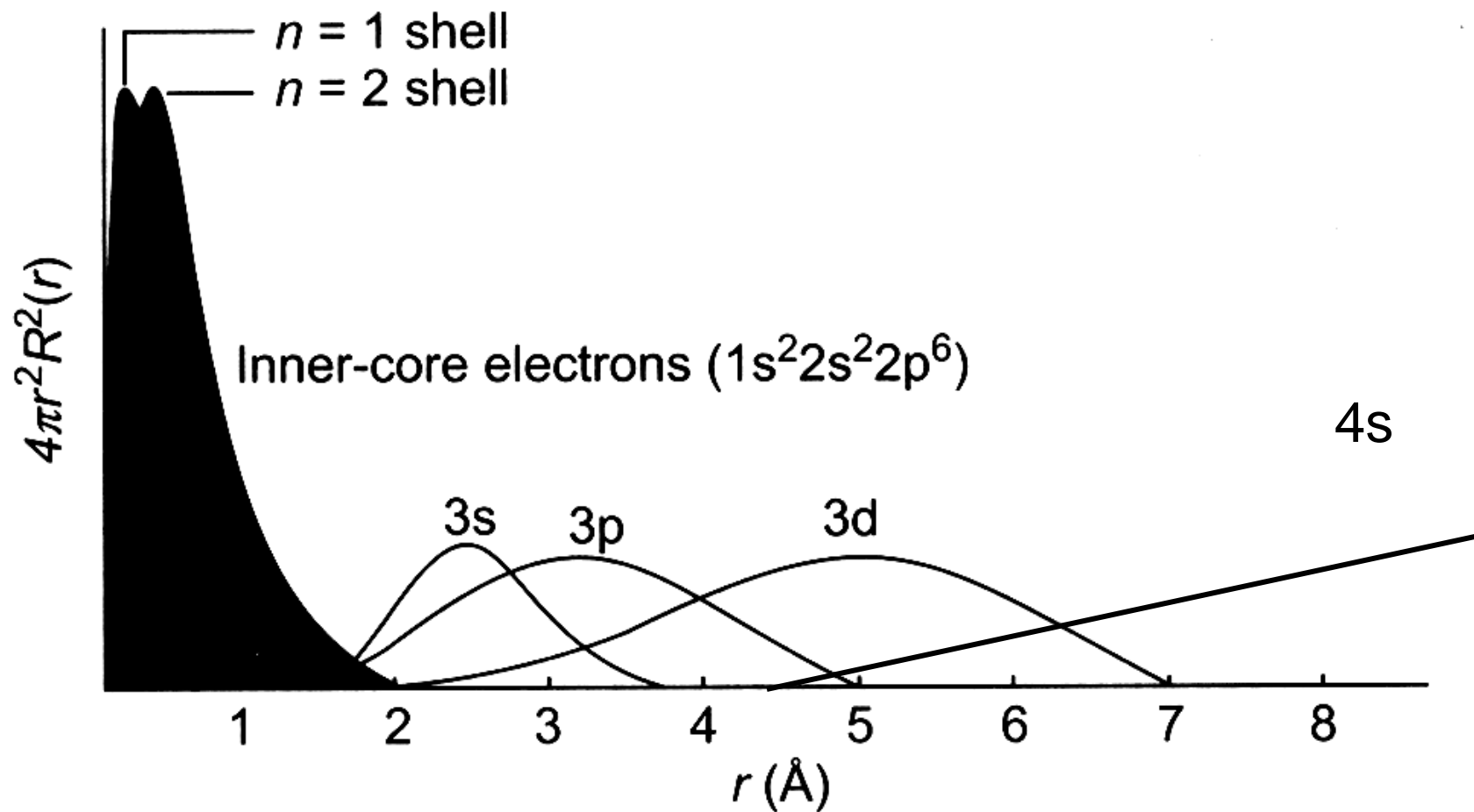


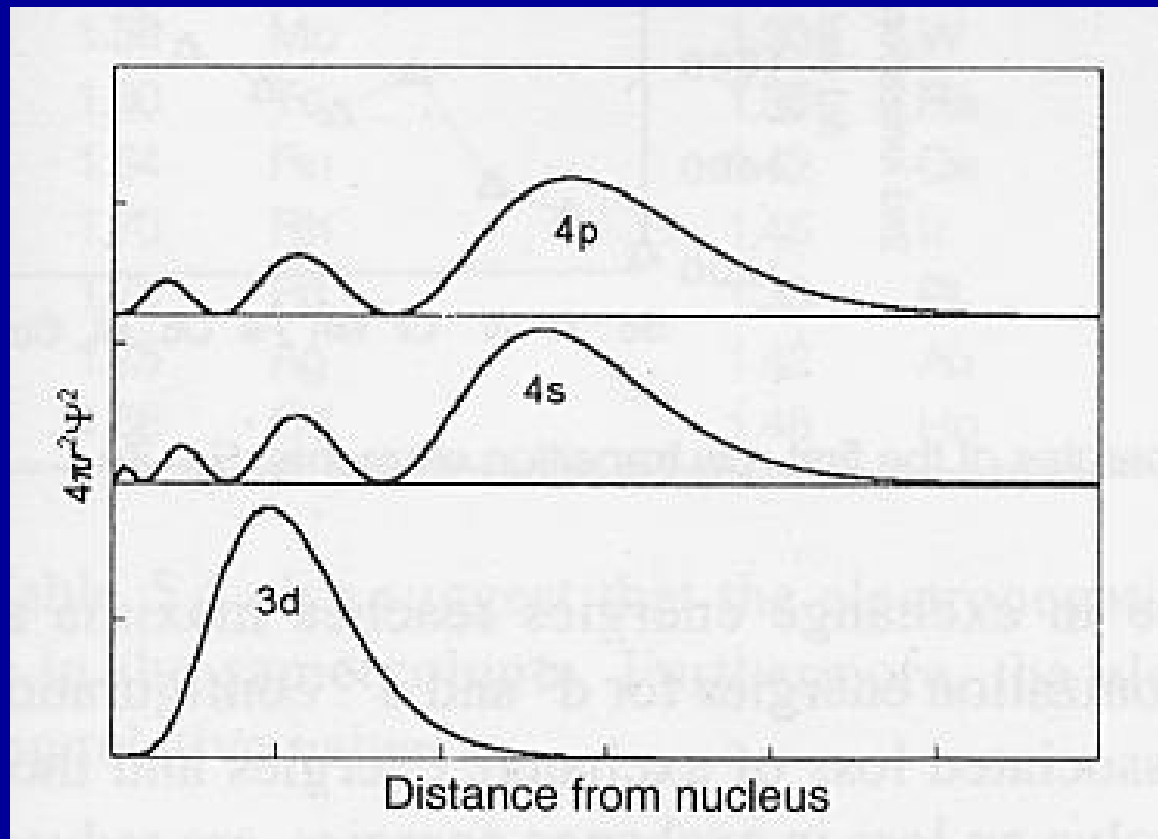
Orbital energy



Placing electrons in orbitals can change order of energy levels

Starting at Sc, $3d$ orbitals have lower energy than $4s$





Electron Configurations of Valence Shell

(Ar)

	3d	4s
Sc		
Ti		
V		
Cr		
Mn		
Fe		
Co		
Ni		
Cu		
Zn		



Ionisation Energies

