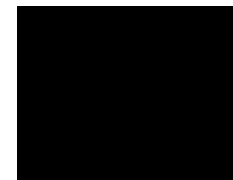
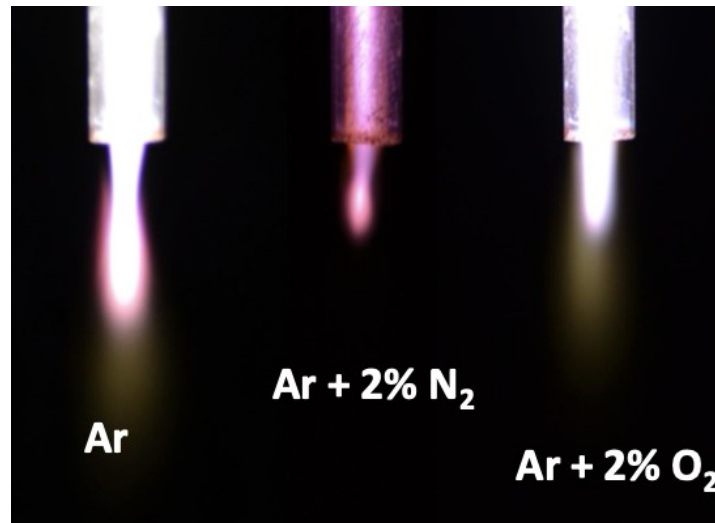


Plasma Physics 2

Lesson 05: High frequency discharges



Lecture series contents

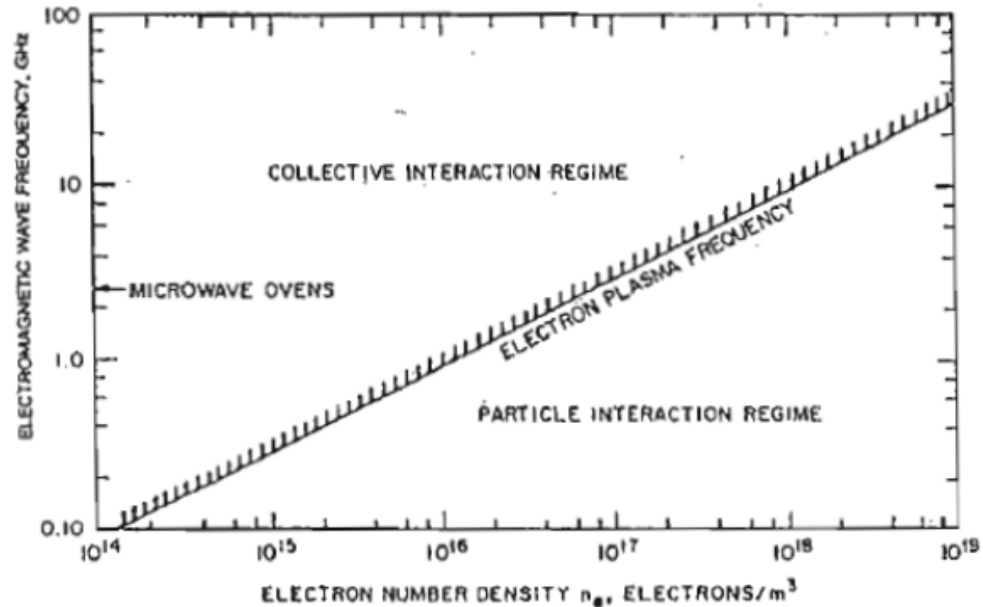
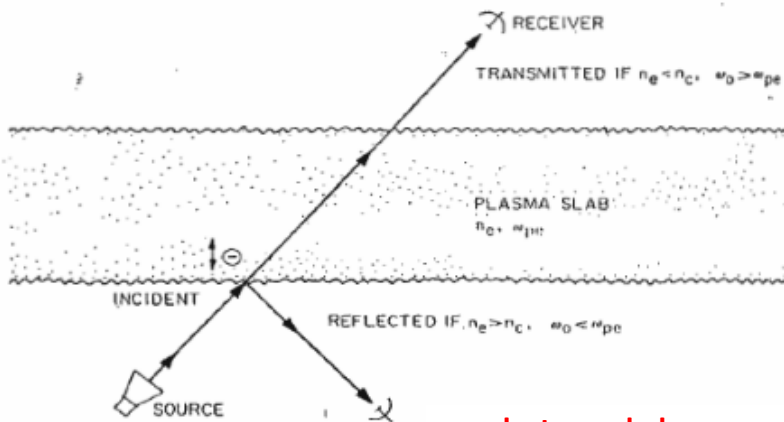
1. Townsend breakdown theory, Paschen's law
2. Glow discharge
3. Electric arc at low and high pressures
4. Magnetized low-pressure plasmas and their role in material deposition methods.
5. **Brief introduction to high-frequency discharges**
6. Streamer breakdown theory, corona discharge, spark discharge
7. Barrier discharges
8. Leader discharge mechanism, ionization and discharges in planetary atmospheres
9. Discharges in liquids, complex and quantum plasmas
10. Thermonuclear fusion, Lawson criterion, magnetic confinement systems, plasma heating and inertial confinement fusion.

High frequency discharges

- enable plasma discharge even for non-conductive electrodes
- electron plasma frequency

$$\omega_{pe} = \left(\frac{ne^2}{m\epsilon_0} \right)^{1/2} \quad (\text{rad/s})$$

- numerically $n_e \approx 10^{10} \text{ cm}^{-3}$ je $f = 0,9 \text{ GHz}$



electrons do have enough time to react to the external electric field!

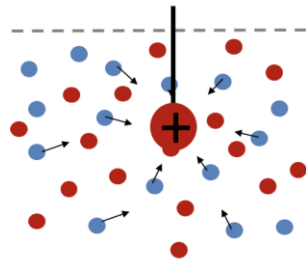
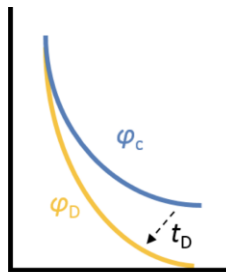
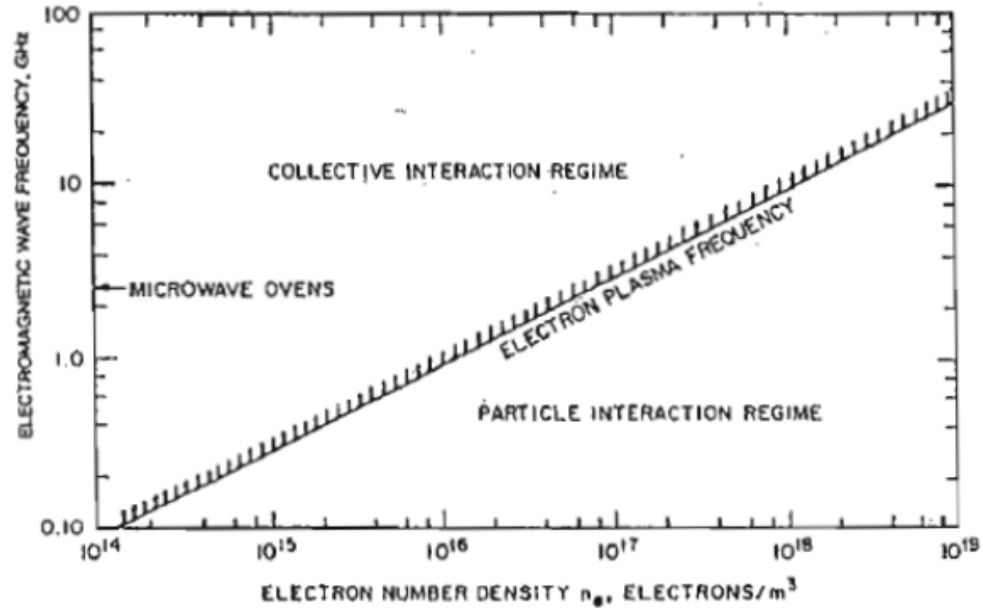
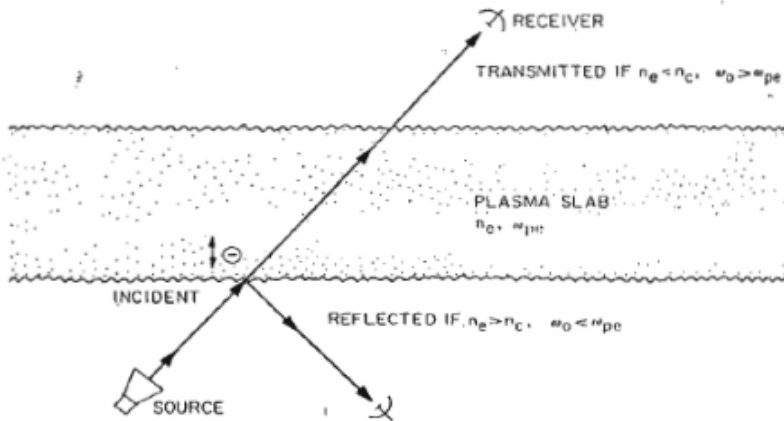
- varieties:
 - MHz = RF plasma, ohřev plazmy posuvným proudem
 - capacitively coupled plasma – collisionall and collisionless heating, can run at very low pressures
 - inductively coupled plasma – heating on the plasma surface or in plasma volume
 - GHz = microwave discharges, surface wave heating or volume heating

High frequency discharges

- enable plasma discharge even for non-conductive electrodes
- electron plasma frequency

$$\omega_{pe} = \left(\frac{ne^2}{m\epsilon_0} \right)^{1/2} \quad (\text{rad/s})$$

- numerically $n_e \approx 10^{10} \text{ cm}^{-3}$ je $f=0,9 \text{ GHz}$



$$t_D \simeq \frac{\lambda_D}{v_{te}} = \left(\frac{\epsilon_0 k_B T_e}{e^2 n_e} \cdot \frac{m}{k_B T_e} \right)^{1/2} = \omega_p^{-1}$$

$$L \gg \lambda_D$$

$$N_D := \frac{4\pi}{3} n \lambda_D^3 \gg 1$$

$$n_e = \sum n_i$$

$$\nu_{en} < \nu_{pl} \quad (kde \nu_{pl} = \omega_{pl}/2\pi)$$

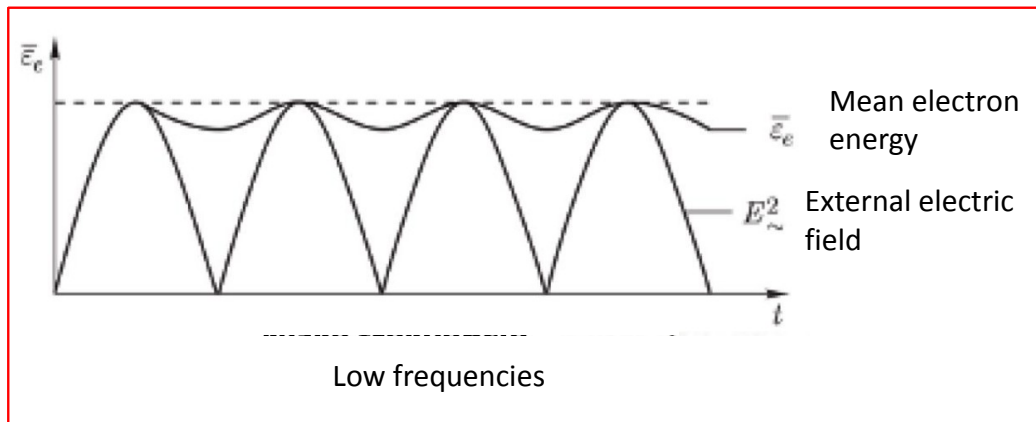
Fig. 4: Response time to form a Debye sheath

HIGH FREQUENCY BREAKDOWN

High frequency breakdown

- At about 100 kHz, the time that the electron needs to transfer between the electrodes starts to be comparable with the period of the external electric field. **The electron avalanche does not reach the electrode and it reverses direction!**

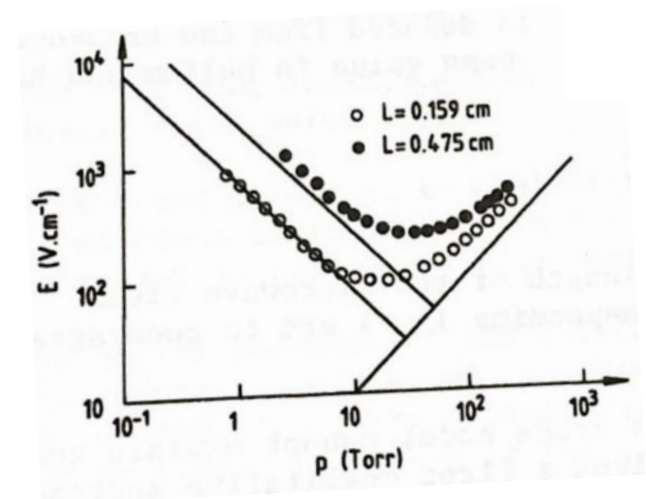
- electron heating in the external field is given as:
$$P = jE = \frac{e^2 E^2}{m} \cdot \frac{n_e v_{en}}{v_{en}^2 + \omega^2}$$



- The optimum power is reached when:

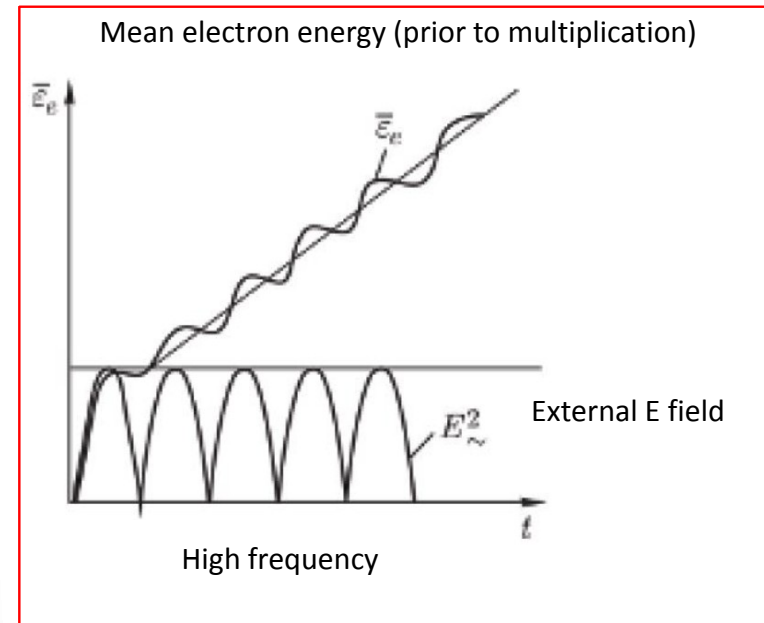
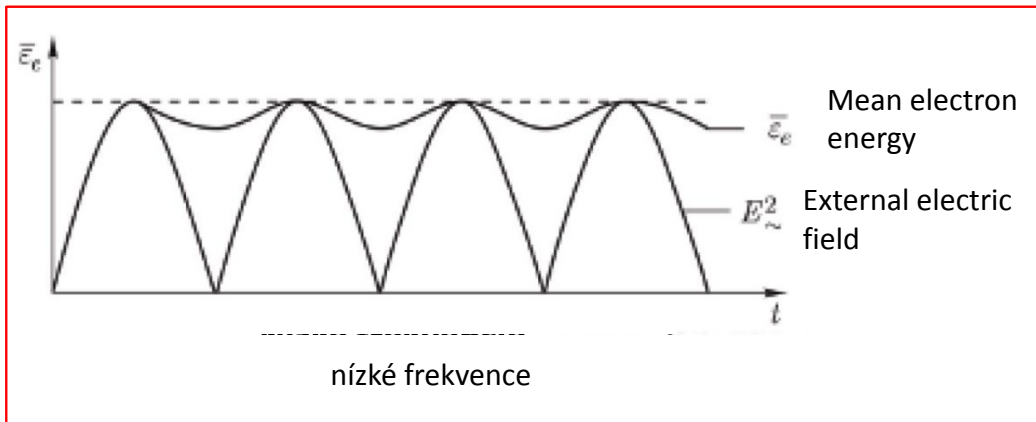
$$\left(\frac{dP}{dv_{en}} = 0 \right)$$

$$\omega = v_{en}$$



High frequency breakdown

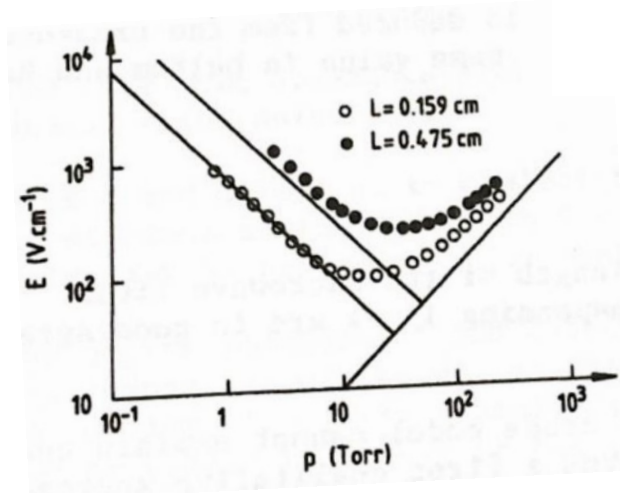
- At about 100 kHz, the time that the electron needs to transfer between the electrodes starts to be comparable with the period of the external electric field. **The electron avalanche does not reach the electrode and it reverses direction!**
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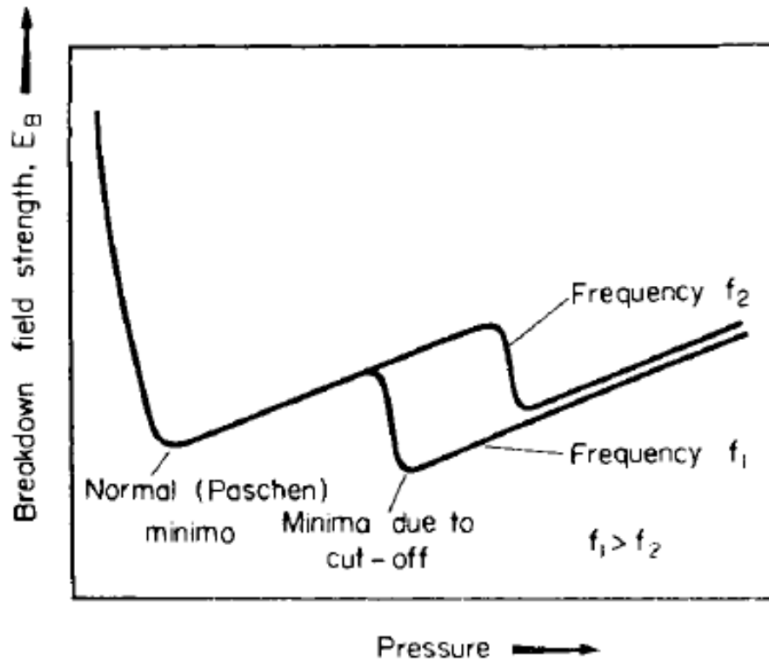
- The optimum power

$$\left(\frac{dP}{dv_{en}} = 0 \right)$$

$$\omega = v_{en}$$

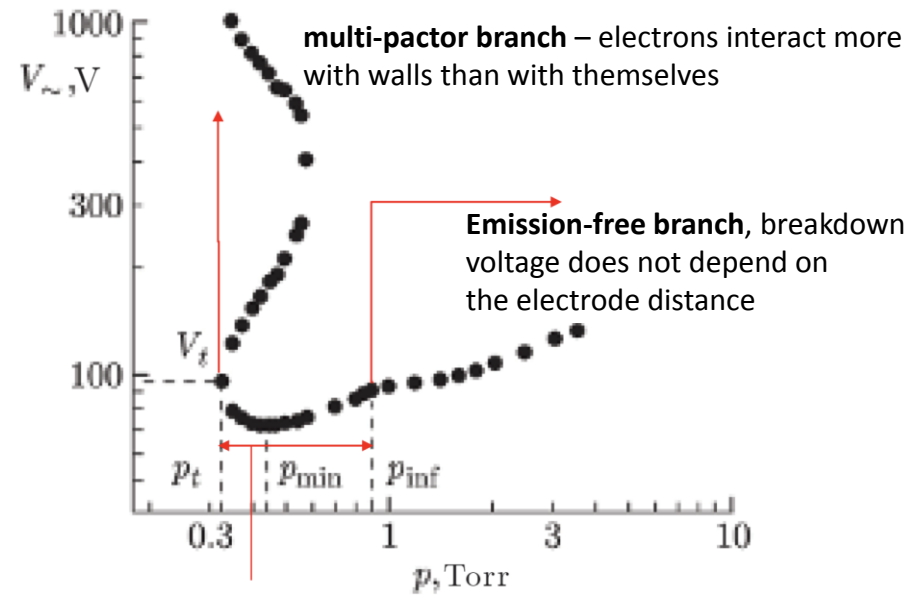


High Frequency breakdown



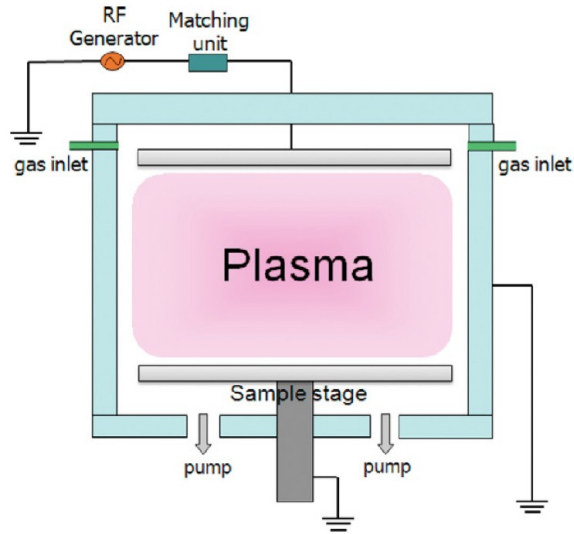
- The Paschen curve for RF breakdown has two minima:
 - 1) **Classical DC-like:** Electron attains some „ideal“ energy when traversing between electrodes
 - 2) **Resonance capture of electrons:** electrons bounce off a plasma sheath that is approaching them.

J. Phys. D: Appl. Phys. **31** (1998) 3349–3357.

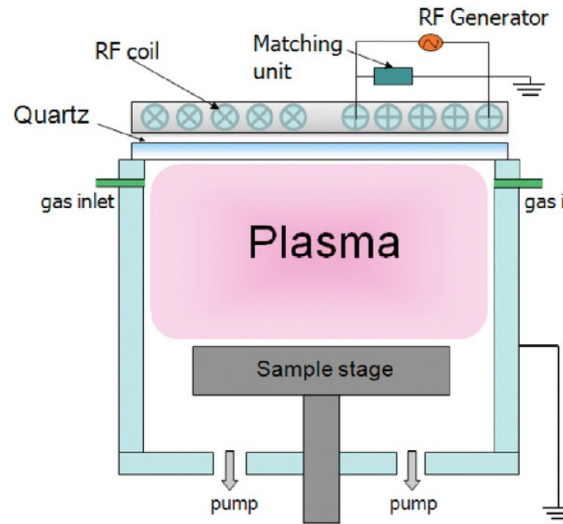


Drift-diffusion branch – secondary emission affects the breakdown voltage

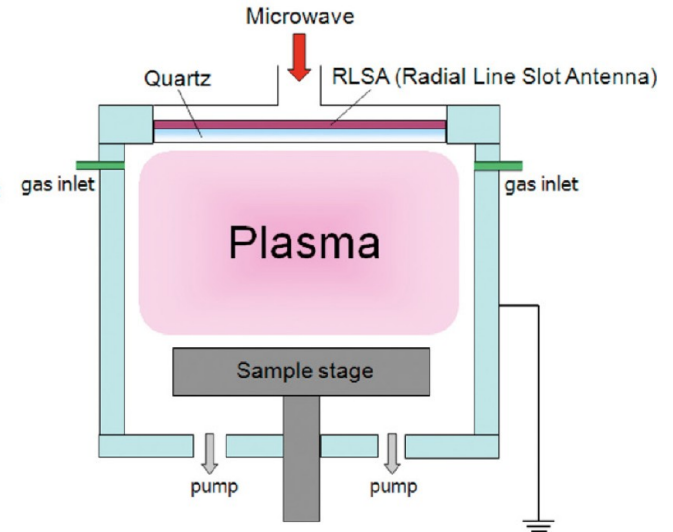
CCP vs ICP vs MW plasma



(a) Capacitively coupled plasma



(b) Inductively coupled plasma



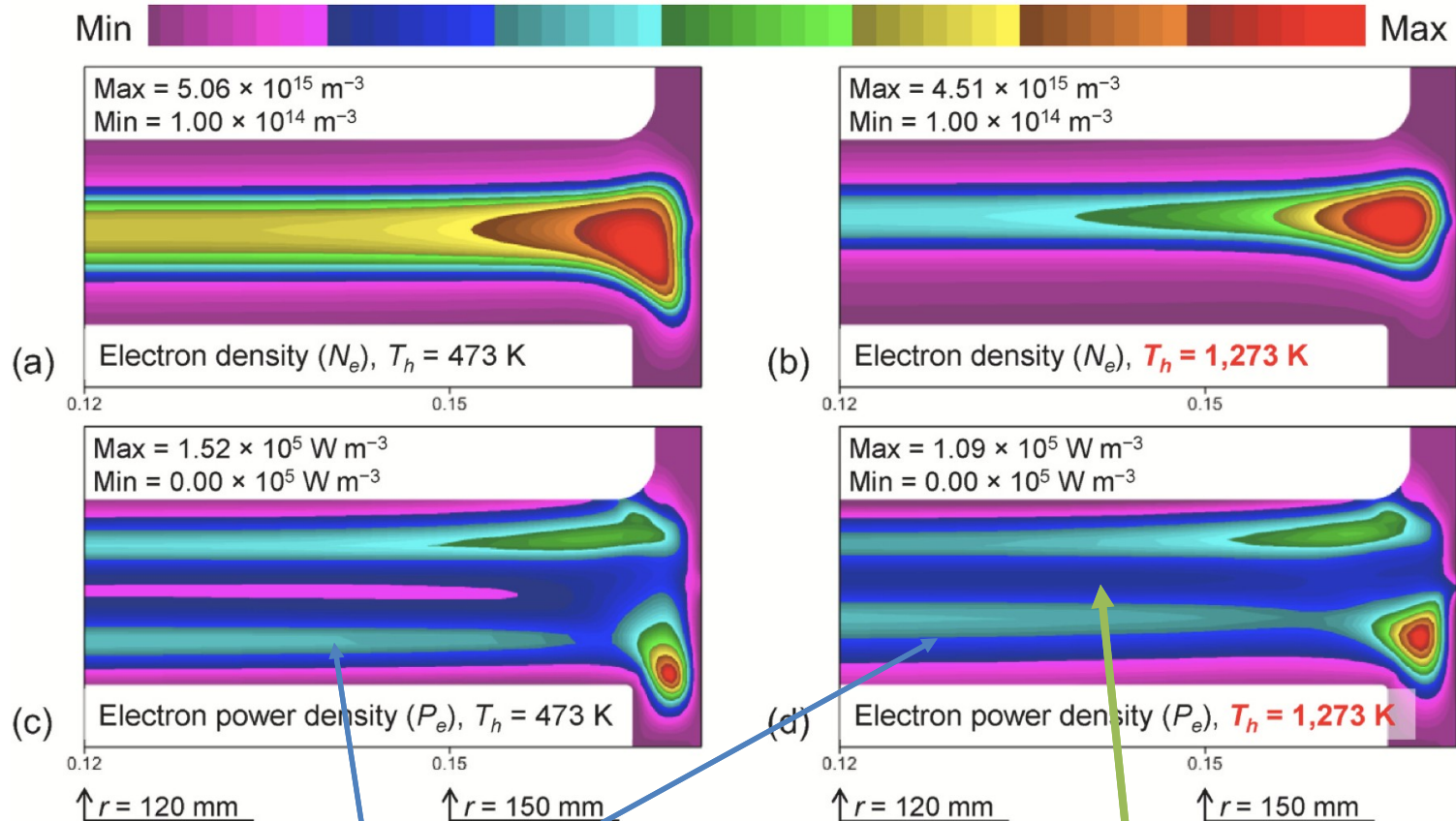
(c) Microwave plasma

CAPACITIVELY COUPLED PLASMA ANATOMY

CCP plasma anatomy

Typical frequencies: 10 MHz – 100 MHz (13.56 MHz "standard")

Typical pressure: 1 Pa – 1000 Pa



Near-electrode electron heating
(stochastic)

Bulk plasma heating (ohmic)

CCP plasma anatomy

- Simplest case are two planar electrodes, covered with dielectric. Potential and density evolution during phase as follows:

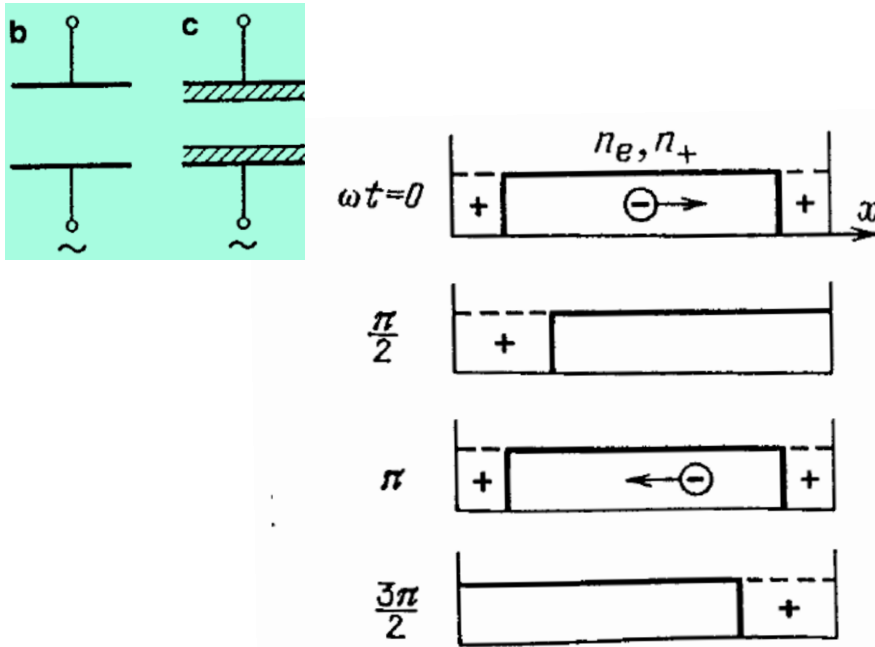


Fig. 13.1. Swings of the electron gas. The gas of ions is assumed to be fixed and uniform. Time is measured from the moment when electrons pass through the equilibrium position when moving to the right. Distributions $n_e(x)$ are shown for every quarter of the period

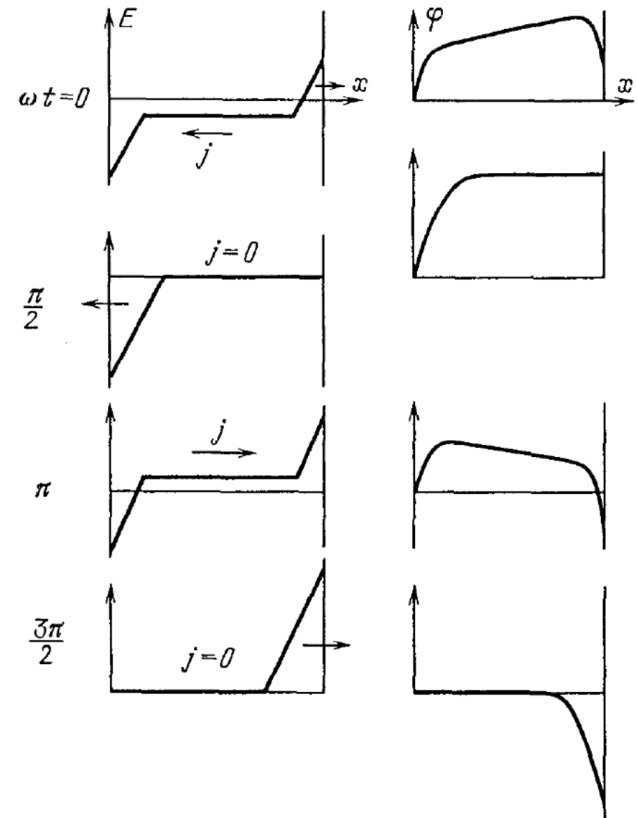


Fig. 13.2. Field and potential distributions in the gap corresponding to Fig. 13.1

CCP plasma – self bias voltage

- In CCP plasmas, we see emergence of the **so-called self-bias DC voltage**.
- empirically, we see: $\frac{V_1}{V_2} = \left(\frac{A_2}{A_1}\right)^q$ $1 \leq q \leq 2.5$ where A_1, A_2 are electrode surface areas

- An approximate derivation can be made as follows – we assume that the ion current on both electrodes has to be the same:

$$I_{i1} = I_{i2}$$

$$A_1 n_i \bar{v}_{i1} = A_2 n_i \bar{v}_{i2}$$

$$\bar{v}_i = \sqrt{\frac{2eV}{m_i}}$$

- And if we further assume that the ion densities in front of both electrodes are the same $n_{i1} = n_{i2}$ and $m_{i1} = m_{i2}$, we obtain

$$A_1 n_i \sqrt{\frac{2eV_1}{m_i}} = A_2 n_i \sqrt{\frac{2eV_2}{m_i}}$$

- And ultimately

$$\frac{A_1}{A_2} = \sqrt{\frac{V_2}{V_1}}$$

The smaller electrode has to have **higher negative voltage** to draw the same ion current => it acts as a cathode and it can be used e.g. for sputtering of material

This derivation is still extremely simplistic, will be done more properly in Plasma Physics 3

INDUCTIVELY COUPLED PLASMA ANATOMY

ICP plasma anatomy

Typical frequencies: 10 MHz – 100 MHz (13.56 MHz "standard")

Typical pressure: 1 Pa – 1000 Pa

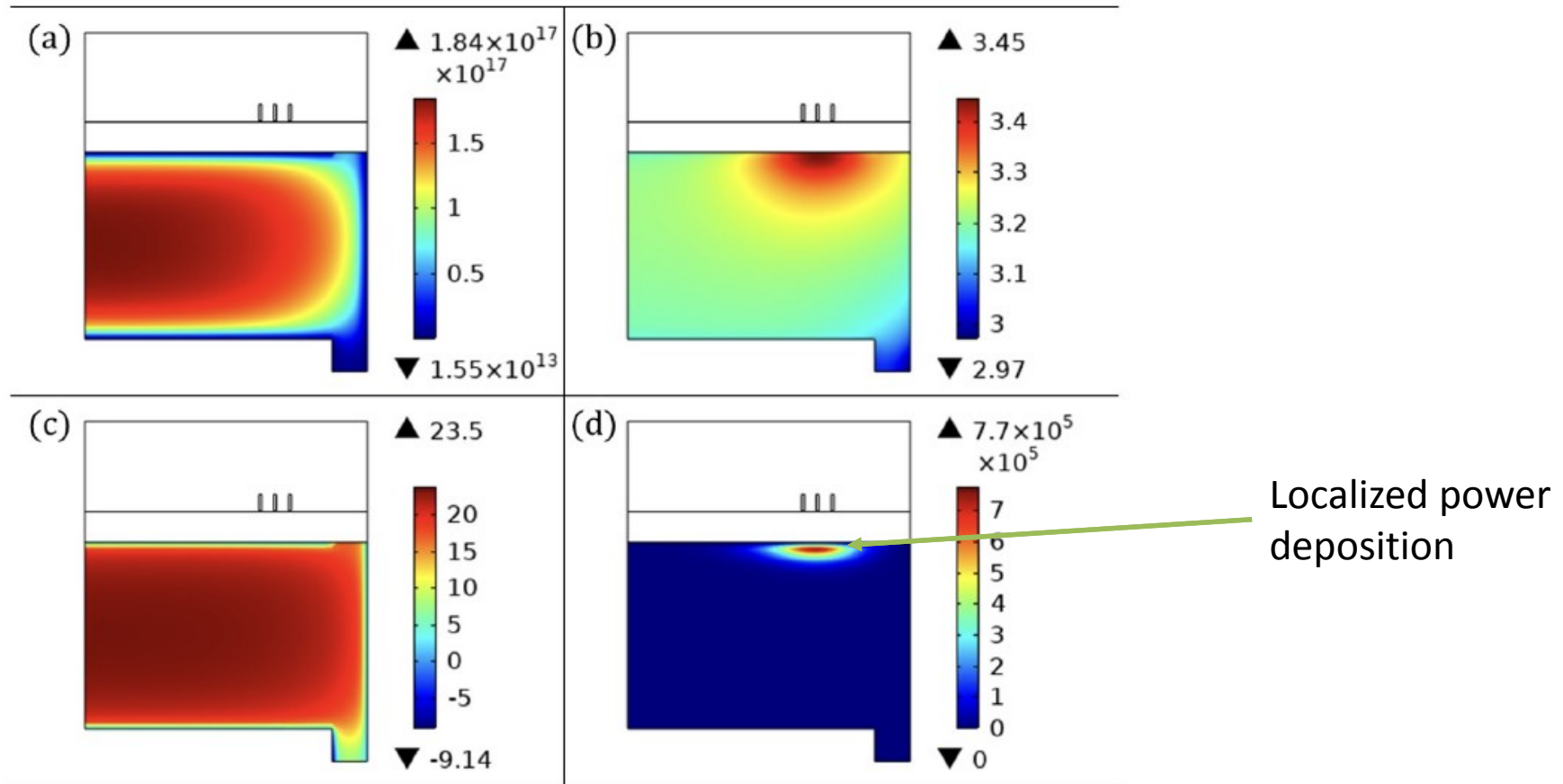


Figure 4. Plasma properties computed for Druyvesteyn EEDF at 10 mTorr and gas flow rate of 100 sccm, and absorbed power of 570 W; (a) electron density ($1/m^3$), (b) electron temperature (eV), (c) plasma potential (V), and (d) power density (W/m^3).

ICP plasma anatomy

- Plasma is generated by **electromagnetic induction** – The RF coil creates H field, which creates poloidal E field in the **volume**
- This wireless transfer of energy makes it **possible to run the system at very high temperatures**, because the walls can be made of ceramics with high thermal resistivity.

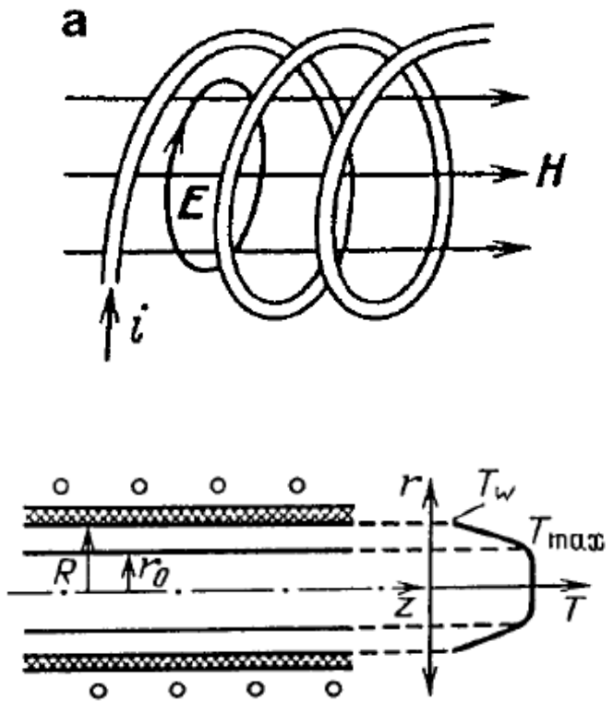


Fig. 11.1. Induction discharge in a tube of radius R placed inside a long solenoid; r_0 is the discharge radius. The radial temperature distribution is given on the *right*

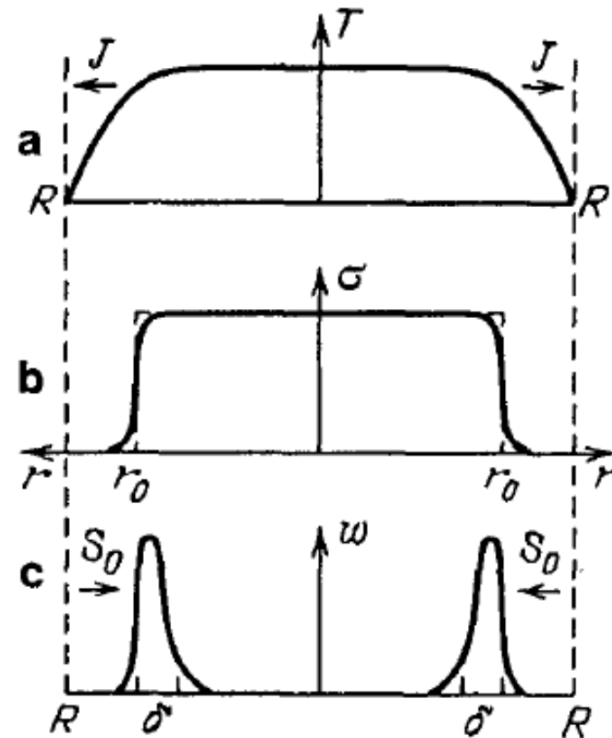


Fig. 11.2. Radial distributions of (a) temperature, (b) conductivity, and (c) Joule heat release in induction discharge. Dashed curves correspond to step function $\sigma(r)$ in the metallic cylinder model. Arrows: J , heat flux; S_0 , electromagnetic energy flux; δ , skin layer thickness

MICROWAVE PLASMA ANATOMY

MW plasma anatomy

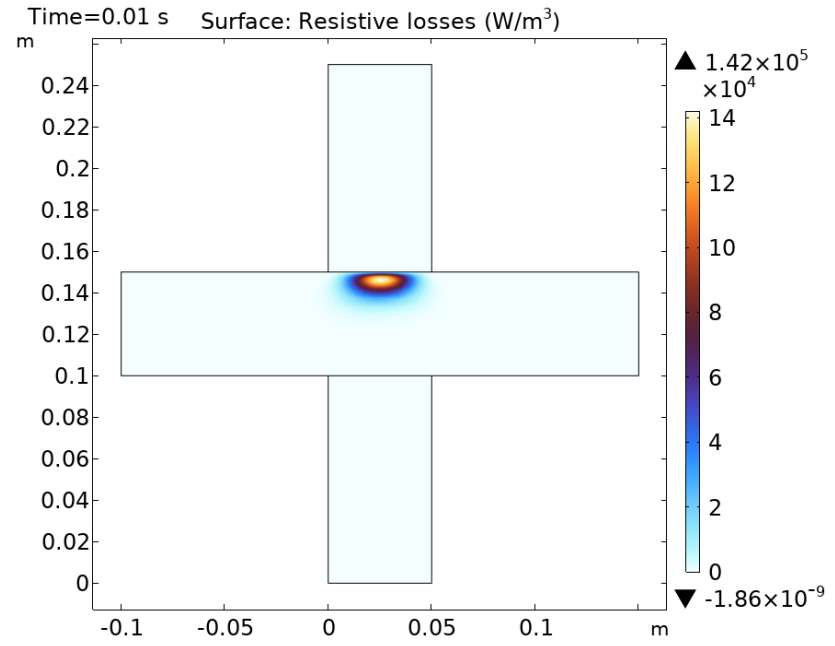
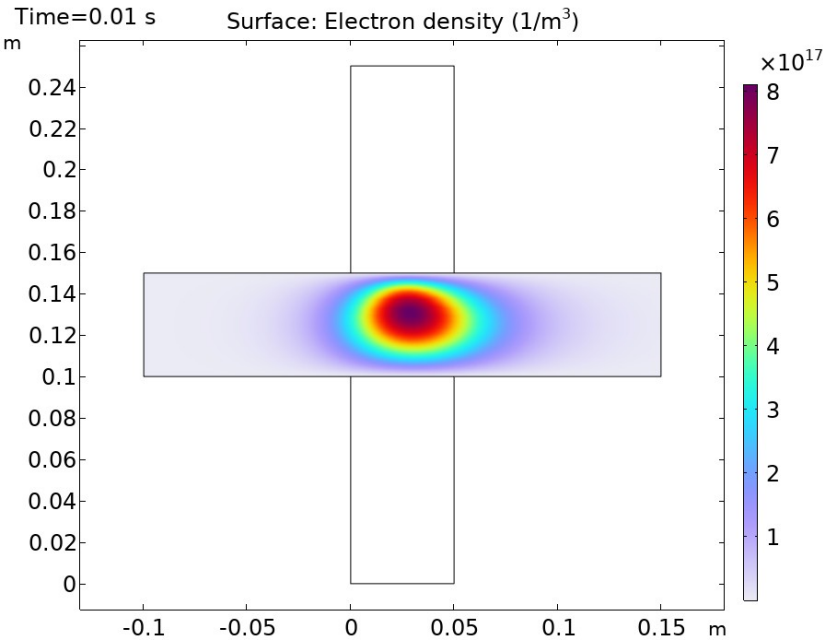
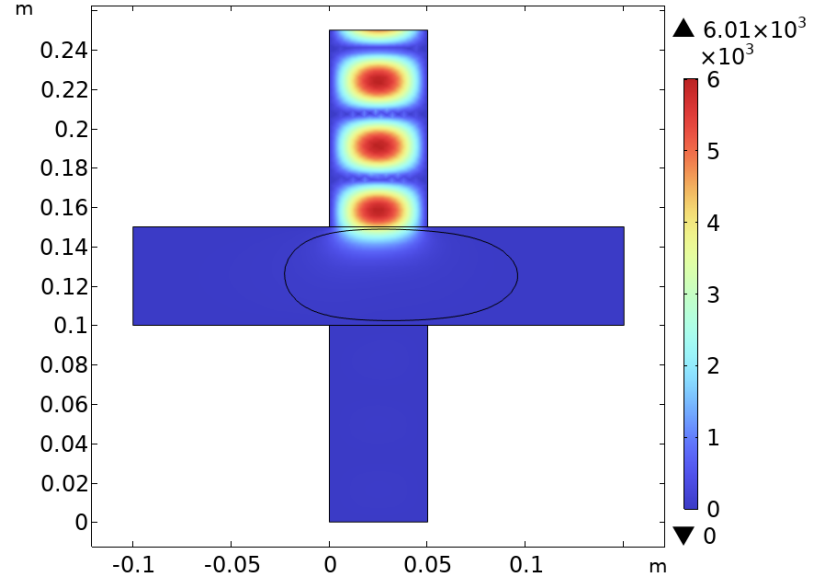
Typical frequencies:

100 MHz – 10 GHz MHz (2.45 GHz "standard")

Gas pressure:

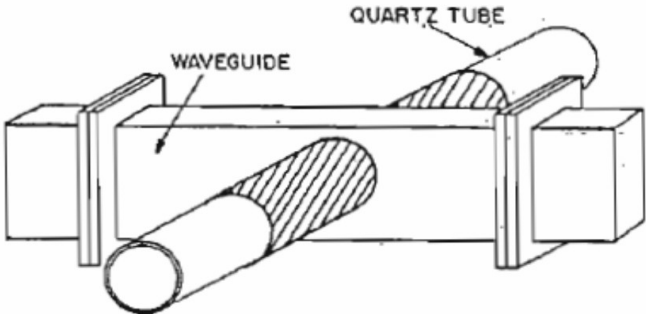
100 Pa – 1 atm

Time=0.01 s Surface: Electric field norm (V/m) Contour: Electron density (1/m³)

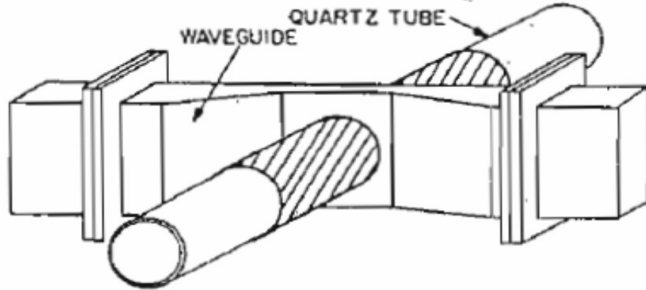


MW plasma anatomy

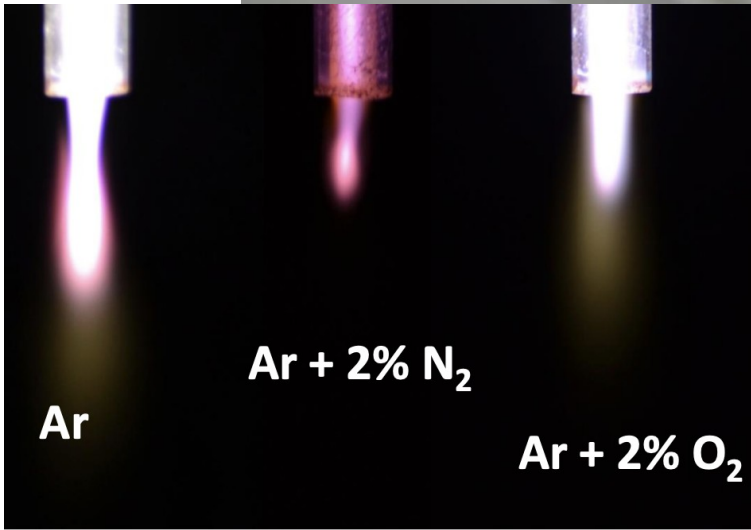
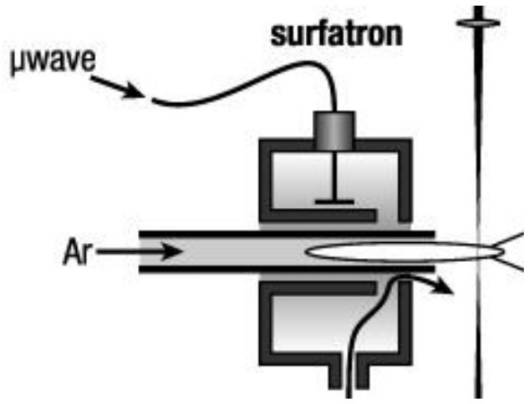
- There are various setups for MW plasma generation – surfatron, surfaguide, coaxial line,



a) STANDARD PLASMA APPLICATOR

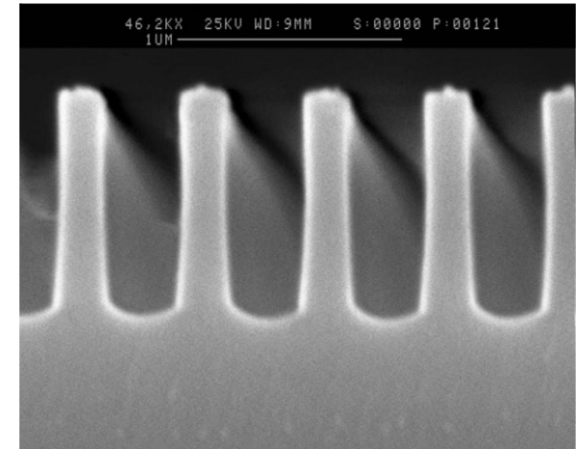
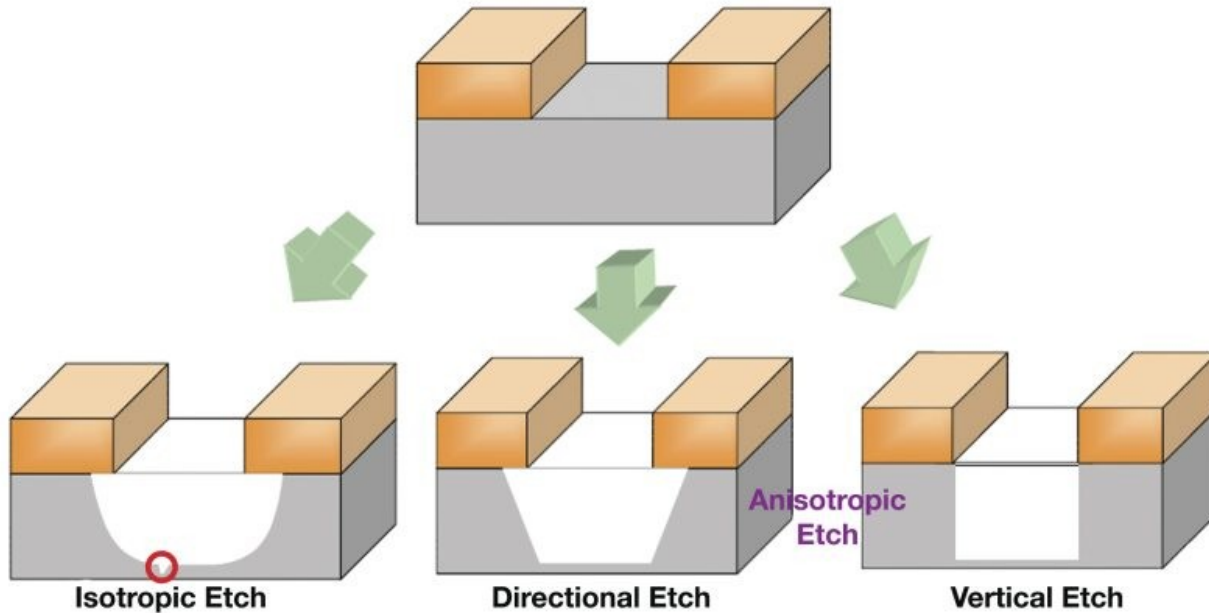
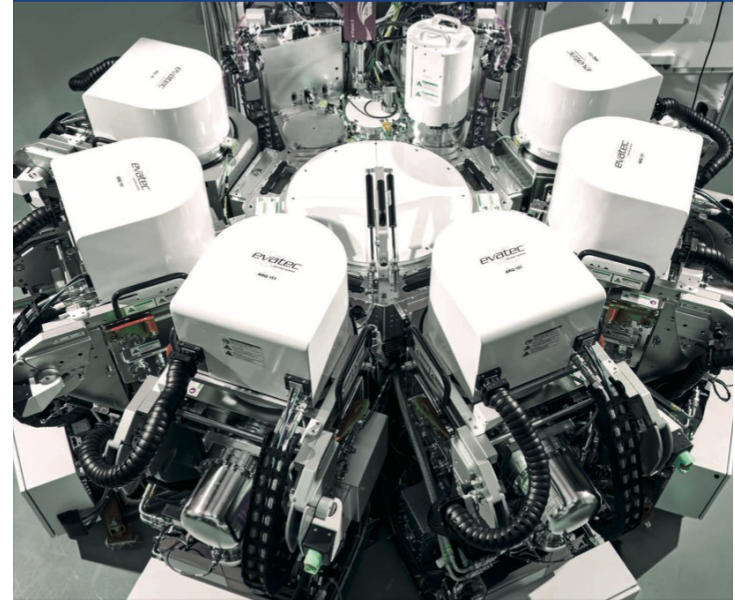
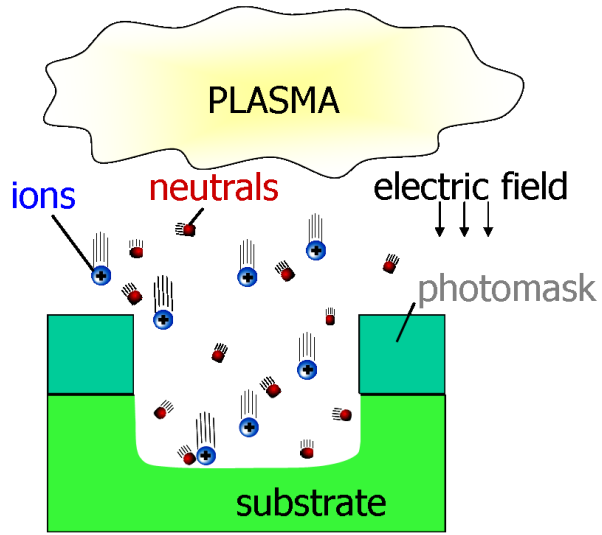


b) HIGH ELECTRIC FIELD APPLICATOR

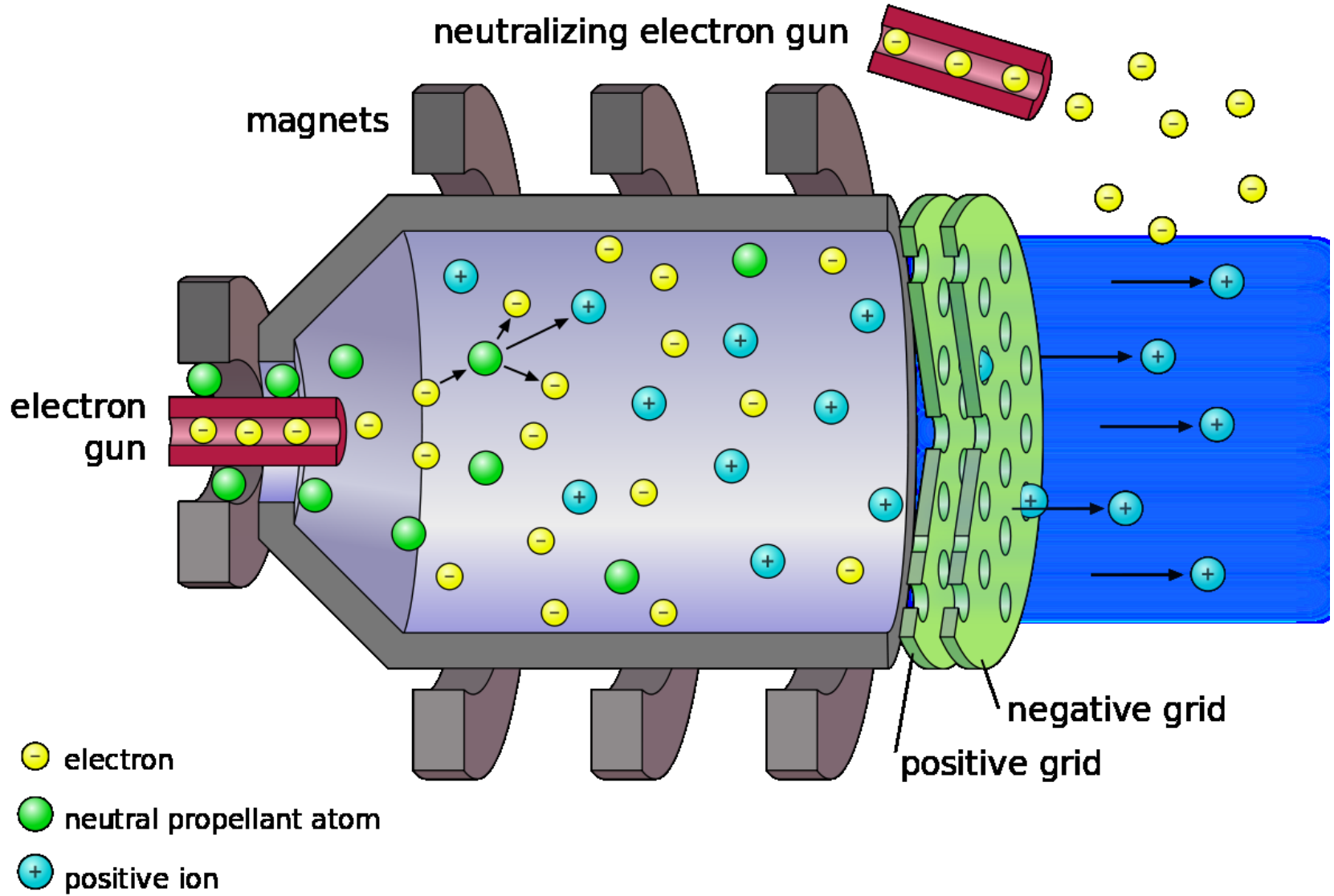


APPLICATIONS OF CCP, ICP, MW

CCP and ICP applications – semiconductor etching

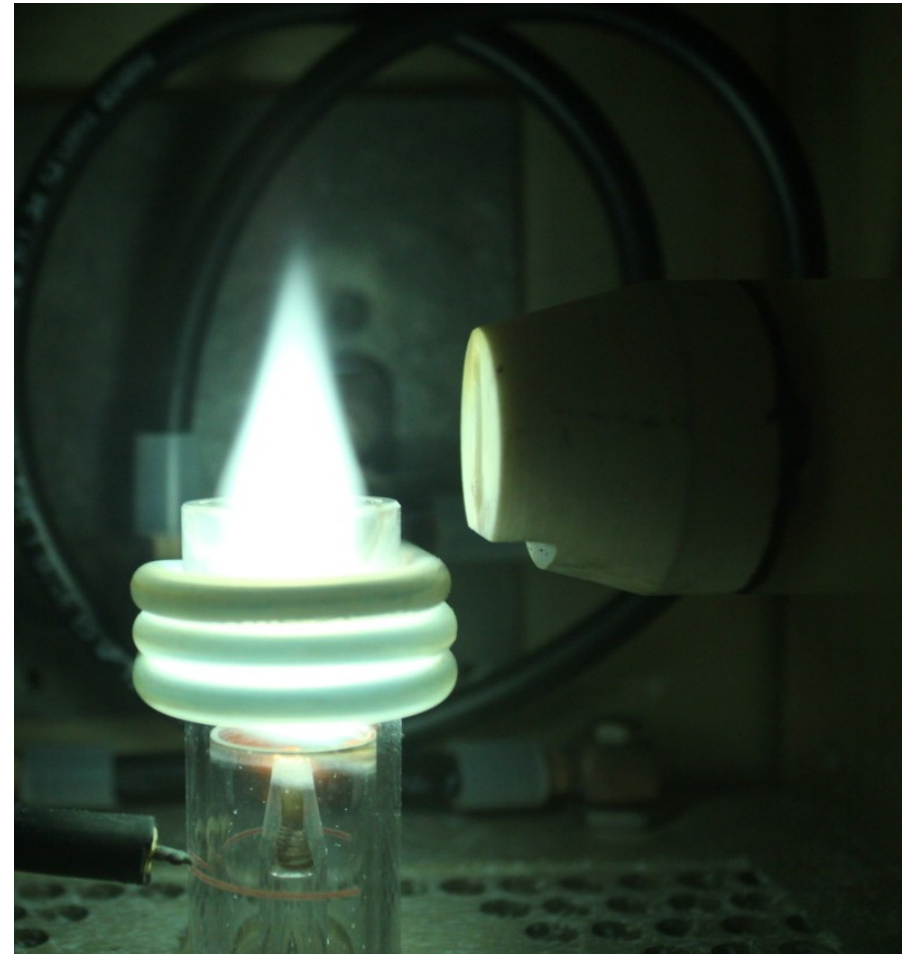
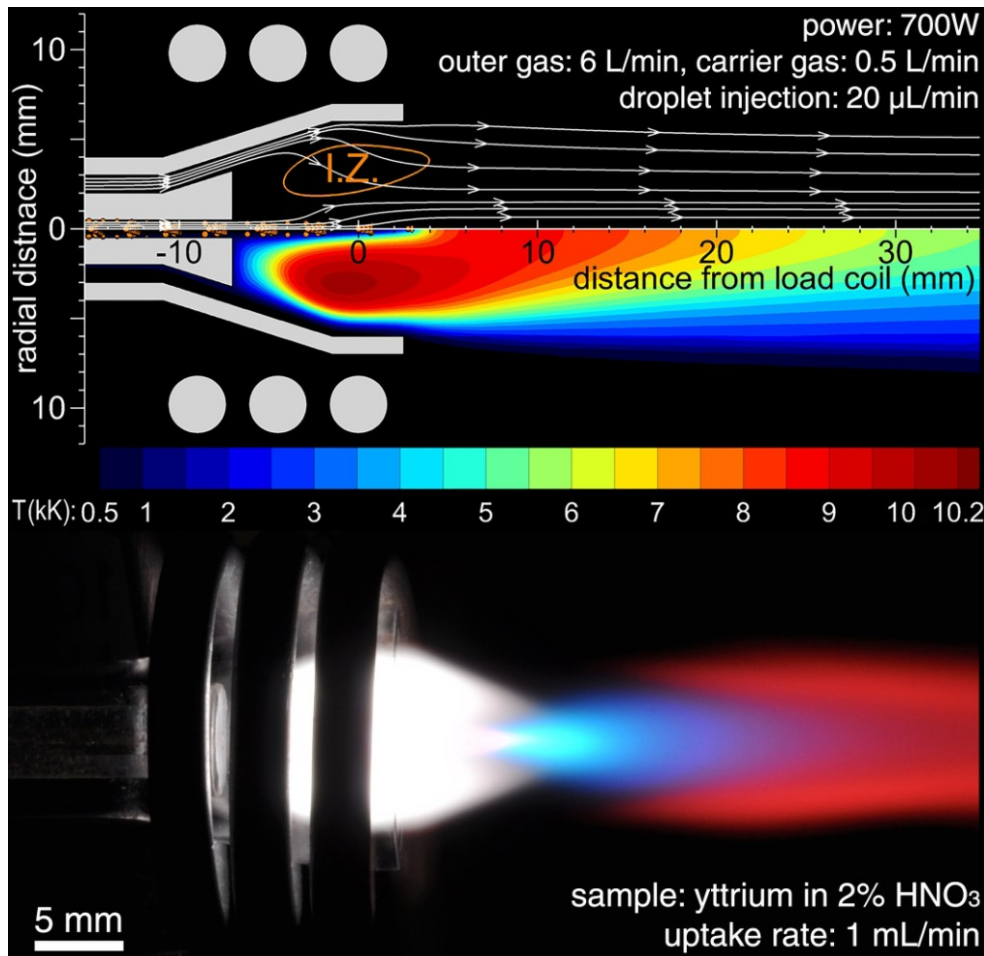


CCP application – gridded ion thruster



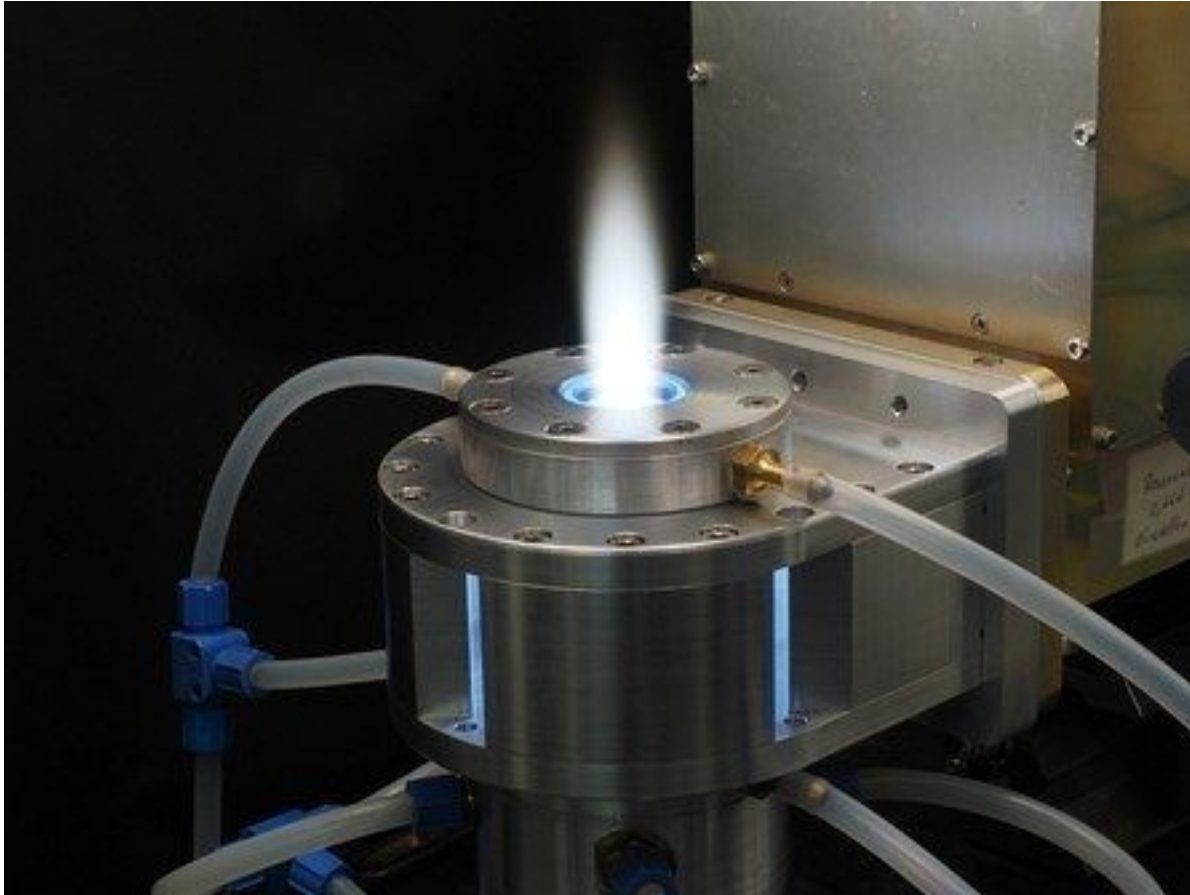
ICP application – DC torch analogue

- At atm. pressure and high power levels, so-called ICP torches can generate fully ionized plasmas. Support similar power ranges (5 kW – 1 MW)
- Advantage over DC torches: no contact with metal



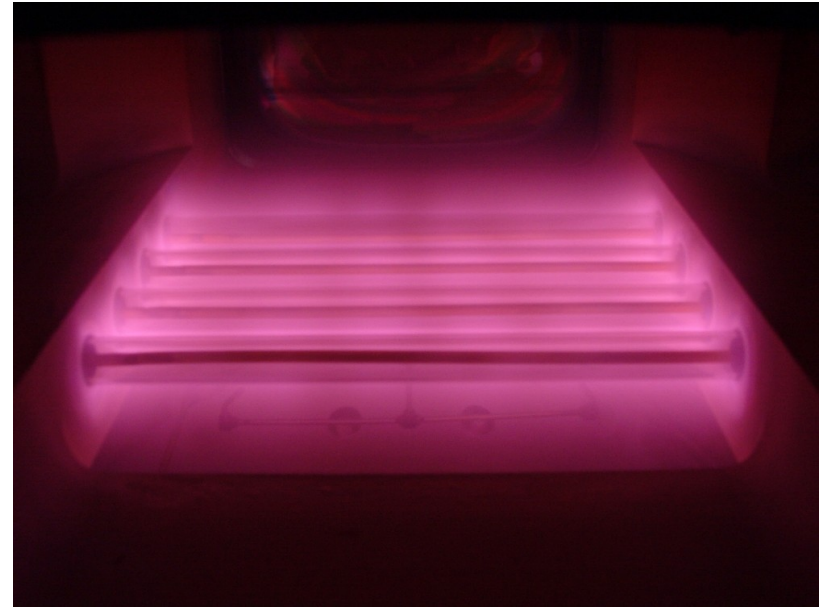
MW application – DC torch analogue

- MW plasma is also being used as an analogue to DC torches
- As opposed to DC or ICP torches, MW plasma typically operates at lower powers (tens of kW)
- Microwave plasma can also be non-LTE (plasma does not have to be so conductive to enable MW absorption) => lower plasma densities achievable



MW application – Plasmaline for solar cells

- At atm. pressure and high power levels, so-called ICP torches can generate fully ionized plasmas.
- Advantage over DC torches: no contact with metal



Takeaways

- High frequency breakdown
- Different layouts for wave-driven discharges
- CCP, ICP, MW – different power absorption mechanisms
- Applications