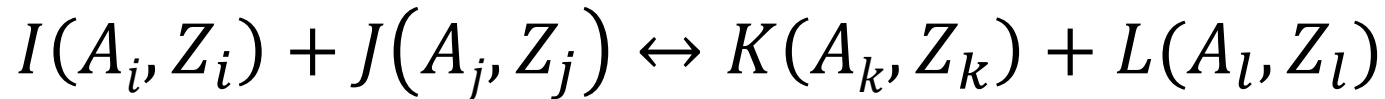


The binding energy of the atomic nucleus

The general description of a nuclear reaction is



Where A_i = *the baryon number, or nucleon number (nuclear mass)*

and Z_i = *the nuclear charge*

The nucleus of any element is uniquely defined by the two integers A_i and Z_i

Recall that in any nuclear reaction the following must be conserved:

1. The baryon number – protons, neutrons and their anti-particles
2. The lepton number – light particles, electrons, positrons, neutrinos, and anti-neutrinos
3. Charge

Note also that the anti-particles have the opposite baryon/lepton number to their particles

The binding energy of the atomic nucleus

- The total mass of a nucleus is known to be less than the mass of the constituent nucleons
- Hence there is a decrease in mass if a companion nucleus is formed from nucleons, and from the Einstein mass-energy relation $E = mc^2$ the mass deficit is released as energy
- This difference is known as the *binding energy* of the compound nucleus
- Thus if a nucleus is composed of Z protons and N neutrons, its binding energy is

$$Q(Z, N) \equiv [Zm_p + Nm_n - m(Z, N)]c^2$$

- For our purposes, a more significant quantity is the total *binding energy per nucleon*
- We can then consider this number relative to the hydrogen nucleus

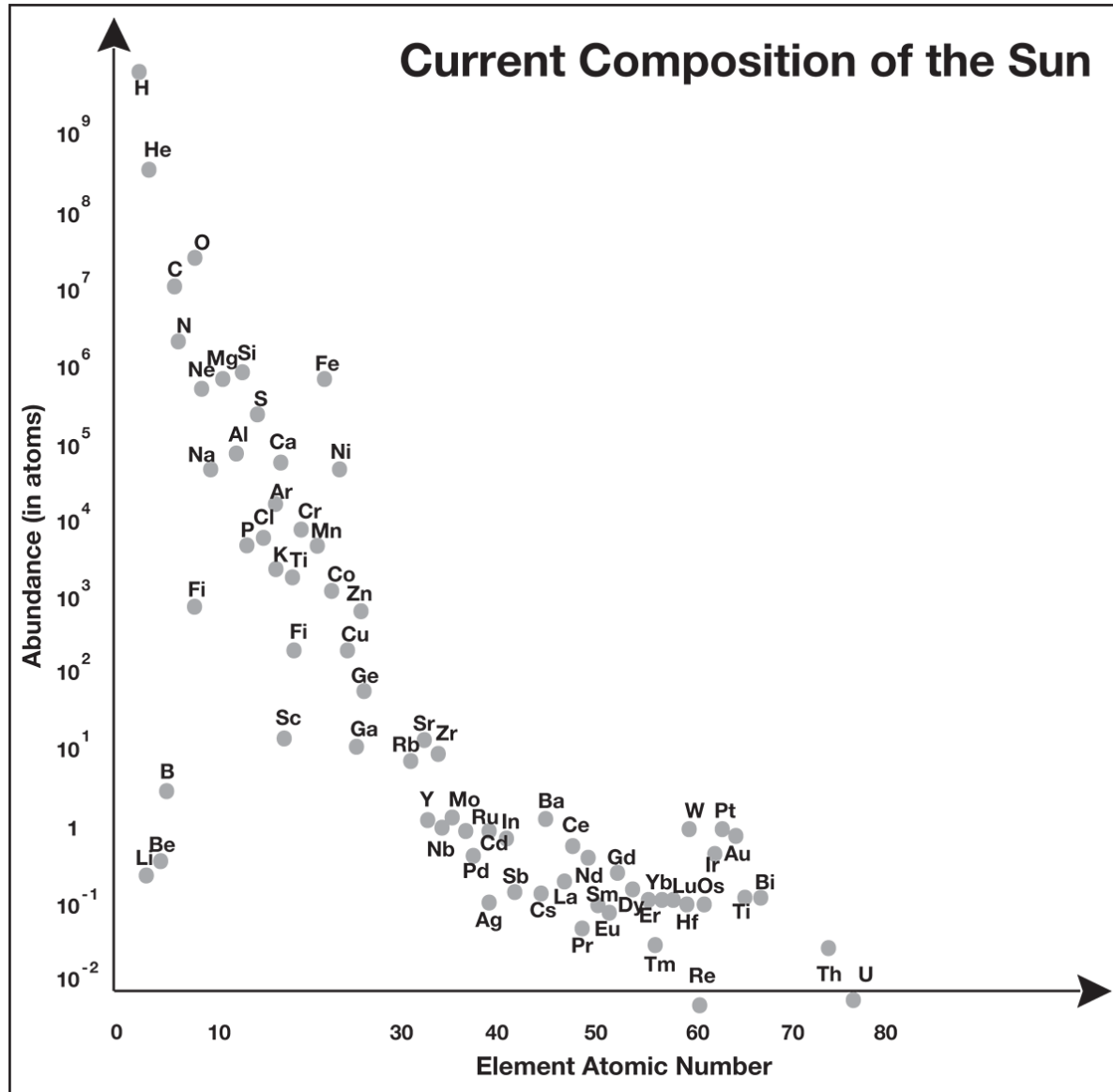
$$\frac{Q(Z, N)}{A}$$

The binding energy per nucleon

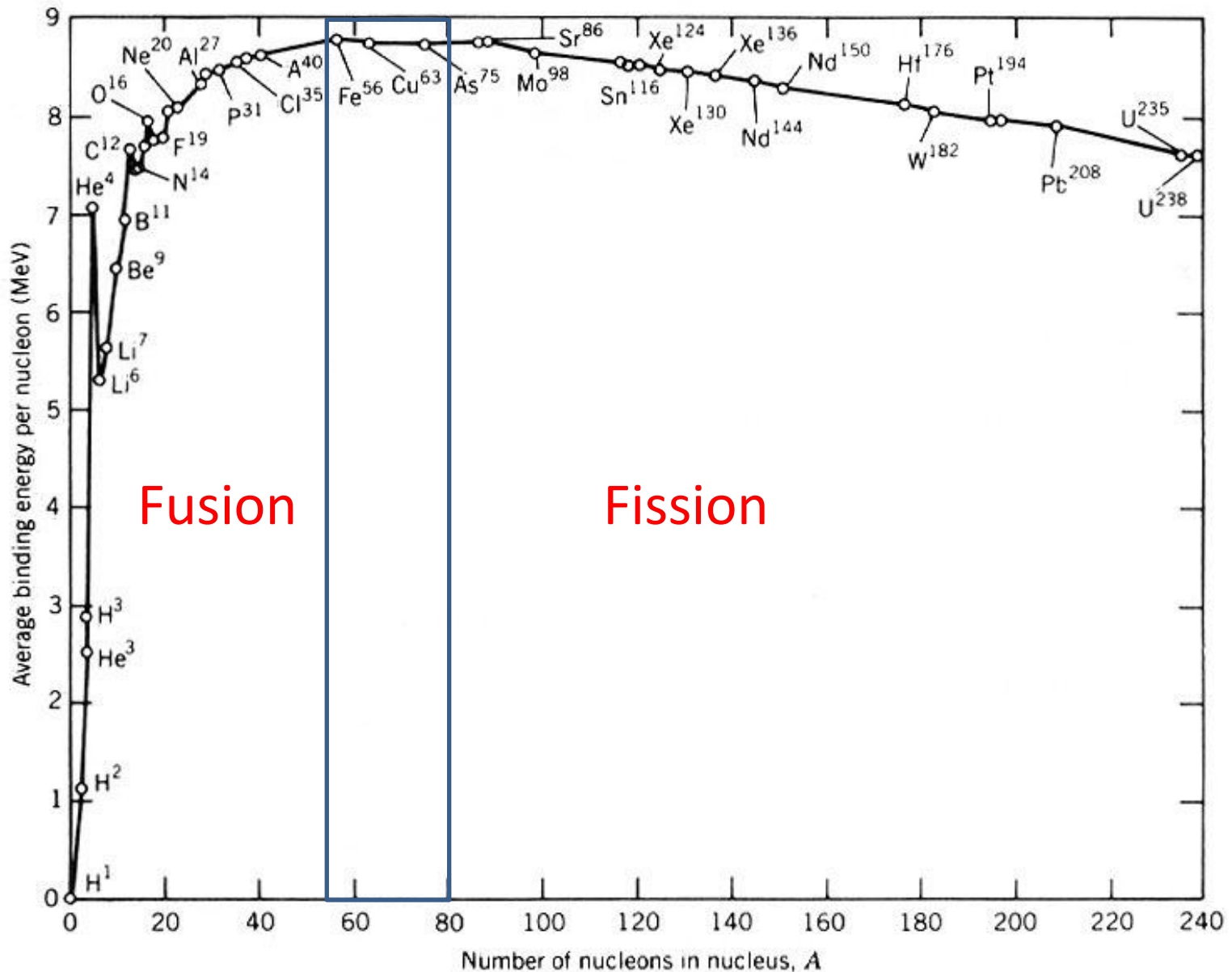
The variation of binding energy per nucleon with baryon number A

- General trend is an increase of Q with atomic mass up to $A= 56$ (Fe), then slow monotonic decline
- There is steep rise from H through ^2H , ^3He , to ^4He => fusion of H to He should release larger amount of energy *per unit mass* than fusion of He to C
- Energy may be gained by *fusion* of light elements to heavier, up to iron
- Or from *fission* of heavy nuclei into lighter ones down to iron

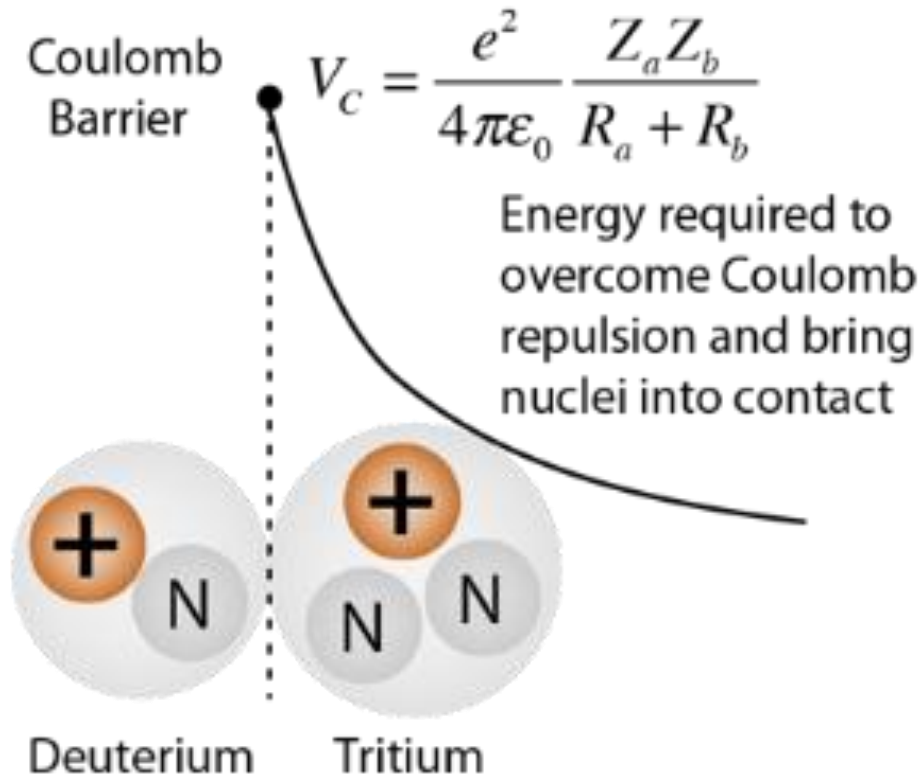
Abundance - Sun



The binding energy per nucleon



Coulomb Barrier



$Z_a Z_b$... number of protons in each nuclei

$R_a R_b$... interaction radii

ϵ_0 ... permittivity of free space ($8.85 \times 10^{-12} \text{ C}^2\text{N}^{-1} \text{ m}^{-2}$)

e ... charge of electron

D-T reaction: V_C is 0.38 MeV

Gas temperature of

$4.4 \times 10^9 \text{ K}$

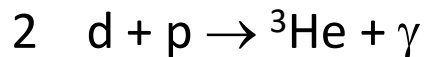
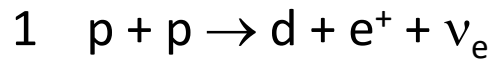
Hydrogen and helium burning

The most important series of fusion reactions are those converting H to He (H-burning). As we shall see this dominates ~90% of lifetime of nearly all stars.

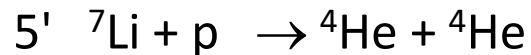
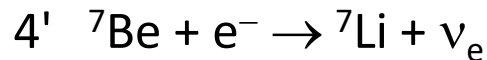
- Fusion of 4 protons to give one ${}^4\text{He}$ is completely negligible
- Reaction proceeds through steps – involving close encounter of 2 particles
- We will consider the main ones: the **PP - chain** and the **CNO cycle**

The PP - chain has three main branches called the PPI, PPII and PPIII chains.

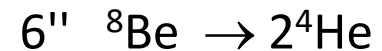
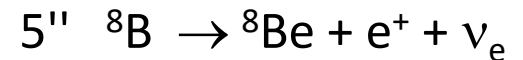
PPI Chain



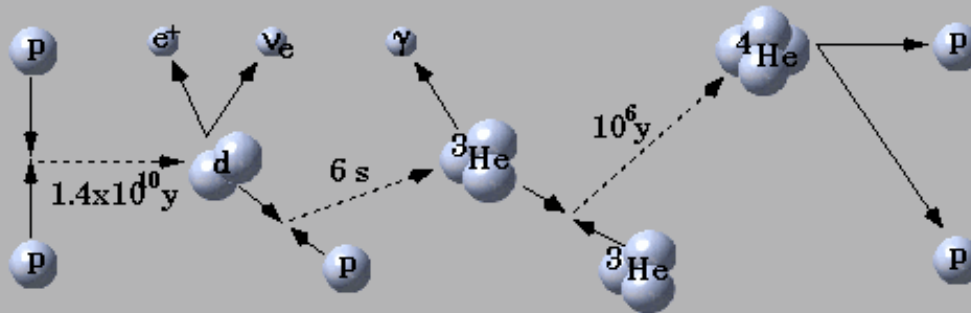
PPII Chain



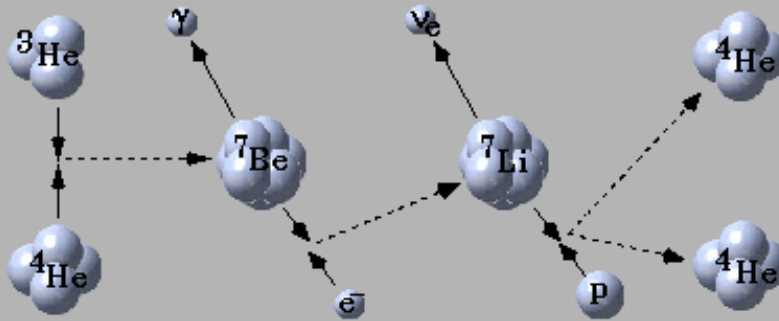
PPIII Chain



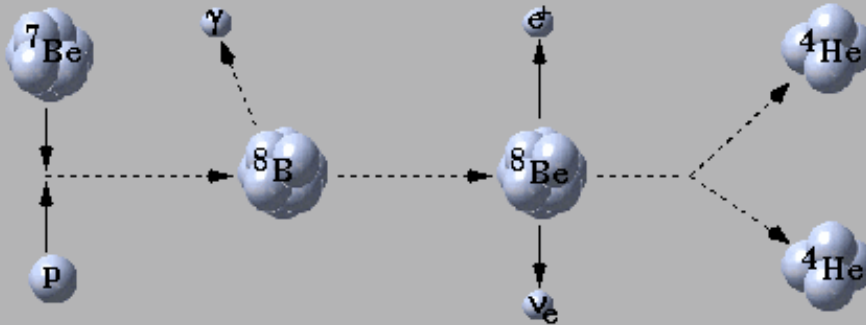
PPI



PPII



PPIII



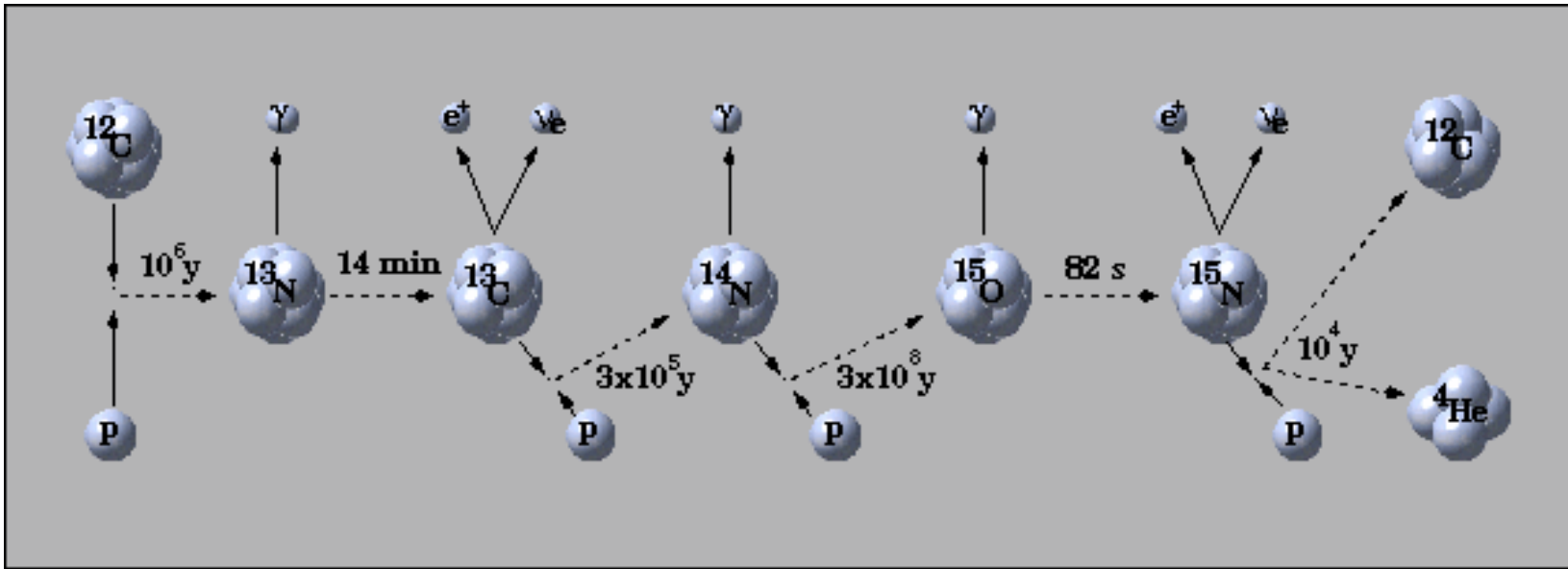
Relative importance of PPI and PPII chains (*branching ratios*) depend on conditions of H-burning (T, ρ , abundances). The transition from PPI to PPII occurs at temperatures in excess of 1.3×10^7 K.

Above 3×10^7 K the PPIII chain dominates over the other two, but another process takes over in this case.

The CNO Cycle

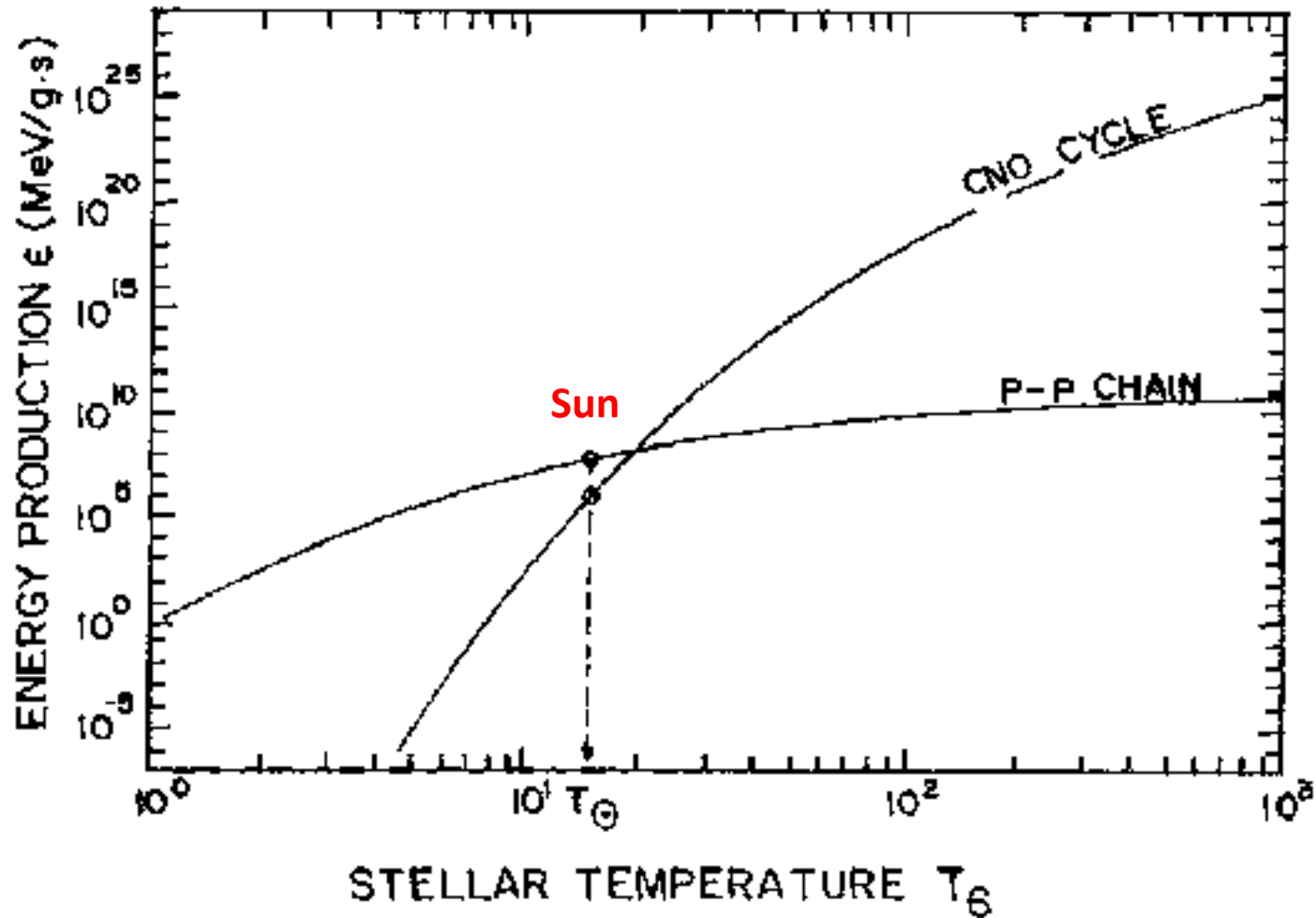
- Remember: $[Z] < 2\%$, most abundant CNO
- CNO induce a chain of H-burning reactions in which they act as catalysts
- The process is known as the CNO Cycle. There are alternative names in the literature:
 1. The CNO bi-cycle
 2. The CNOF cycle
 3. The CN and NO cycles
 4. The CN and NO bi-cycles
- In this course we will just refer to it all as the CNO cycle

The main CNO branch



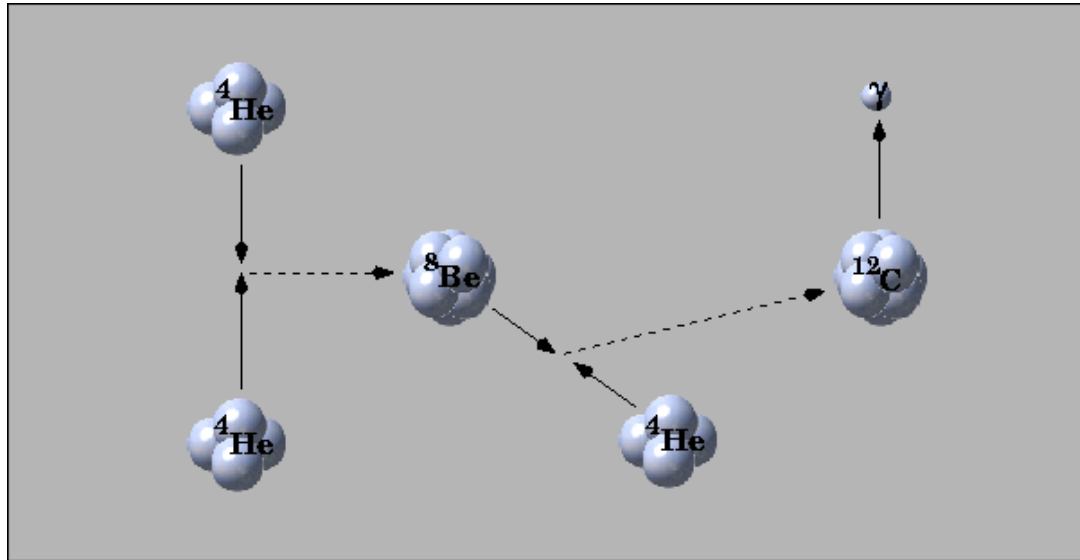
- 1 $^{12}\text{C} + \text{p} \rightarrow ^{13}\text{N} + \gamma$
- 2 $^{13}\text{N} \rightarrow ^{13}\text{C} + e^+ + \nu_e$
- 3 $^{13}\text{C} + \text{p} \rightarrow ^{14}\text{N} + \gamma$
- 4 $^{14}\text{N} + \text{p} \rightarrow ^{15}\text{O} + \gamma$
- 5 $^{15}\text{O} \rightarrow ^{15}\text{N} + e^+ + \nu_e$
- 6 $^{15}\text{N} + \text{p} \rightarrow ^{12}\text{C} + ^4\text{He}$

In the steady state case, the abundances of isotopes must take values such that the isotopes which react more slowly have higher abundance. The slowest reaction is p capture by ^{14}N . Hence most of ^{12}C is converted to ^{14}N .

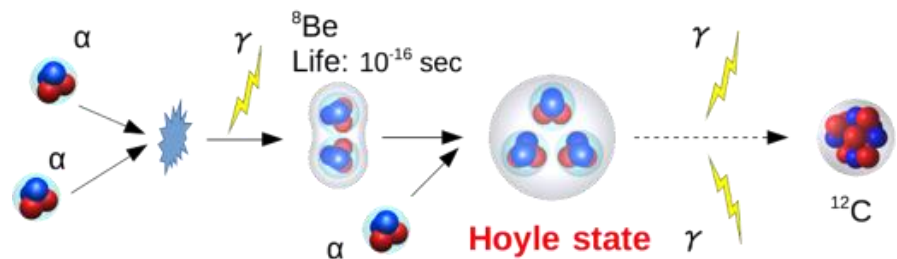
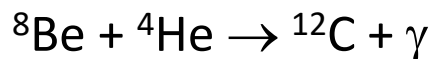


CNO cycle for stars with $M > 1.2 M_{\odot}$ dominant

Helium Burning: the triple - α reaction



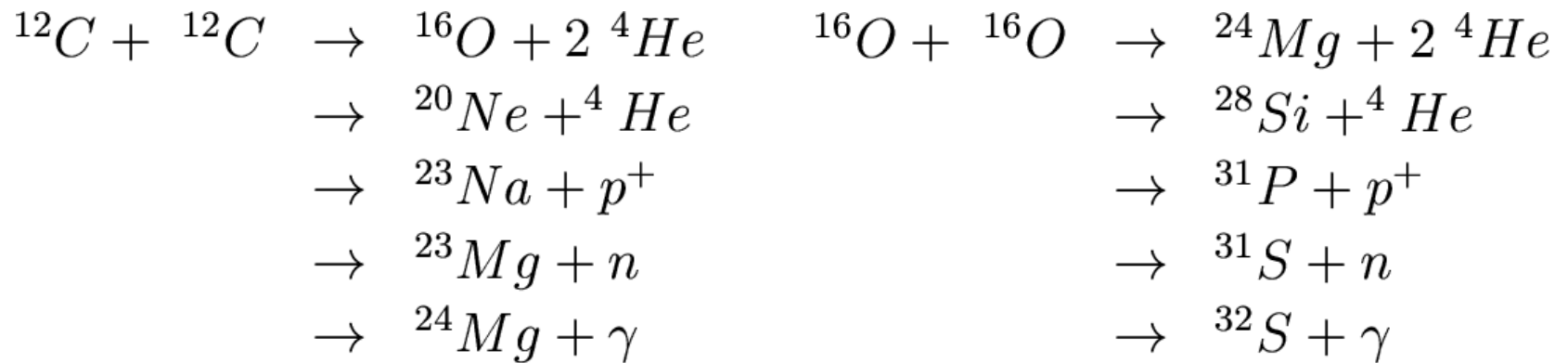
- Simplest reaction should be the fusion of two helium nuclei
- ${}^4\text{He} + {}^4\text{He} \rightarrow {}^8\text{Be}$
- There is no stable configuration with $A = 8$.
- Beryllium isotope ${}^8\text{Be}$ has a lifetime of only 2.6×10^{-16} s
- But a third helium nucleus can be added to ${}^8\text{Be}$ before decay, forming ${}^{12}\text{C}$ by the “triple-alpha” reaction



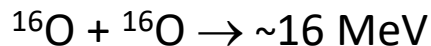
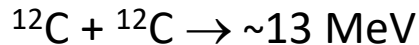
Carbon and oxygen burning

Carbon burning (fusion of two C nuclei) requires temperatures above 5×10^8 K, and oxygen burning in excess of 10^9 K.

Interactions of C and O nuclei are negligible – as at the intermediate temperatures required by the coulomb barrier the C nuclei are quickly destroyed by interacting with themselves



The branching ratios for these reactions are temperature dependent probabilities.

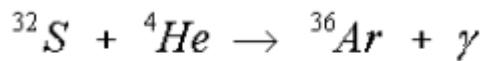
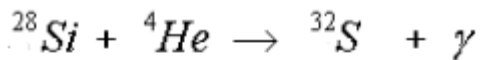


These reactions produce p, n, He, which are immediately captured by heavy nuclei, thus many isotopes created by secondary reactions.

Silicon burning

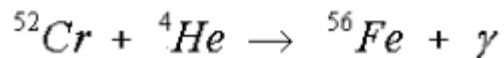
Two Si nuclei could fuse to create ^{56}Fe – the end of the fusion chain

But now very high Coulomb barrier, at T above O burning, but below that required for Si burning, **photodisintegration** takes place



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Si disintegration occurs around 3×10^9 K, and the light particles emitted are recaptured by other Si nuclei.

Although the reactions tend to a state of equilibrium, a leakage occurs towards the stable iron group nuclei (Fe, Co, Ni), which resist photodisintegration up to 7×10^9 K.

Summary - nuclear burning processes

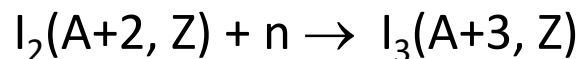
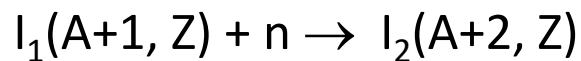
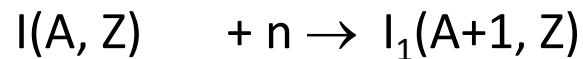
- Release of energy by consumption of nuclear fuel
- Rates of energy release vary enormously
- Nuclear processes can also absorb energy from radiation field

Nuclear Fuel	Process	$T_{\text{threshold}}$ 10^6K	Products	Energy per nucleon (MeV)
H	PP	~4	He	6.55
H	CNO	15	He	6.25
He	3α	100	C, O	0.61
C	C+C	600	O, Ne, Na, Mg	0.54
O	O+O	1000	Mg, S, P, Si	~0.3
Si	Nuc eq.	3000	Co, Fe, Ni	<0.18

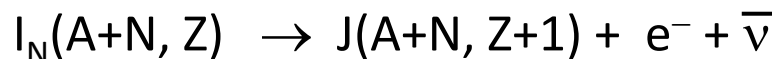
The *s*-process and *r*-process

Interaction between nuclei and free neutrons (neutron capture) – the neutrons are produced during C, O and Si burning.

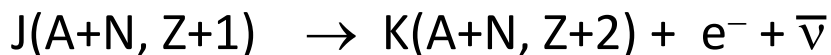
Neutrons capture by heavy nuclei is not limited by the Coulomb barrier – so could proceed at relatively low temperatures. The obstacle is the scarcity of free neutrons. If enough neutrons available, chain of reactions possible:



If a radioactive isotope is formed it will undergo β – decay, creating new element.



If new element stable, it will resume neutron capture, otherwise undergo series of β -decays



The *s*-process and *r*-process

- Two types of reactions and two types of nuclei
 1. Neutron captures and β -decays
 2. Stable and unstable nuclei
- Stable nuclei may undergo only neutron captures
- Unstable ones may undergo both
- Outcome depending on the timescales for the two processes
- Neutron capture reactions may proceed more ***slowly*** or more ***rapidly*** than the competing β – decays
- The different chains of reactions and products are called the ***s – process*** and ***r – process***

The *s*-process and *r*-process

