

# Solar models, neutrinos and helioseismology

Aldo Serenelli

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Institute of  
Space Sciences



# Outline

## (Standard) Solar models

Solar neutrinos

current status

origin of solar luminosity

Helioseismology

solar abundance problem

energy transport (opacity)

opacity – composition degeneracy

Near-term prospects for CNO neutrinos measurement

# Basic facts about the Sun

Luminosity – $L_{\odot}$	$3.842 \times 10^{33}$ erg/s
Radius – $R_{\odot}$	$6.9598 \times 10^{10}$ cm
Mass – $M_{\odot}$	$1.9891 \times 10^{33}$ g
Age (solar system oldest meteorites) – $\tau_{\odot}$	$4.57 \times 10^9$ yr
Surface metal to hydrogen abundance ratio – $(Z/X)_{\odot}$	0.018 – 0.024

The Sun is a typical middle aged low-mass star

# Quantitative predictions: Solar models

Compute the evolution of a  $1M_{\odot}$  stellar model

initial homogenous composition

evolve up to  $\tau_{\odot}$

Adjust initial composition (two parameters) and one free parameter of convection to match

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Age (solar system oldest meteorites) – $\tau_{\odot}$	$4.57 \times 10^9 \text{ yr}$
Surface metal to hydrogen abundance ratio – $(Z/X)_{\odot}$	$0.018 - 0.024$

Result: model of present day structure of the Sun – internal thermodynamic profiles

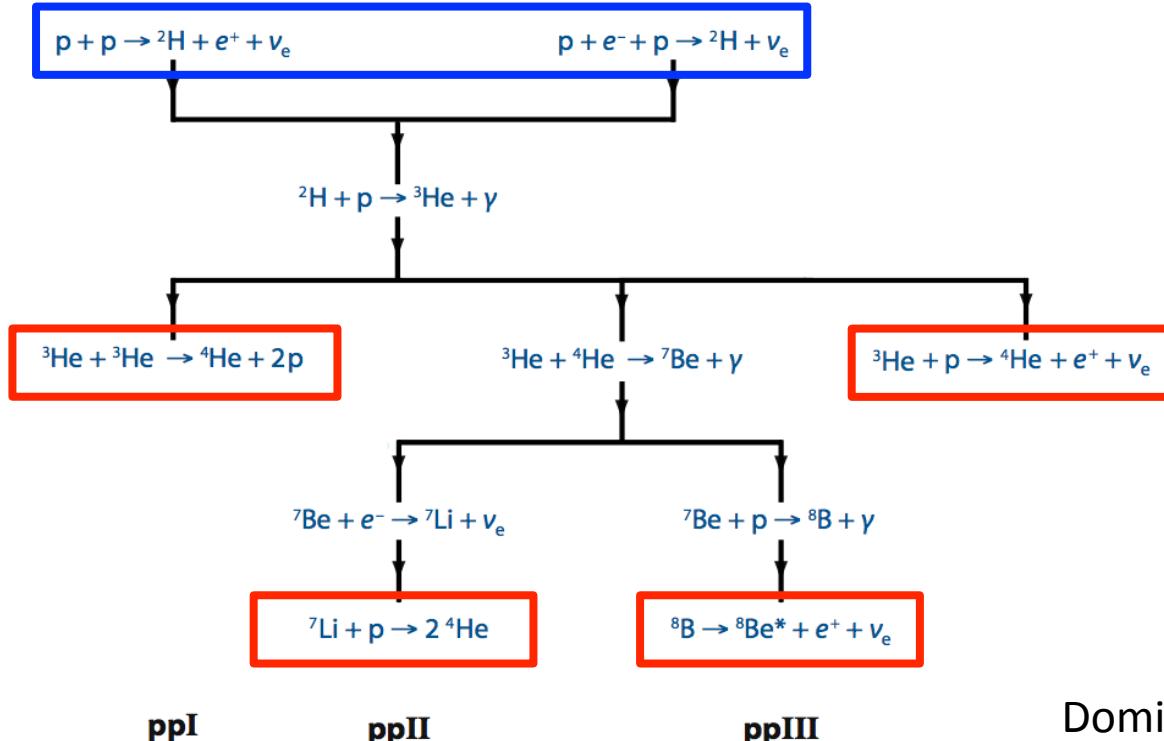
Two main ways to test models: solar neutrinos & helioseismology

# Hydrogen burning & solar neutrinos



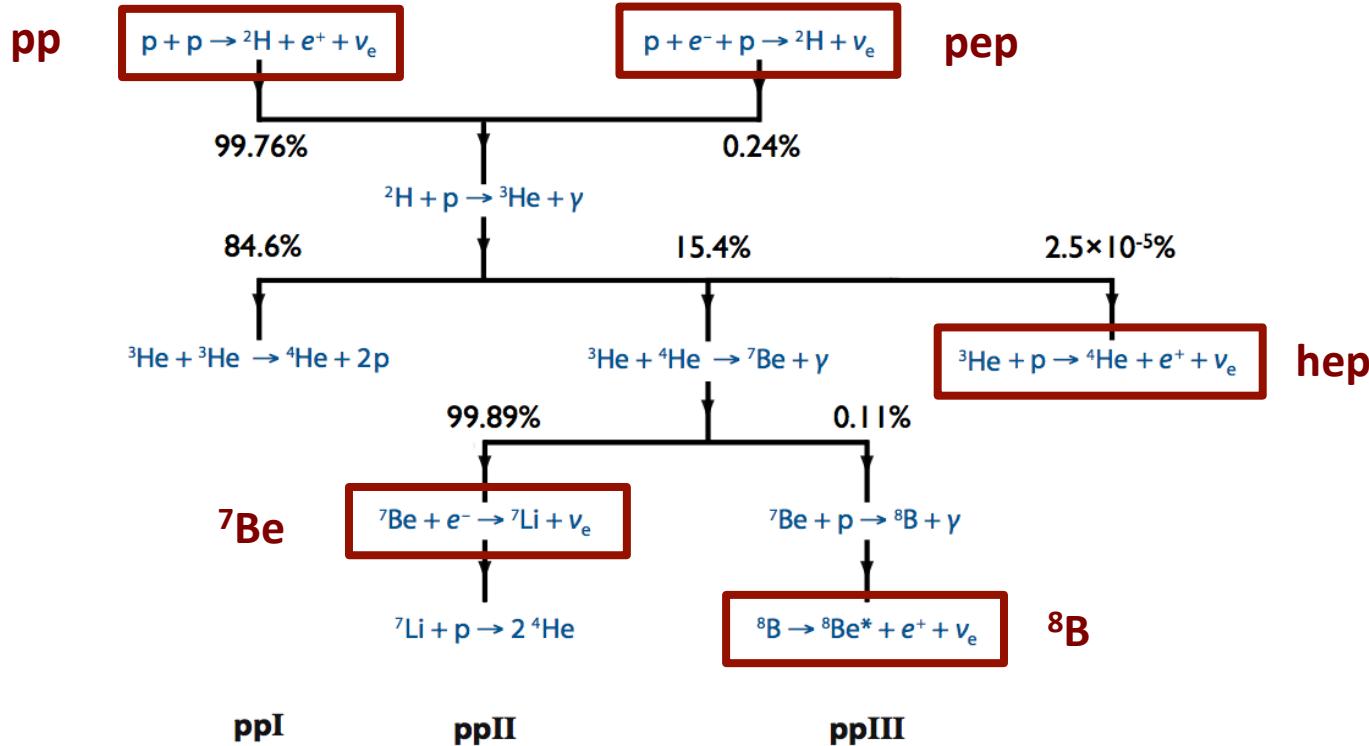
# Hydrogen burning – pp-chains

Several paths to:  $4p \rightarrow ^4He + 2\nu_e + 2e^+$



# Hydrogen burning – pp-chains

Five different neutrino sources



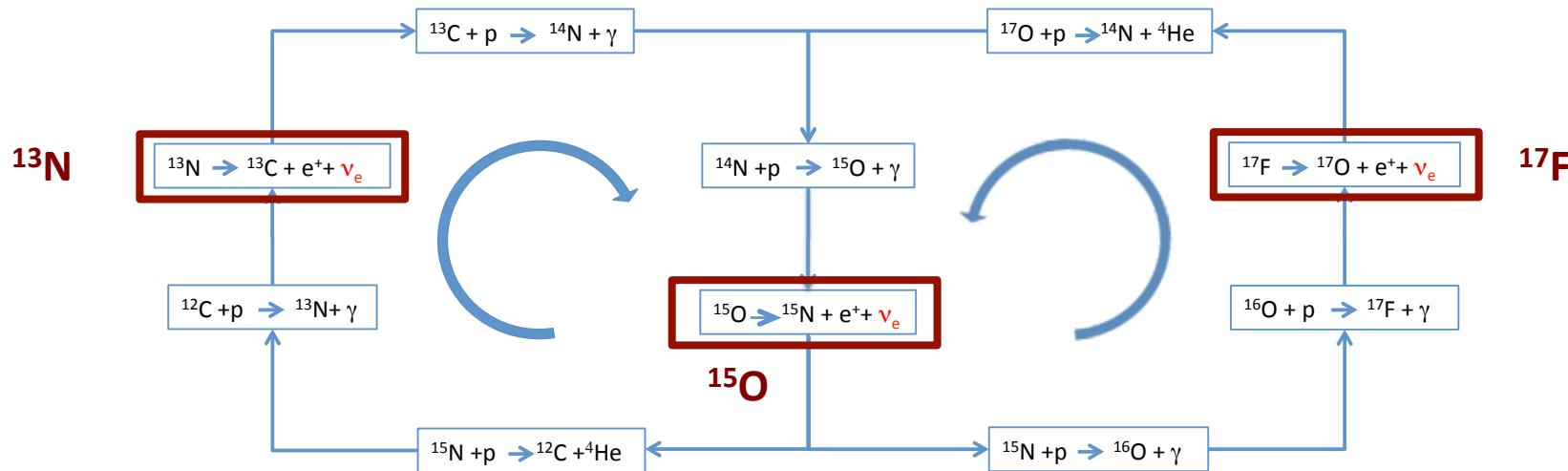
# Hydrogen burning – CNO-cycles

CNO-cycles – dominates at higher temperatures (more massive stars)

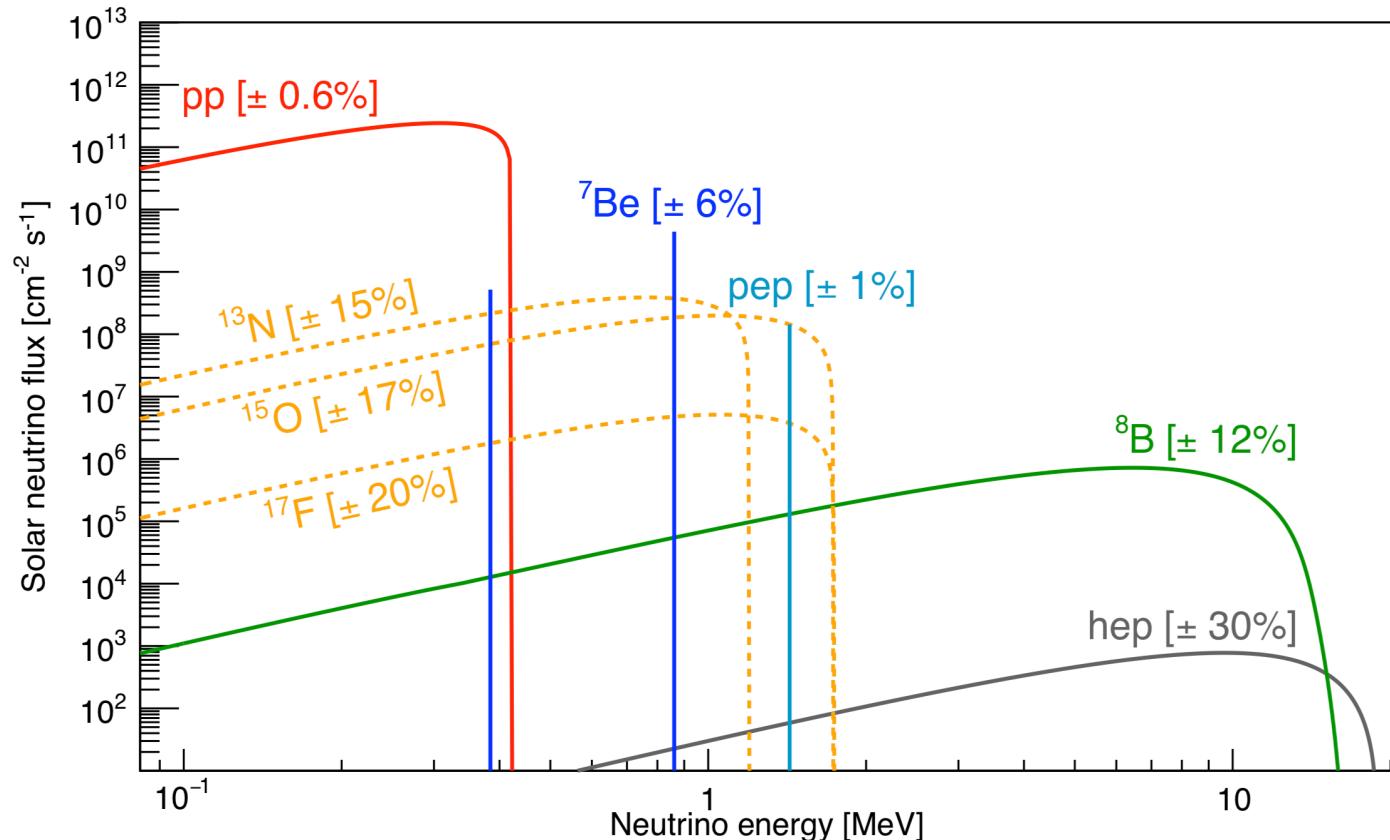
Marginal in the Sun

Three different neutrino sources

Catalyzed by C+N(+O) abundance



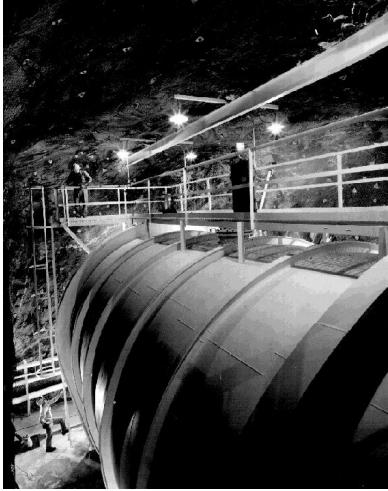
# Solar neutrino spectrum



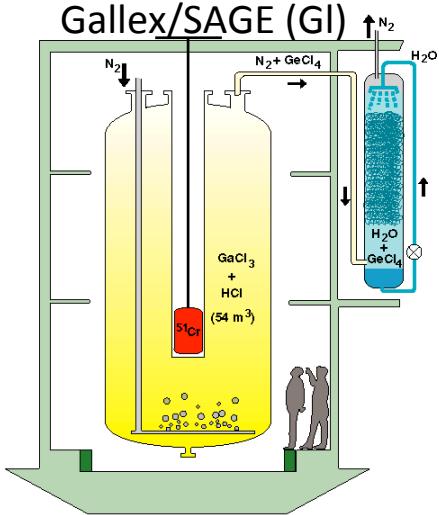
Flux	SFII-GS98 [ $\text{cm}^{-2} \text{s}^{-1}$ ]
pp	$5.98 \times 10^{10}$
pep	$1.44 \times 10^8$
hep	$8.04 \times 10^3$
${}^7\text{Be}$	$5.00 \times 10^9$
${}^8\text{B}$	$5.58 \times 10^6$
${}^{13}\text{N}$	$2.96 \times 10^8$
${}^{15}\text{O}$	$2.23 \times 10^8$
${}^{17}\text{F}$	$5.52 \times 10^6$

# 50 years of solar neutrinos experiments

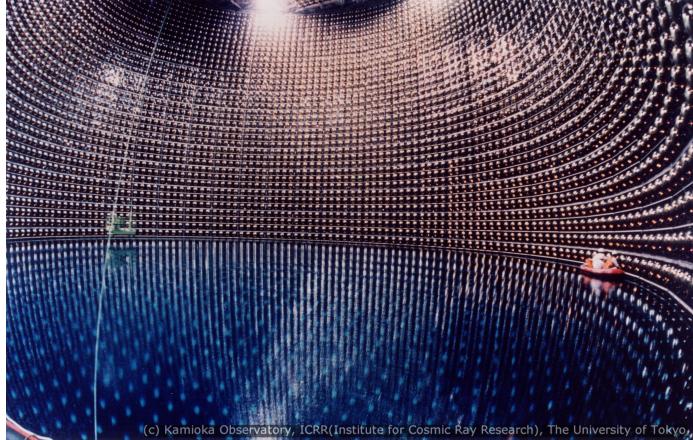
Homestake (Cl)



Gallex/SAGE (GI)



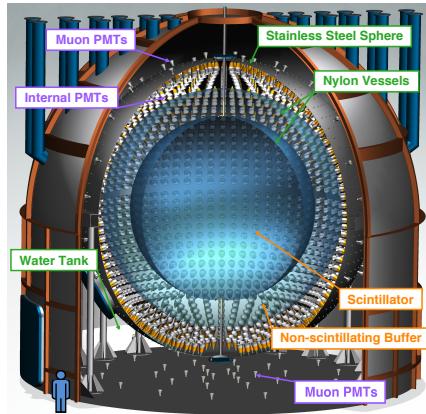
Kamiokande/SuperK (H2O)



SNO (D2O)



Borexino (Scint.)



# Coming full circle

How does the Sun shine?  
Which is the solar core temperature?

Solar vs as sensitive  
probes of solar core

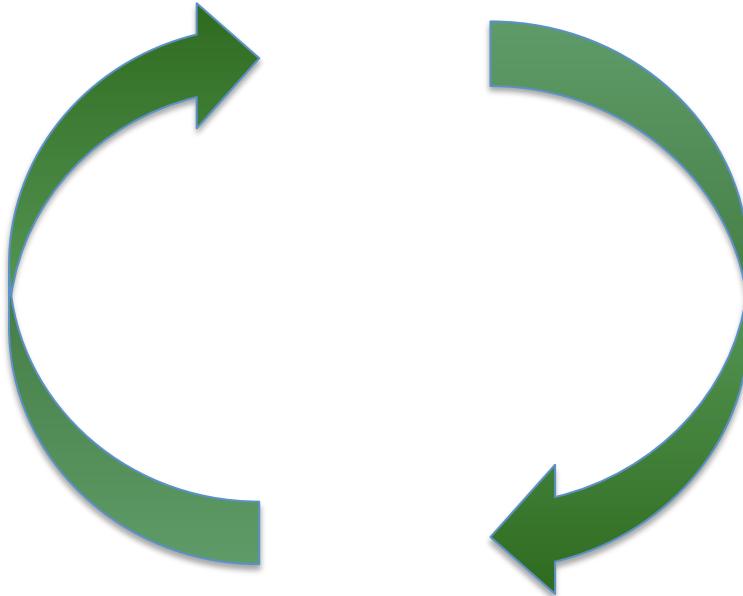
SuperK & SNO  
ν oscillations

Radiochemical experiments  
Kamiokande

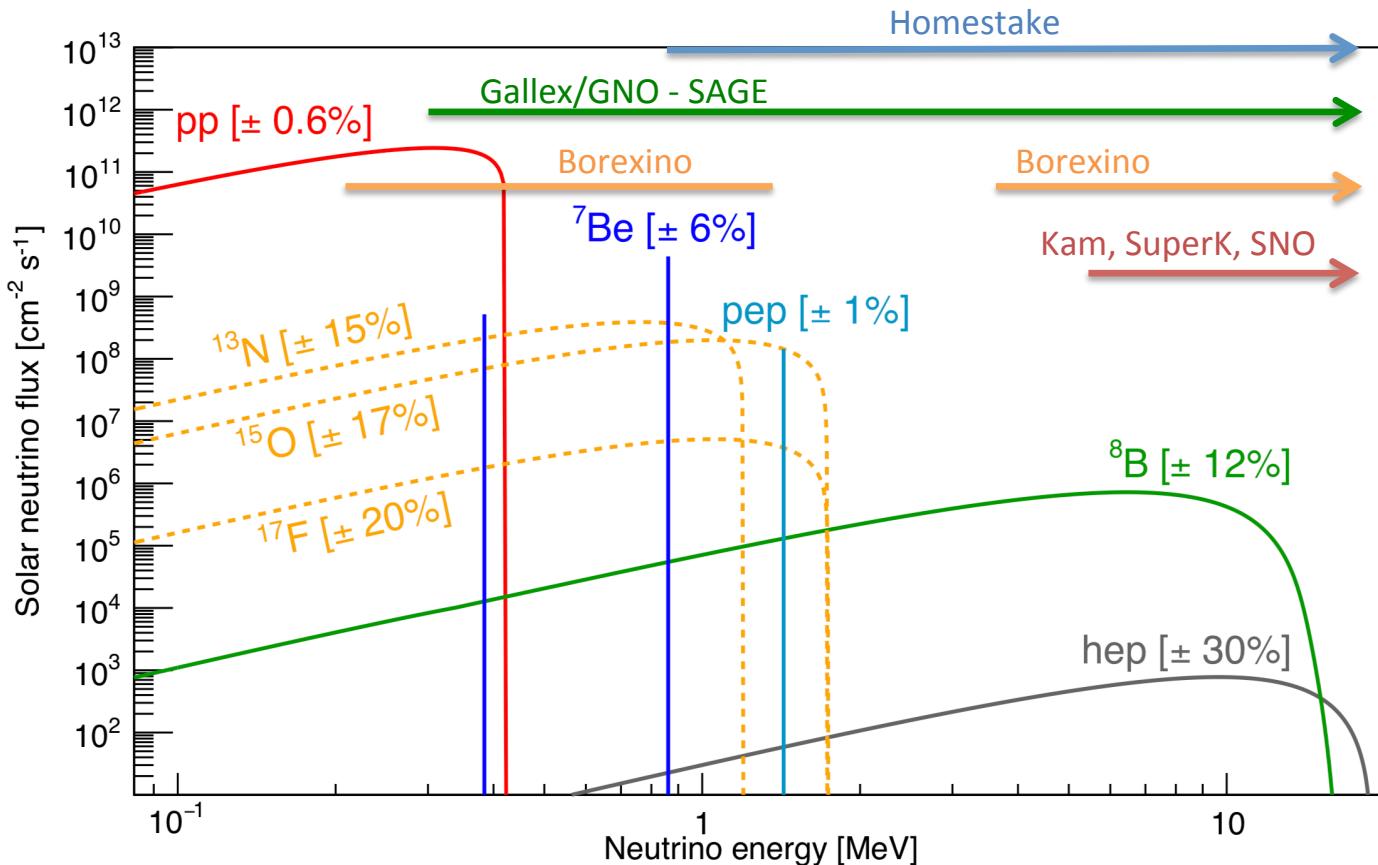
Solar neutrino problem

Standard model(s) crises!

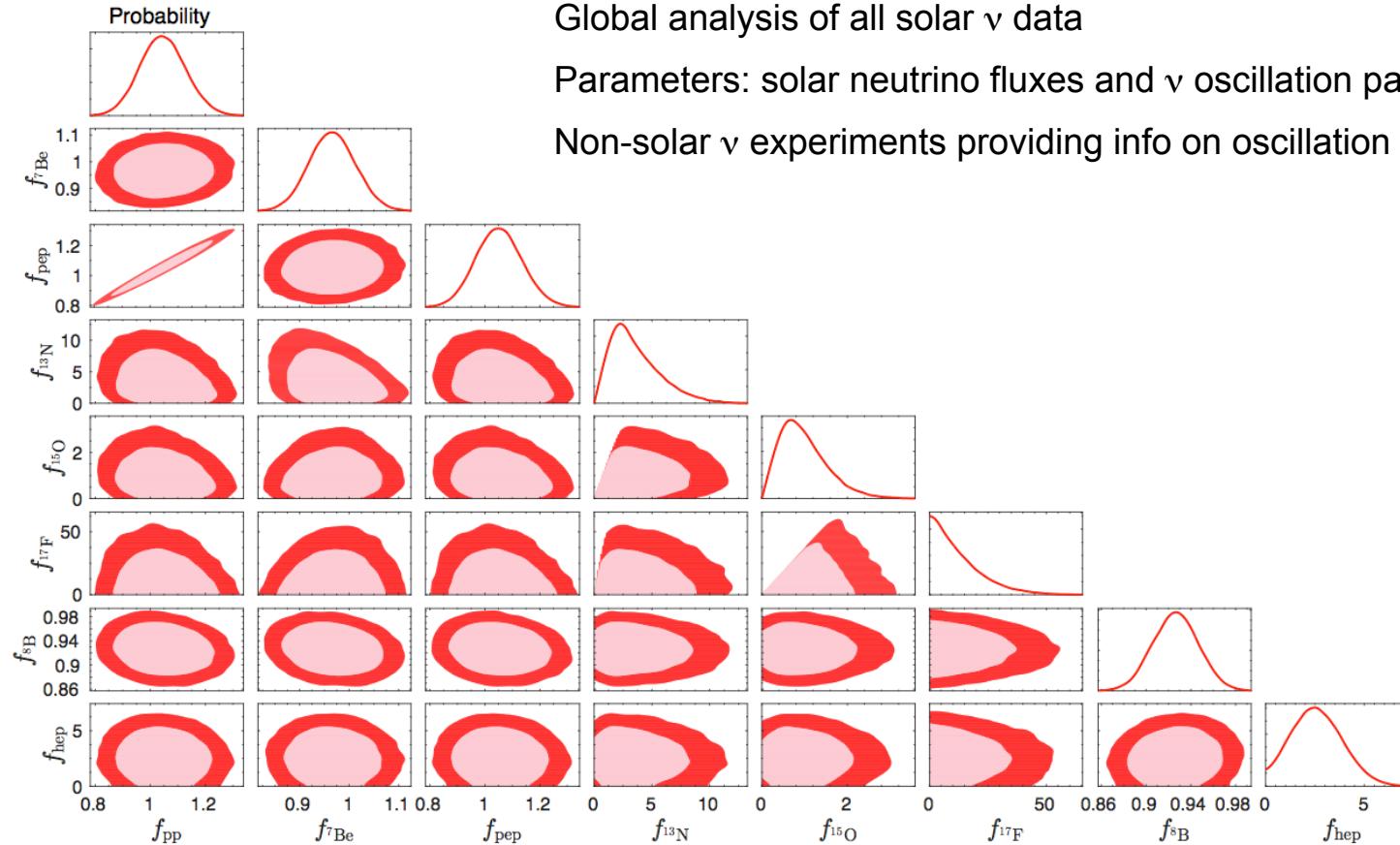
Helioseismic inference of  
solar structure → agreement with solar models  
→ Physics solution (circa 1996)



# Solar neutrino spectrum

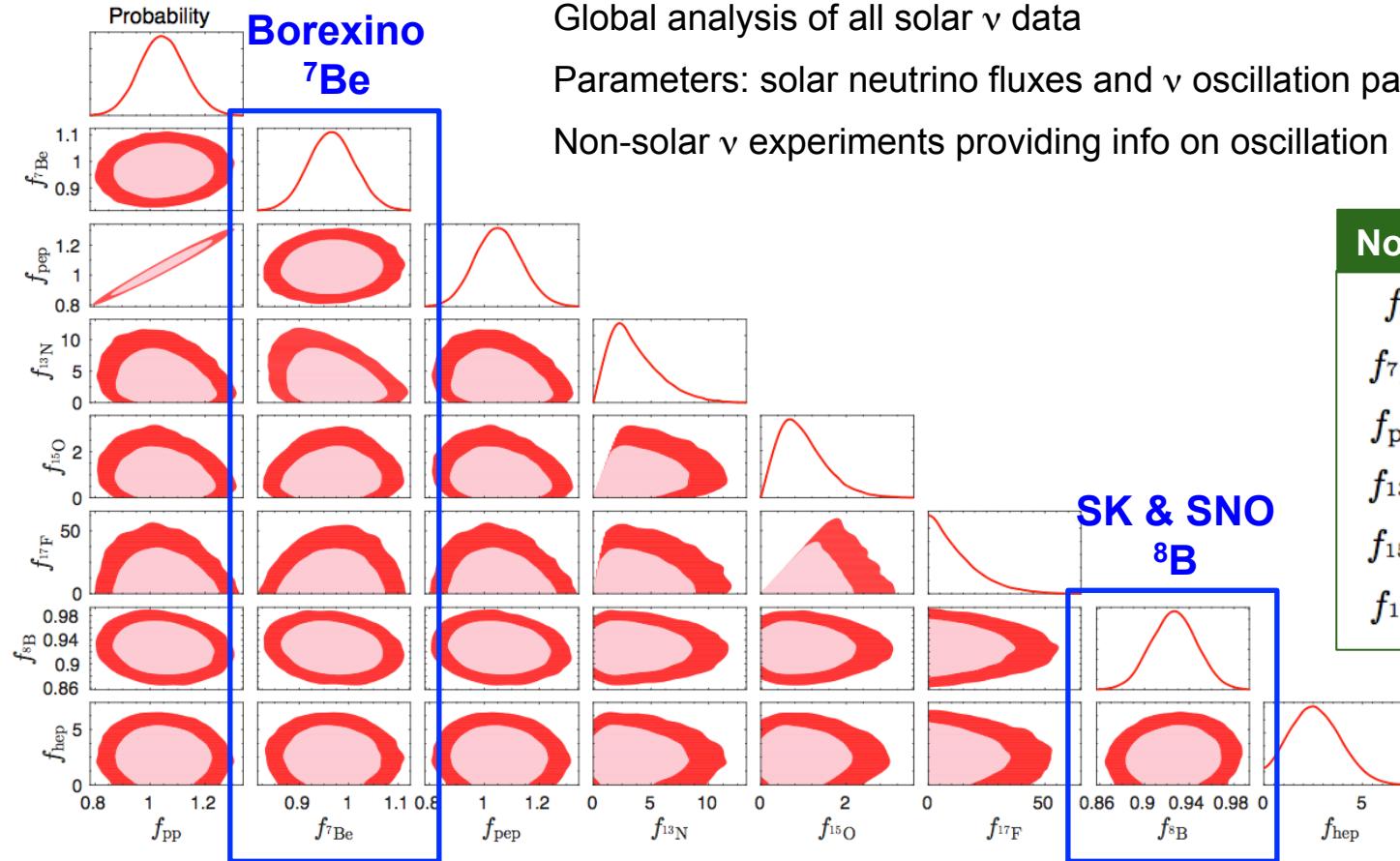


# How does the Sun shine?



Bergstrom et al. 2016

# How does the Sun shine?



No luminosity constraint

$$f_{\text{pp}} = 1.04 \pm 0.08 [{}^{+0.22}_{-0.20}] ,$$

$$f_{^7\text{Be}} = 0.97^{+0.04}_{-0.05} [\pm 0.12] ,$$

$$f_{\text{pep}} = 1.05 \pm 0.08 [{}^{+0.23}_{-0.20}] ,$$

$$f_{^{13}\text{N}} = 1.7^{+2.8}_{-1.0} [{}^{+8.4}_{-1.6}] ,$$

$$f_{^{15}\text{O}} = 0.6^{+0.7}_{-0.4} [\leq 2.6] ,$$

$$f_{^{17}\text{F}} \leq 15 [47] .$$

Bergstrom et al. 2016

# How does the Sun shine?

Simple linear relation linking all neutrino fluxes to nuclear energy generation rate

$\alpha_i$  depend only on Q values of reactions and shape of neutrino spectra

$$\frac{L_{\text{nuc}}}{4\pi(\text{AU})^2} = \sum_i \alpha_i \Phi_i$$

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**Purely experimental result – no solar model information**

$$\frac{L_{\text{pp-chain}}}{L_\odot} = 1.03^{+0.08}_{-0.07} [{}^{+0.21}_{-0.18}] \quad \text{and} \quad \frac{L_{\text{CNO}}}{L_\odot} = 0.008^{+0.005}_{-0.004} [{}^{+0.014}_{-0.007}] .$$

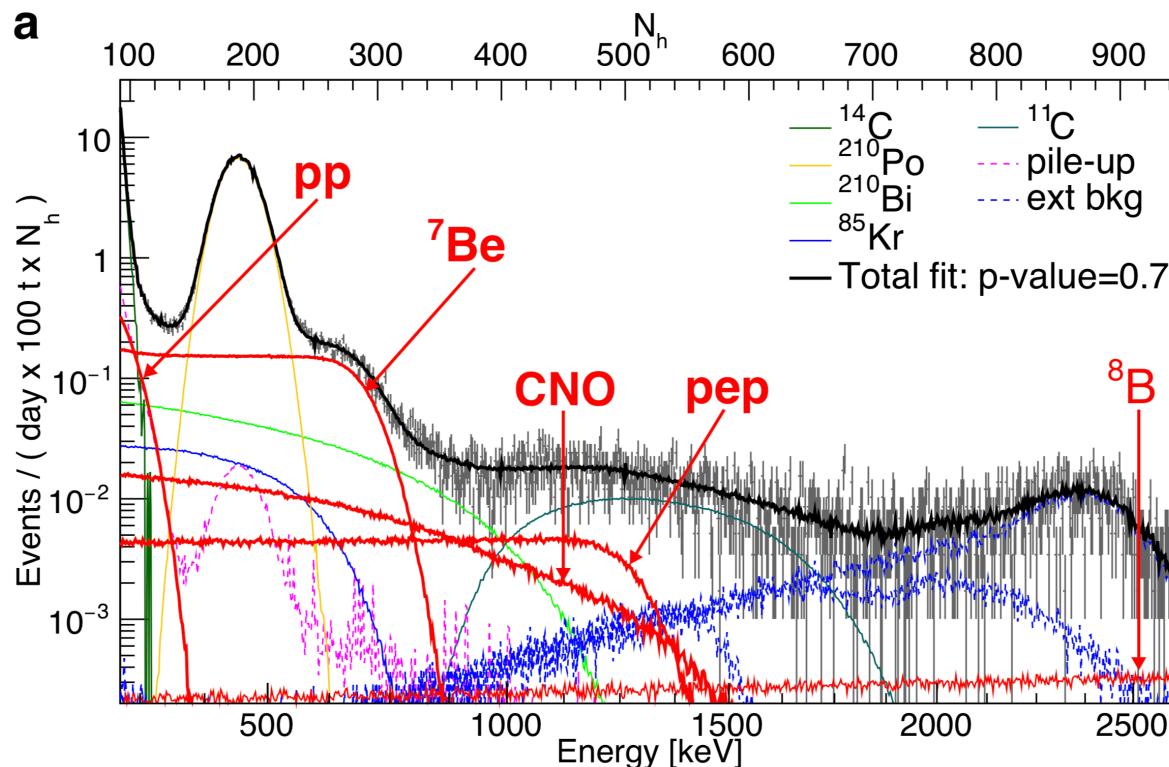
$$\frac{L_{\text{nuc}}(\text{neutrino-inferred})}{L_\odot} = 1.04 [{}^{+0.07}_{-0.08}] [{}^{+0.20}_{-0.18}]$$

Bergstrom et al. 2016

# Latest results from Borexino

Data taking for more than 10 years

Observed neutrino spectrum – Caccianaga et al. 2018 (Borexino Collaboration)



Some highlights from Borexino

$^7\text{Be}$  measured to 3%

pp measured to 10%

pep measured to 15%

$^8\text{B}$  measured to lowest energy

Caccianaga et al. 2018

# Latest results from Borexino

**Table 2 | Borexino experimental solar-neutrino results**

Solar neutrino	Rate (counts per day per 100 t)	Flux ( $\text{cm}^{-2} \text{s}^{-1}$ )	Flux-SSM predictions ( $\text{cm}^{-2} \text{s}^{-1}$ )
$p\bar{p}$	$134 \pm 10^{+6}_{-10}$	$(6.1 \pm 0.5^{+0.3}_{-0.5}) \times 10^{10}$	$5.98(1.0 \pm 0.006) \times 10^{10}$ (HZ) $6.03(1.0 \pm 0.005) \times 10^{10}$ (LZ)
$^7\text{Be}$	$48.3 \pm 1.1^{+0.4}_{-0.7}$	$(4.99 \pm 0.11^{+0.06}_{-0.08}) \times 10^9$	$4.93(1.0 \pm 0.06) \times 10^9$ (HZ) $4.50(1.0 \pm 0.06) \times 10^9$ (LZ)
$p\bar{p}\text{p}$ (HZ)	$2.43 \pm 0.36^{+0.15}_{-0.22}$	$(1.27 \pm 0.19^{+0.08}_{-0.12}) \times 10^8$	$1.44(1.0 \pm 0.01) \times 10^8$ (HZ) $1.46(1.0 \pm 0.009) \times 10^8$ (LZ)
$p\bar{p}\text{p}$ (LZ)	$2.65 \pm 0.36^{+0.15}_{-0.24}$	$(1.39 \pm 0.19^{+0.08}_{-0.13}) \times 10^8$	$1.44(1.0 \pm 0.01) \times 10^8$ (HZ) $1.46(1.0 \pm 0.009) \times 10^8$ (LZ)
$^8\text{B}_{\text{HER-I}}$	$0.136^{+0.013+0.003}_{-0.013-0.003}$	$(5.77^{+0.56+0.15}_{-0.56-0.15}) \times 10^6$	$5.46(1.0 \pm 0.12) \times 10^6$ (HZ) $4.50(1.0 \pm 0.12) \times 10^6$ (LZ)
$^8\text{B}_{\text{HER-II}}$	$0.087^{+0.080+0.005}_{-0.010-0.005}$	$(5.56^{+0.52+0.33}_{-0.64-0.33}) \times 10^6$	$5.46(1.0 \pm 0.12) \times 10^6$ (HZ) $4.50(1.0 \pm 0.12) \times 10^6$ (LZ)
$^8\text{B}_{\text{HER}}$	$0.223^{+0.015+0.006}_{-0.016-0.006}$	$(5.68^{+0.39+0.03}_{-0.41-0.03}) \times 10^6$	$5.46(1.0 \pm 0.12) \times 10^6$ (HZ) $4.50(1.0 \pm 0.12) \times 10^6$ (LZ)
CNO	$<8.1$ (95% C.L.)	$<7.9 \times 10^8$ (95% C.L.)	$4.88(1.0 \pm 0.11) \times 10^8$ (HZ) $3.51(1.0 \pm 0.10) \times 10^8$ (LZ)
hep	$<0.002$ (90% C.L.)	$<2.2 \times 10^5$ (90% C.L.)	$7.98(1.0 \pm 0.30) \times 10^3$ (HZ) $8.25(1.0 \pm 0.12) \times 10^3$ (LZ)

Experimental results after accounting for  $\nu$  oscillations

Caccianaga et al. 2018

## Oscillation parameters

	Normal Ordering (best fit)	
	bfp $\pm 1\sigma$	$3\sigma$ range
$\sin^2 \theta_{12}$	$0.306^{+0.012}_{-0.012}$	$0.271 \rightarrow 0.345$
$\theta_{12}/^\circ$	$33.56^{+0.77}_{-0.75}$	$31.38 \rightarrow 35.99$
$\sin^2 \theta_{23}$	$0.441^{+0.027}_{-0.021}$	$0.385 \rightarrow 0.635$
$\theta_{23}/^\circ$	$41.6^{+1.5}_{-1.2}$	$38.4 \rightarrow 52.8$
$\sin^2 \theta_{13}$	$0.02166^{+0.00075}_{-0.00075}$	$0.01934 \rightarrow 0.02392$
$\theta_{13}/^\circ$	$8.46^{+0.15}_{-0.15}$	$7.99 \rightarrow 8.90$
$\frac{\Delta m_{21}^2}{10^{-5} \text{ eV}^2}$	$7.50^{+0.19}_{-0.17}$	$7.03 \rightarrow 8.09$
$\frac{\Delta m_{3\ell}^2}{10^{-3} \text{ eV}^2}$	$+2.524^{+0.039}_{-0.040}$	$+2.407 \rightarrow +2.643$

Esteban et al. 2017

## Borexino experimental result

$$\frac{L_{\text{nuc}}(\text{neutrino-inferred})}{L_\odot} = 1.01 \left[ {}^{+0.09}_{-0.11} \right]$$

# How does the Sun shine?

$$\frac{L_{\text{nuc}}(\text{neutrino-inferred})}{L_{\odot}} = 1.01 \left[ {}^{+0.09}_{-0.11} \right]$$

$$\frac{L_{\text{nuc}}(\text{neutrino-inferred})}{L_{\odot}} = 1.04 \left[ {}^{+0.07}_{-0.08} \right] \left[ {}^{+0.20}_{-0.18} \right]$$

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Standard solar models

$$L_{\odot} = \int \frac{\partial L}{\partial m} dm = \int (\varepsilon_{\text{nuc},\nu} + \varepsilon_g) dm = \int \varepsilon_{\text{nuc},\nu} dm \longrightarrow L_{\odot} = L_{\text{nuc}}$$

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But, what if there is an energy source/sink not recognized in standard solar models ...

$$L_{\odot} = \int (\varepsilon_{\text{nuc},\nu} + \varepsilon_g + \varepsilon_x) dm = L_{\text{nuc}} + L_x \longrightarrow L_{\odot} \neq L_{\text{nuc}}$$

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$$\frac{L_{\text{nuc}}(\text{neutrino-inferred})}{L_{\odot}} = 1.01 \left[ {}^{+0.09}_{-0.11} \right]$$

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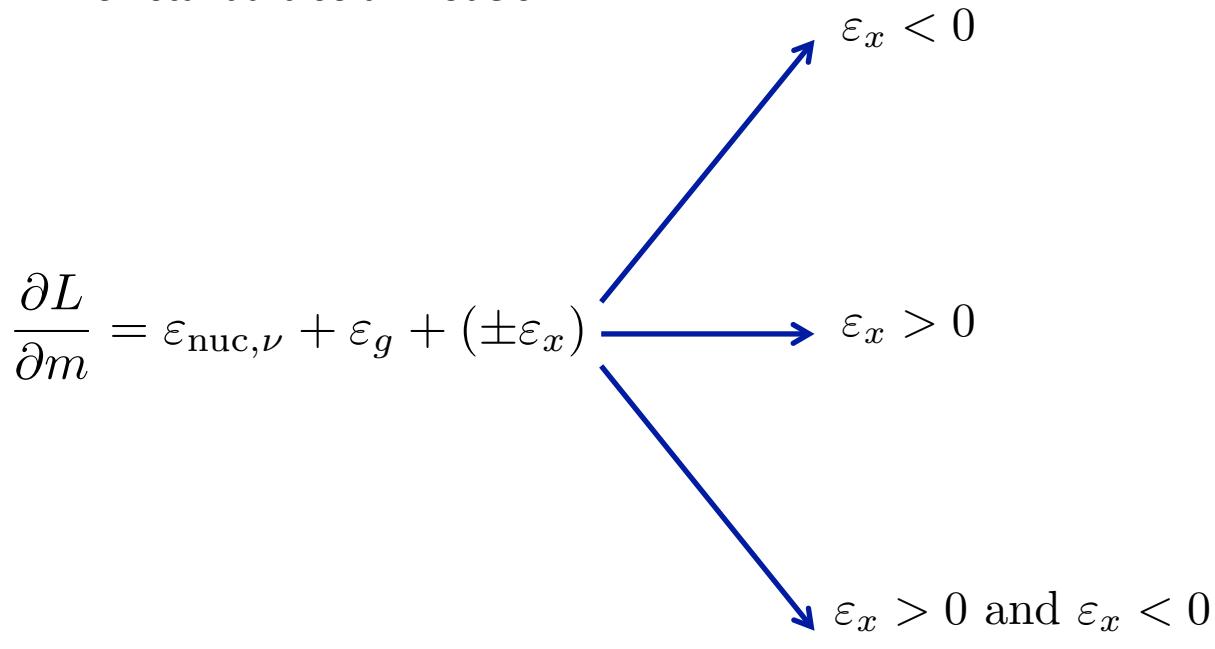
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A complete measurement of solar neutrino fluxes offers the only model independent limit on non-standard energy sources in the Sun (and stars)

Present-day limit: 8%

# How does the Sun shine?

Non standard solar models



$x$  particles escape from the star, carrying away energy, analogous to neutrinos, e.g. axions –  $L_\odot < L_{\text{nuc}}$

$x$  particles produce energy in the star, e.g. self-annihilating DM –  $L_\odot > L_{\text{nuc}}$

transport of energy within the star,  $< 0$  where energy is extracted and  $> 0$  where it is deposited, e.g. non-annihilating DM

talks today by P. Scott and A. Vincent

# Helioseismology

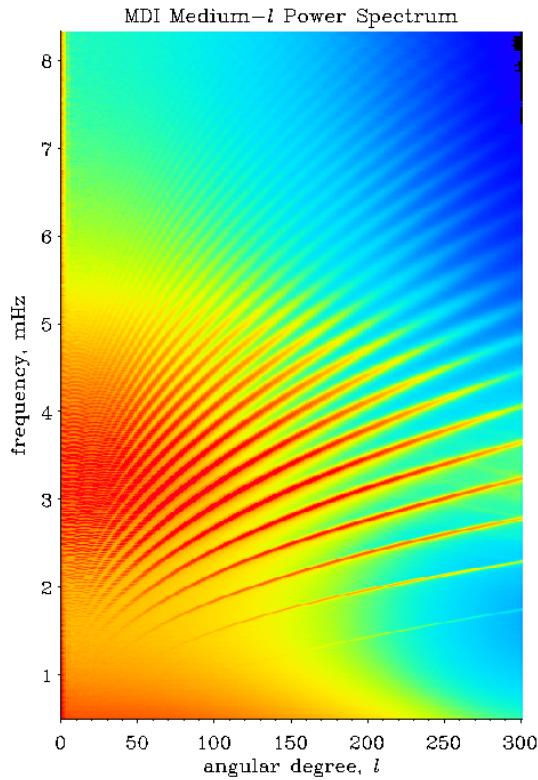
Global solar oscillations – acoustic waves

Key quantities:  $c^2$ ,  $p$ ,  $r$ ,  $\Gamma_1$

$10^5$  non-radial individual modes

cm/s amplitudes in radial velocity

ppm amplitudes in brightness



# Helioseismology

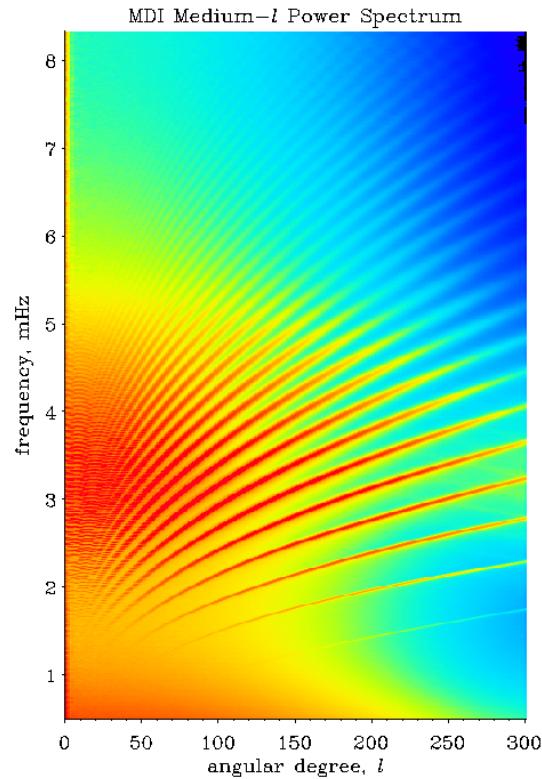
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Key quantities:  $c^2$ ,  $p$ ,  $r$ ,  $\Gamma_1$

$10^5$  non-radial individual modes

cm/s amplitudes in radial velocity

ppm amplitudes in brightness



Inversion methods allow reconstructing internal profiles

$$\frac{\delta\omega_i}{\omega_i} = \int K_{c^2, \rho}^i(r) \frac{\delta c^2}{c^2}(r) dr + \int K_{\rho, c^2}^i(r) \frac{\delta \rho}{\rho}(r) dr + F_{surf}(\omega_i)$$

# Back to solar models: solar composition

Luminosity – $L_{\odot}$	$3.842 \times 10^{33}$ erg/s
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Mass – $M_{\odot}$	$1.9891 \times 10^{33}$ g
Age (solar system oldest meteorites) – $\tau_{\odot}$	$4.57 \times 10^9$ yr
Surface metal to hydrogen abundance ratio – $(Z/X)_{\odot}$	0.018 – 0.024

# Solar composition

Change of paradigm in solar composition:

Grevesse & Sauval 1998 → Asplund et al. 2005, 2009, Scott et al. 2015 – Caffau et al. 2011

- 3D solar atmosphere models
- refined atomic data and line selection
- non-LTE treatment of line formation

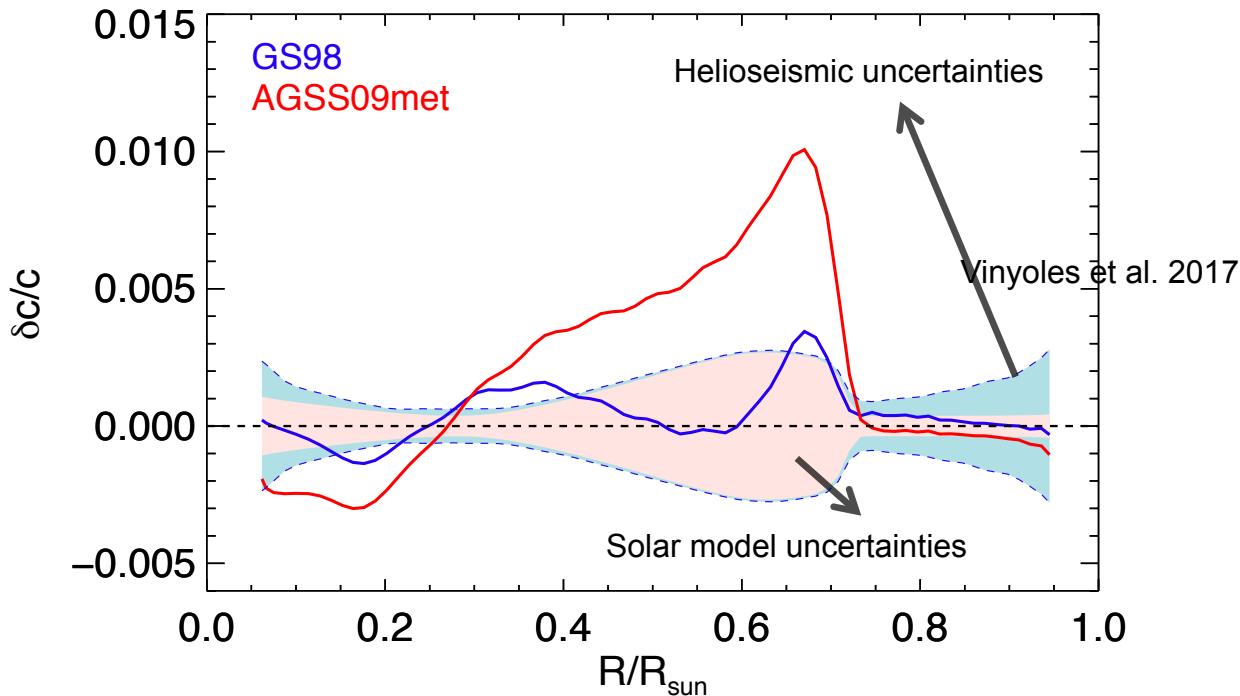
Elem.	GS98	AGSS09met	Change
C	$8.52 \pm 0.06$	$8.52 \pm 0.05$	23%
N	$7.92 \pm 0.06$	$7.83 \pm 0.05$	23%
O	$8.83 \pm 0.06$	$8.69 \pm 0.05$	38%
Ne	$8.08 \pm 0.06$	$7.93 \pm 0.10$	41%
Mg	$7.58 \pm 0.01$	$7.53 \pm 0.01$	12%
Si	$7.56 \pm 0.01$	$7.51 \pm 0.01$	12%
S	$7.20 \pm 0.06$	$7.15 \pm 0.02$	12%
Fe	$7.50 \pm 0.06$	$7.45 \pm 0.01$	12%
$(Z/X)_{\odot}$	0.0229	0.0178	29%

Impact of SSM calibration			
	$L_{\odot}$	$R_{\odot}$	$(Z/X)_{\odot}$
$\alpha_{\text{mlt}}$	0.06	-0.19	0.06
$Y_{\text{ini}}$	2.35	0.56	0.08
$Z_{\text{ini}}$	-0.73	-0.14	1.11

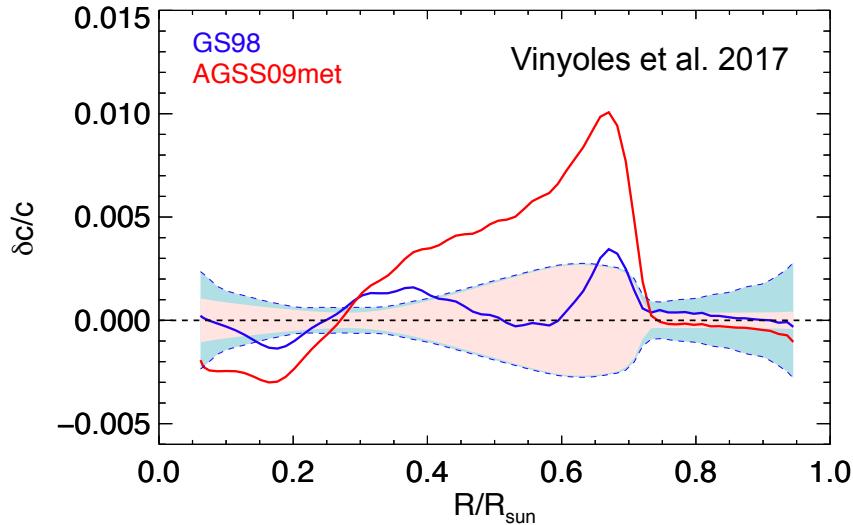
Z/X determines solar model  
composition:  $Z_{\text{ini}}$  &  $Y_{\text{ini}}$

# Solar models - helioseismology

## Sound speed inversion



# Solar models - helioseismology



Other helioseismic probes:

Density inversion

Depth of convective envelope

Solar core through frequency ratios

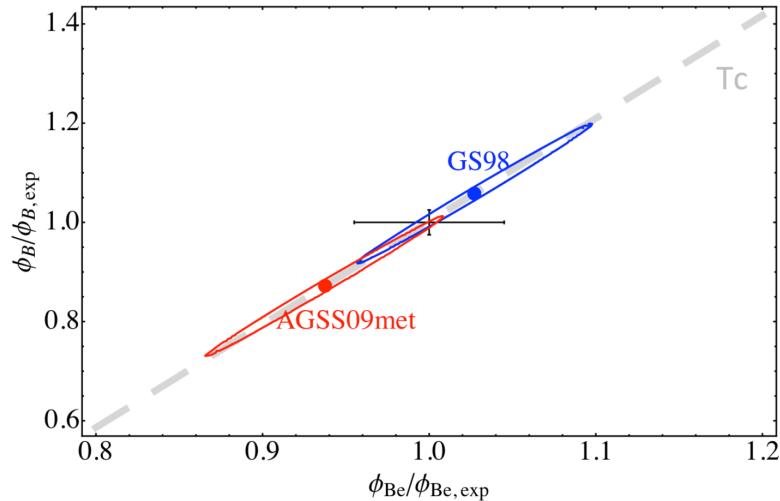
Qnt.	B16-GS98	B16-AGSS09met	Solar
$Y_S$	$0.2426 \pm 0.0059$	$0.2317 \pm 0.0059$	$0.2485 \pm 0.0035$
$R_{\text{CZ}}/R_{\odot}$	$0.7116 \pm 0.0048$	$0.7223 \pm 0.0053$	$0.713 \pm 0.001$
$\langle \delta c/c \rangle$	$0.0005^{+0.0006}_{-0.0002}$	$0.0021 \pm 0.001$	$0^{\text{a}}$

} 2-3  $\sigma$  discrepancy for low Z

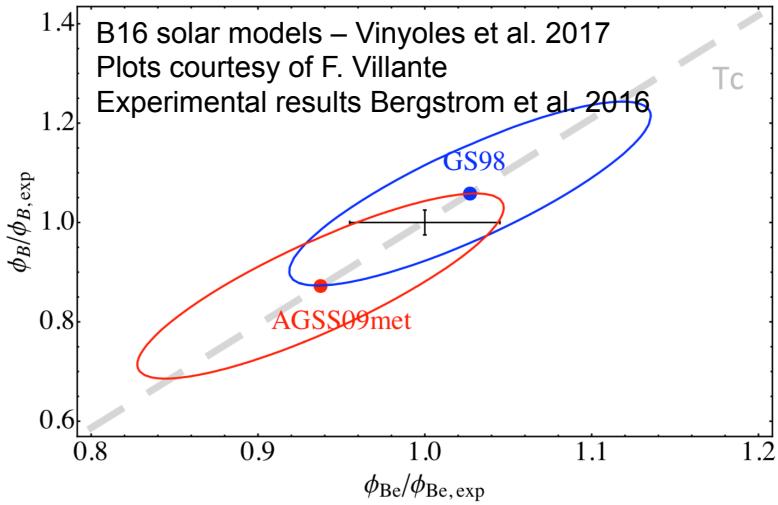
Discrepancy between solar models (new composition) and all helioseismic inferences:  
solar abundance problem (lots of literature about this)

# Solar composition: neutrinos

Environmental (temperature) uncertainties  
**composition, opacity**, age, luminosity, etc

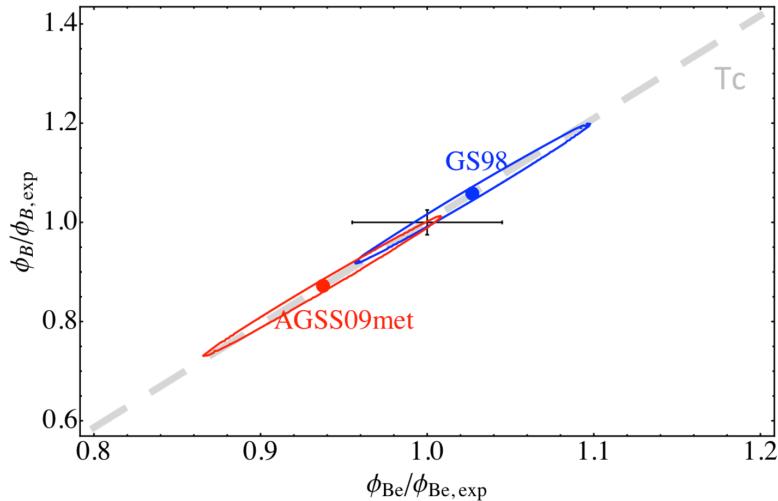


+ nuclear rate uncertainties

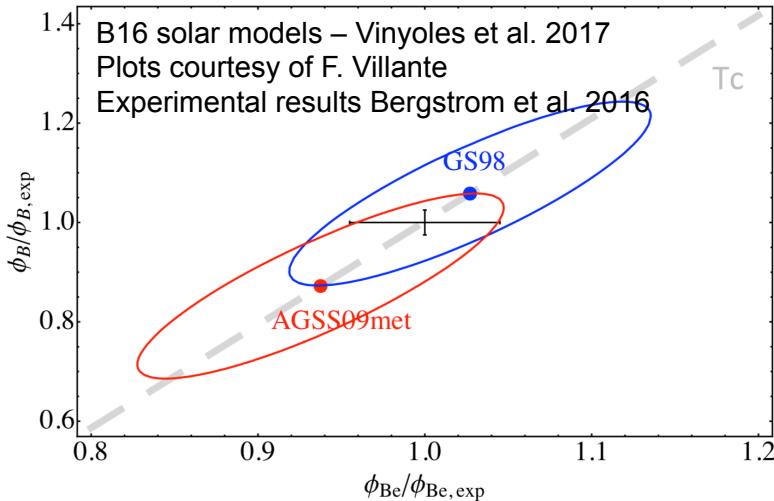


# Solar composition: neutrinos

Environmental (temperature) uncertainties  
**composition, opacity**, age, luminosity, etc



+ nuclear rate uncertainties



Composition → affects pp-chain fluxes through  $T_c$  change

- determines opacity
- pp-fluxes sensitive to opacity (i.e. temperature, only indirectly to composition)
- composition and atomic opacities are degenerate in pp-chain fluxes (and helioseismology)

# Solar models – overall status

Case	dof	GS98		AGSS09met	
		$\chi^2$	p-value ( $\sigma$ )	$\chi^2$	p-value ( $\sigma$ )
$Y_S + R_{CZ}$ only	2	0.9	0.5	6.5	2.1
$\delta c/c$ only	30	58.0	3.2	76.1	4.5
$\delta c/c$ no-peak	28	34.7	1.4	50.0	2.7
$\Phi(^7\text{Be}) + \Phi(^8\text{B})$	2	0.2	0.3	1.5	0.6
All $\nu$ -fluxes	8	6.0	0.5	7.0	0.6
Global	40	65.0	2.7	94.2	4.7
Global no-peak	38	40.5	0.9	67.2	3.0

High-Z models favored by helioseismology

Vinyoles et al. 2017

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$\nu$ -fluxes comparably good agreement  
(model uncertainties dominate)  
because they are from pp-chains

Vinyoles et al. 2017

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Vinyoles et al. 2017

Global comparison favors high-Z models

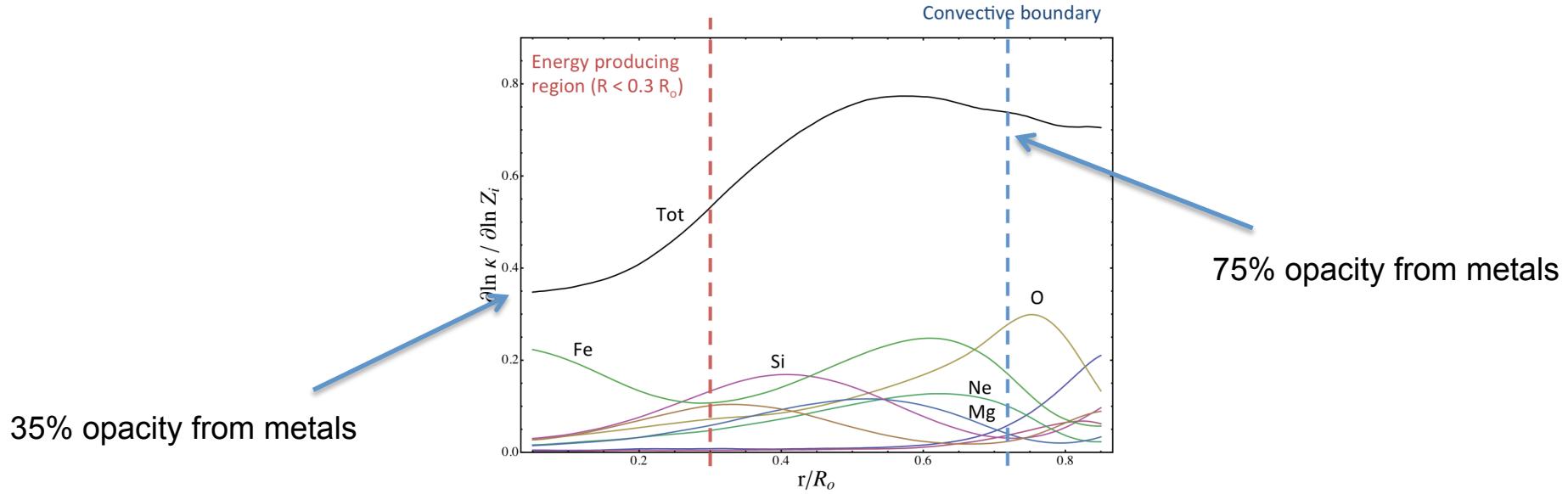
i.e. models with  $(P, \rho)$  or  $(T, \mu)$  profiles  
consistent with high-Z models

But interpretation in terms of solar composition is hampered by  
degeneracy between composition and opacity

# Energy transport: Metals & Opacity

In solar interior ( $R < 0.7 R_{\odot}$ ) energy transport by radiation – radiative opacity fundamental quantity

Lack of metals = lack of opacity : hard to disentangle



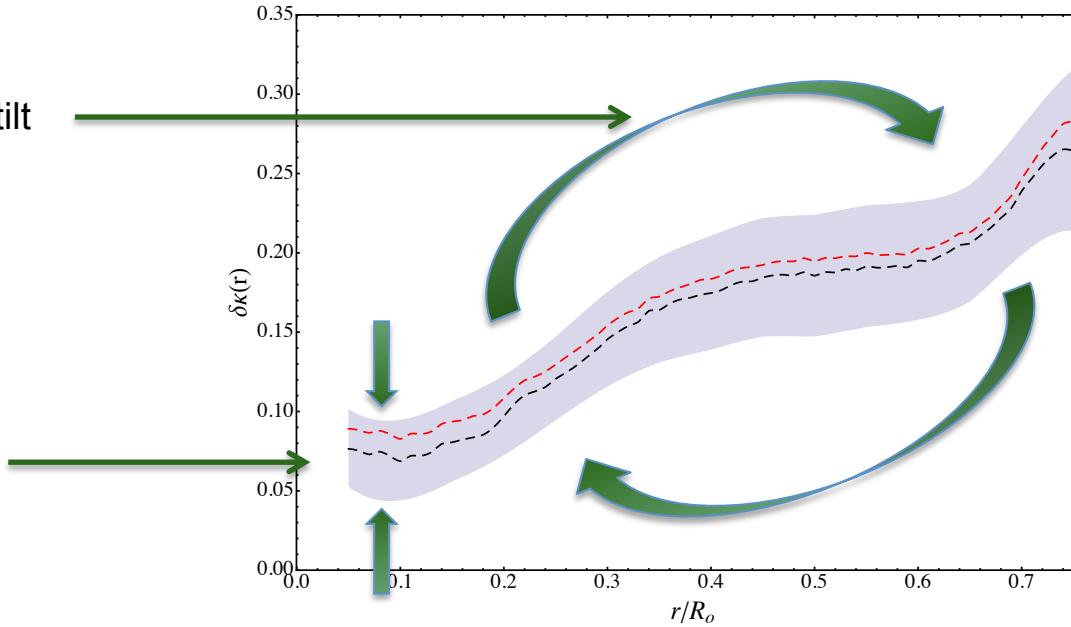
Intrinsic uncertainty + composition induced variation  
( $\delta$  = fractional variation)

$$\delta \kappa = \delta \kappa_I + \sum \frac{\partial \log \kappa}{\partial \log z_i} \delta z_i$$

# Solar opacity from $\nu S$ and helioseismology

Helioseismology fixes the tilt

Solar neutrinos and  $Y_S$  the scaling (core)

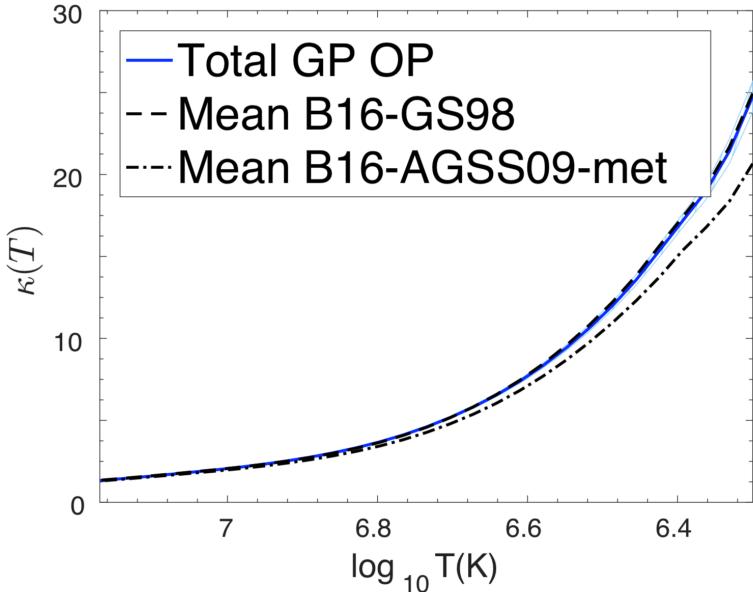


Villante et al. 2014

# Solar opacity from $\nu s$ and helioseismology

$\delta\kappa_I$  is an unknown function  $\rightarrow$  Gaussian Process

$$\delta\kappa = \delta\kappa_I + \sum \frac{\partial \log \kappa}{\partial \log z_i} \delta z_i$$



Song et al. 2018

Bayesian analysis – composition free to vary

Opacity solar profile (posterior distrib)

Very close to that from GS98 model (unsurprisingly)

If AGSS09 composition  $\rightarrow$  20% opacity increase at base of convective zone

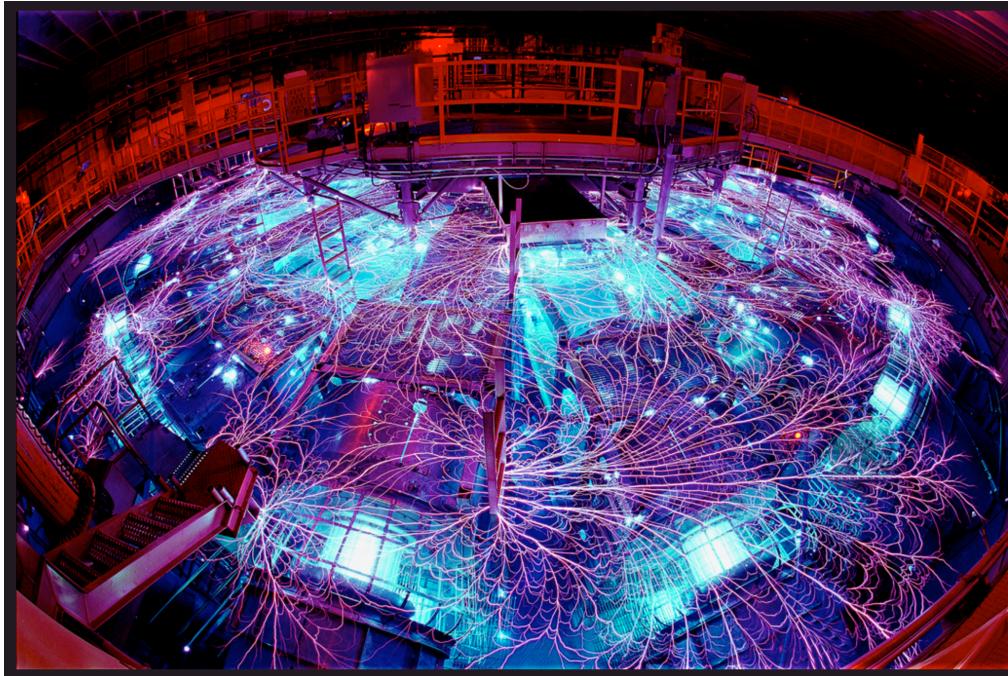
Few % opacity increase in solar core

Determine ‘effective’ opacity profile: cannot disentangle contributions (atomic, composition, other mechanisms, e.g. dark matter)

# Opacities – Experimental result

Z-pinch experiment at Sandia Lab

First ever measurement at conditions close to base of the solar convective envelope

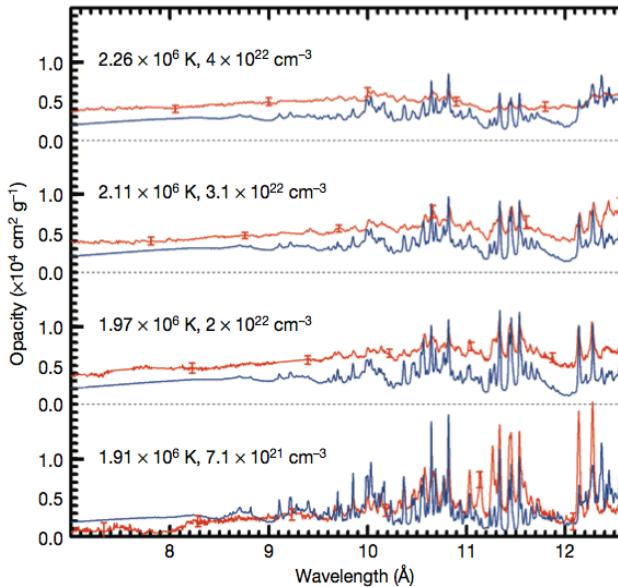


Bailey et al. 2015

# Opacities – Experimental result

First ever opacity measurement at conditions close to base of the solar convective envelope

Fe opacity @Sandia Lab -- > 7% increase of Rosseland mean opacity



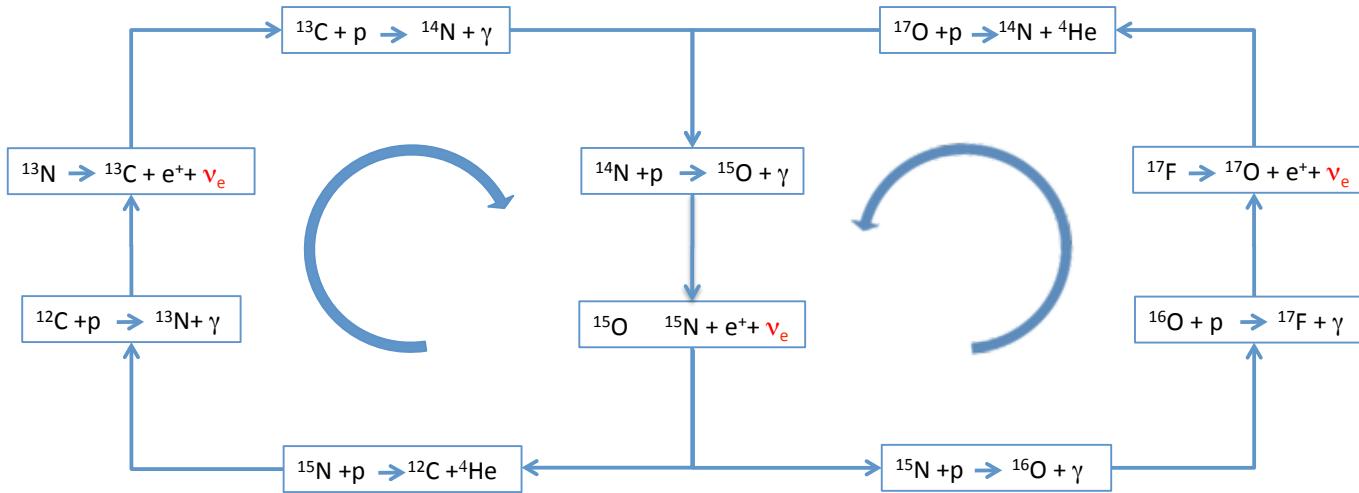
Bailey et al. 2015

$$T \sim T_{\text{CZ}}$$

$$N_e \sim 1/4 N_{e\text{CZ}}$$

Experimental hint of higher opacity than theoretical calculations predict – but situation unclear

# CN-vs and solar composition



CN-cycle marginal in the Sun → intrinsic changes in its rate do not alter background state (e.g. T,  $\rho$ , X)

→ linear dependence on core C+N abundance and  $S_{1,14}$  (for  $^{15}\text{O}$ )

$$\frac{\Phi(^{15}\text{O})_{\text{HZ}} - \Phi(^{15}\text{O})_{\text{LZ}}}{\Phi(^{15}\text{O})_{\text{LZ}}} \approx 40\%$$

# CN-vs and solar composition

Converting  $\Phi(^{15}\text{O})$  measurement into C+N core measurement – use  $\Phi(^8\text{B})$  as thermometer

$$\Phi(^8\text{B}) \propto T_c^{25} \longrightarrow \text{SuperK+SNO} \longrightarrow \delta T_c/T_c \approx 0.1\%$$

Removes many solar model uncertainties (environmental) in predictions of CN neutrino fluxes

# CN-vs and solar composition

Converting  $\Phi(^{15}\text{O})$  measurement into C+N core measurement – use  $\Phi(^8\text{B})$  as thermometer

$$\frac{\phi(^{15}\text{O})}{\phi(^{15}\text{O})_{\text{SSM}}} \Big/ \left[ \frac{\phi(^8\text{B})}{\phi^{\text{SSM}}(^8\text{B})} \right]^{0.785} = x_C^{0.794} x_N^{0.212} D^{0.172} \\ \times [L_\odot^{0.515} O^{-0.016} A^{0.308}] \\ \times [S_{11}^{-0.831} S_{33}^{0.342} S_{34}^{-0.685} S_{17}^{-0.785} S_{e7}^{0.785} S_{114}^{0.995}] \\ \times [x_O^{0.003} x_{\text{Ne}}^{-0.005} x_{\text{Mg}}^{-0.003} x_{\text{Si}}^{-0.001} x_S^{-0.001} x_{\text{Ar}}^{0.001} x_{\text{Fe}}^{0.003}]$$

- Temperature dependence
- Nuclear rates
- Temperature dependence (opacity)

Serenelli et al. 2013

# CN-vs and solar composition

Converting  $\Phi(^{15}\text{O})$  measurement into C+N core measurement – use  $\Phi(^8\text{B})$  as thermometer

$$\frac{\phi(^{15}\text{O})}{\phi(^{15}\text{O})^{\text{SSM}}} / \left[ \frac{\phi(^8\text{B})}{\phi^{\text{SSM}}(^8\text{B})} \right]^{0.785} = x_C^{0.794} x_N^{0.212} D^{0.172} \\ \times [L_{\odot}^{0.515} O^{-0.016} A^{0.308}] \\ \times [S_{11}^{-0.831} S_{33}^{0.342} S_{34}^{-0.685} S_{17}^{-0.785} S_{e7}^{0.785} S_{114}^{0.995}] \\ \times [x_O^{0.003} x_{\text{Ne}}^{-0.005} x_{\text{Mg}}^{-0.003} x_{\text{Si}}^{-0.001} x_S^{-0.001} x_{\text{Ar}}^{0.001} x_{\text{Fe}}^{0.003}]$$

→ Temperature dependence  
→ Nuclear rates  
→ Temperature dependence (opacity)

Reduces to

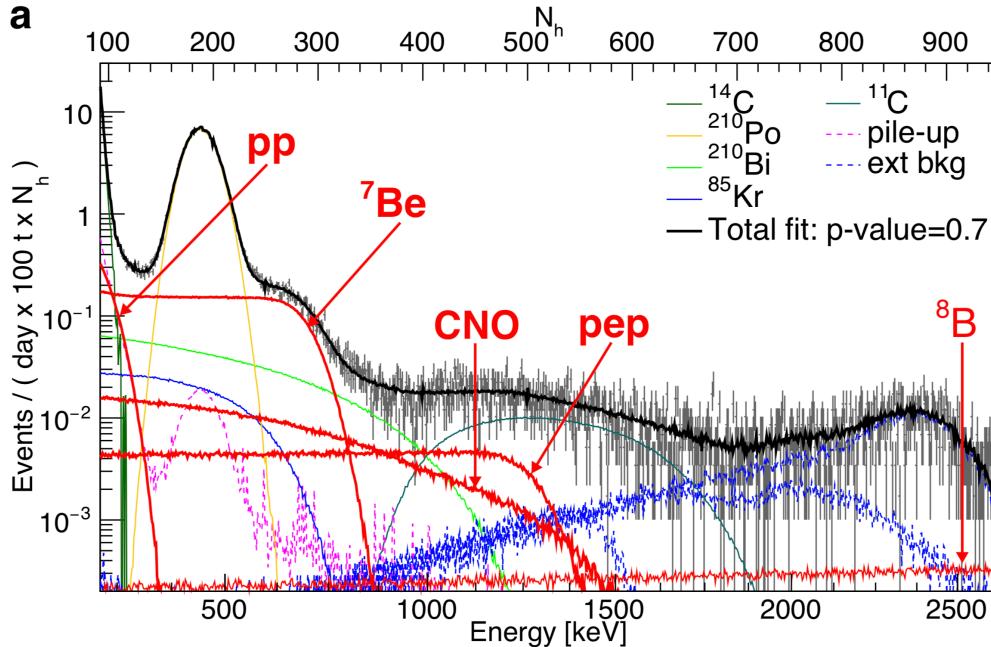
$$\frac{\phi(^{15}\text{O})}{\phi(^{15}\text{O})^{\text{SSM}}} / \left[ \frac{\phi(^8\text{B})}{\phi^{\text{SSM}}(^8\text{B})} \right]^{0.785} = \left[ \frac{C + N}{C^{\text{SSM}} + N^{\text{SSM}}} \right] (1 \pm 0.4\% \text{ (env)} \pm 2.6\% \text{ (D)} \pm 10\% \text{ (nucl)})$$

Nuclear uncertainty dominant:  $S_{1,14}$  (7%) &  $S_{17}$  (5%) – can be potentially reduced further

(Almost) direct measurement of C+N in solar core

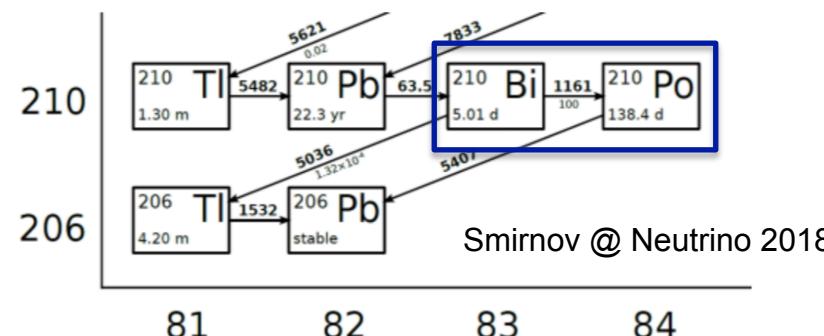
Serenelli et al. 2013

# CN-vs at Borexino



CN flux hidden below  $^{210}\text{Bi}$  background

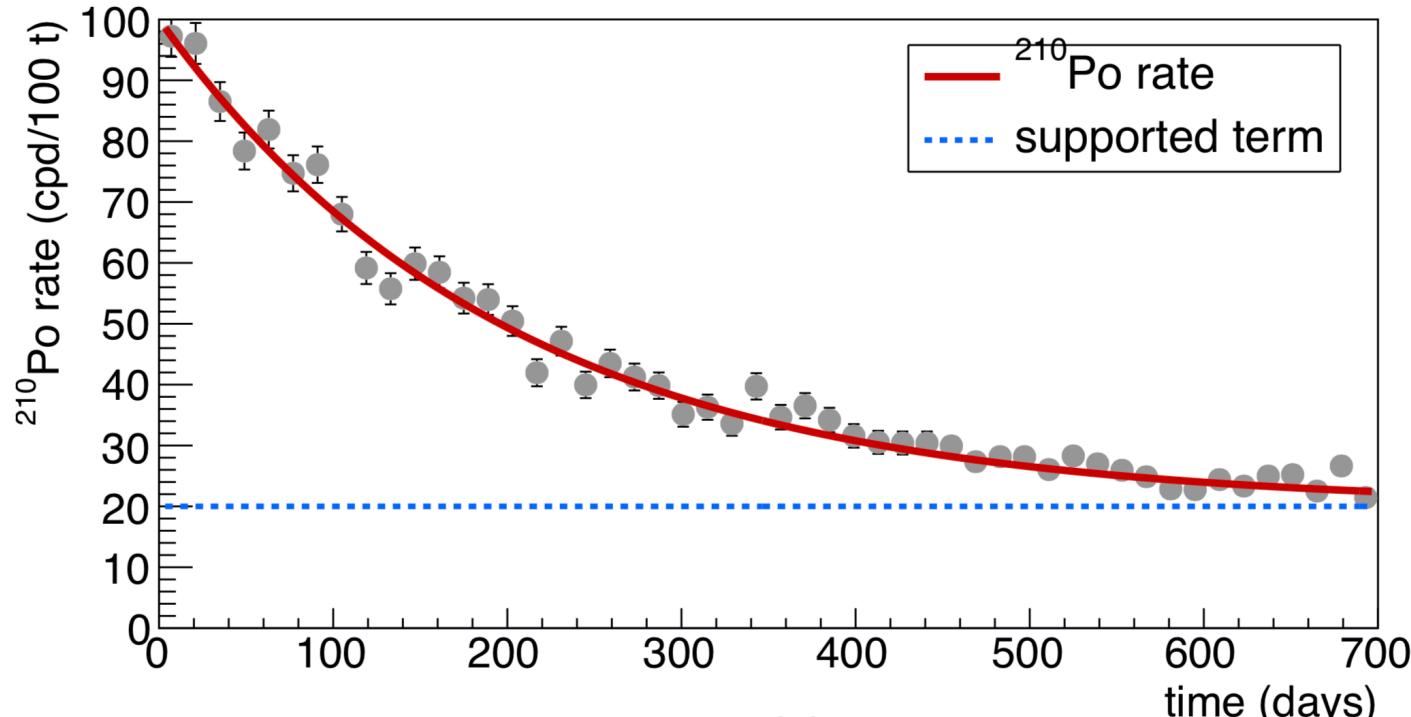
Indirect measurement of  $^{210}\text{Bi}$  by evolution of  $^{210}\text{Po}$  (Villante et al. 2011) provided  
 $^{210}\text{Bi} \rightarrow ^{210}\text{Po}$  only source of  $^{210}\text{Po}$



But, slow convection in the scintillator was bringing  $^{210}\text{Po}$  from the nylon vessel to the fiducial volume

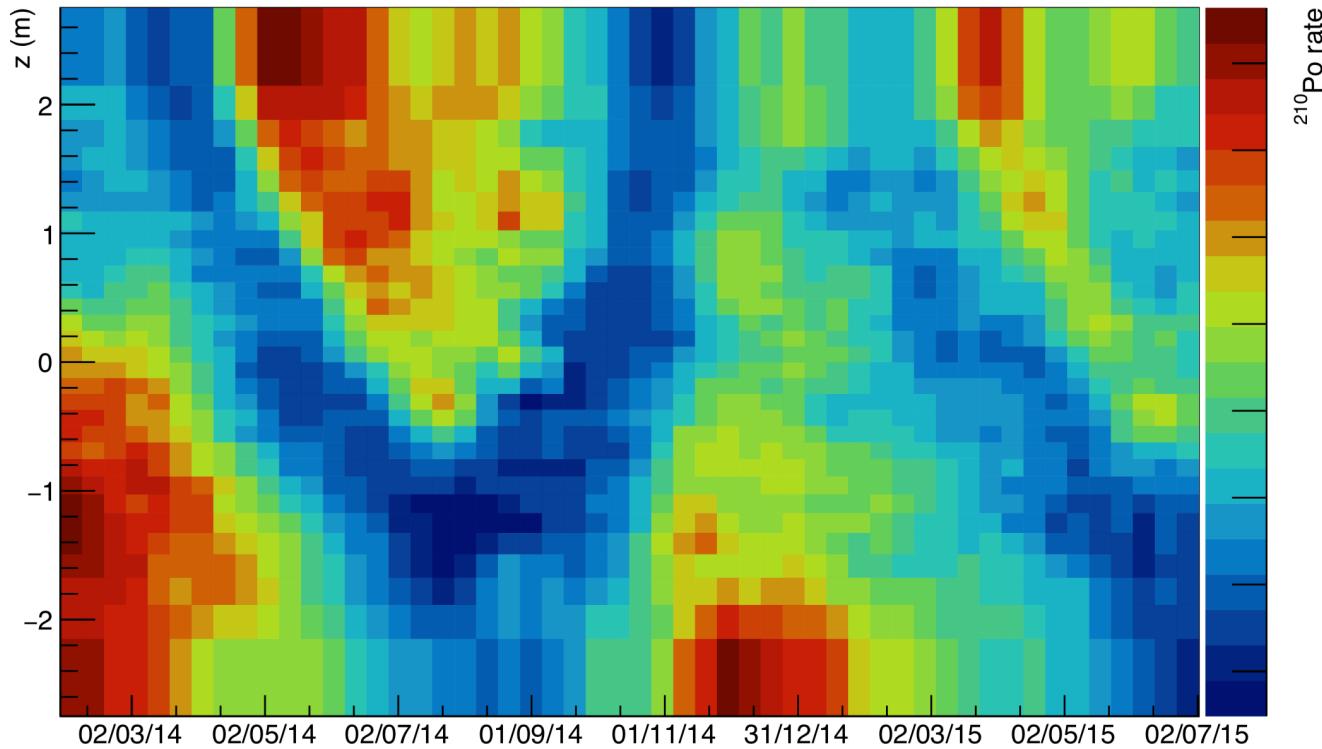
# CN-vs at Borexino

Po Time Evolution (example)



Borexino coll.

# CN-vs at Borexino



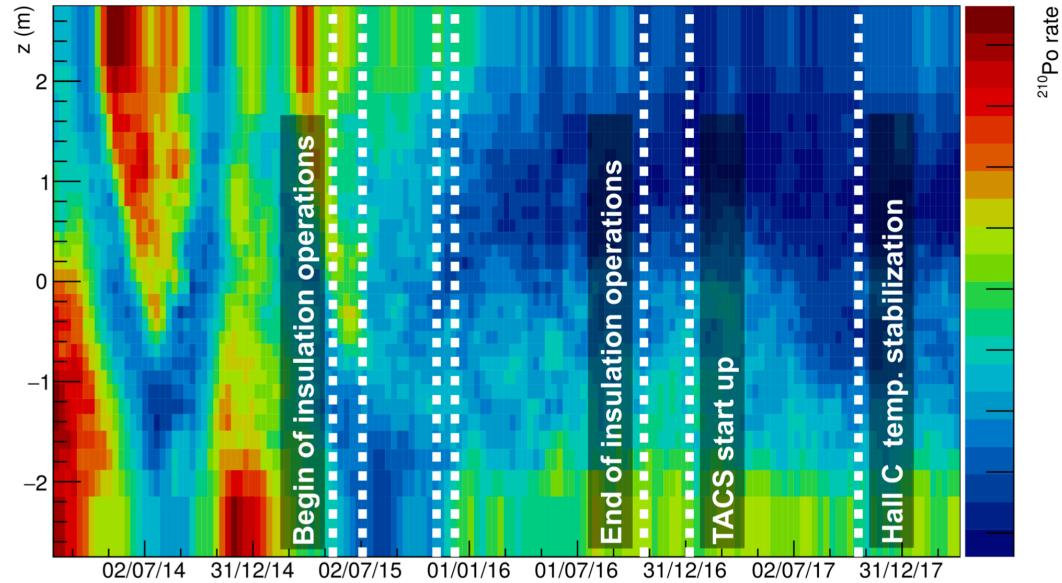
Guffanti 2018 (Borexino coll.) @ 5<sup>th</sup> International Solar Neutrino Conference

# CN-vs at Borexino

Thermal insulation



$^{210}\text{Po}$  rate evolution



after insulation

Guffanti 2018

# Summary

- The Sun shines by pp burning :  $1.03 \pm 0.08 L_{\odot}$  – all neutrino experiments
- The Sun shines by pp burning :  $1.01 \pm 0.10 L_{\odot}$  – Borexino experiment
- Open question: pp neutrinos measurement to 1% needed to test other energy sources in the Sun
- Open question: direct detection of CN fluxes
- Precise determination of solar opacity profile from vs (core opacity) and seismic data
- Solar abundance/model problem remains: opacity  $\leftrightarrow$  composition degeneracy  
Exporting abundance-opacity problem to other stars: systematic errors in ages by 10-15%
- First experimental opacity measurement @ solar conditions  
hints of higher Fe opacity at right place : 7%  
not enough :  $\sim 20\%$  needed
- Open question: are there other mechanisms of energy transport at work (e.g. ADM), modifying ‘effective’ radiative opacity?
- CN fluxes remain necessary and only way to break degeneracy  
Whatever the result, very important measurement  $\rightarrow$  core C+N abundance  
40% solar abundance problem, 15% chemical mixing processes in the Sun (surface/core difference)

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# SSM – neutrinos

No luminosity constraint – purely experimental result

$$\frac{L_{\text{pp-chain}}}{L_{\odot}} = 1.03^{+0.08}_{-0.07} \left[ {}^{+0.21}_{-0.18} \right] \quad \text{and} \quad \frac{L_{\text{CNO}}}{L_{\odot}} = 0.008^{+0.005}_{-0.004} \left[ {}^{+0.014}_{-0.007} \right].$$

$$\frac{L_{\odot}(\text{neutrino-inferred})}{L_{\odot}} = 1.04^{+0.07}_{-0.08} \left[ {}^{+0.20}_{-0.18} \right].$$

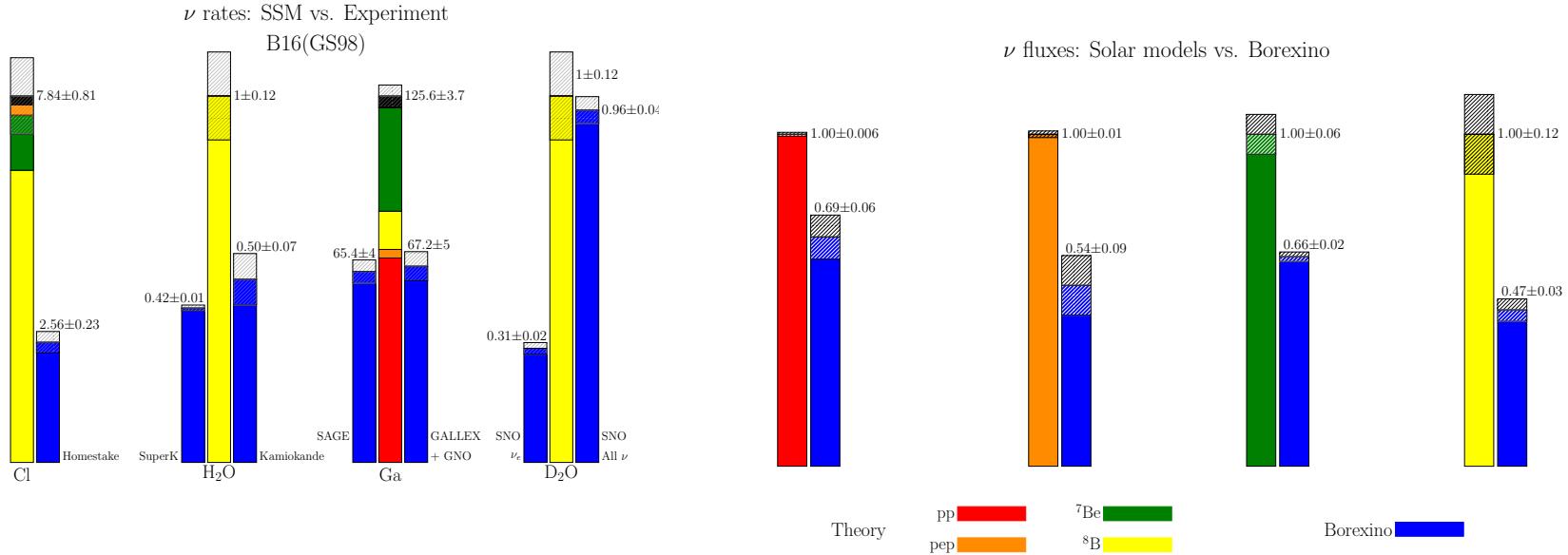
With luminosity constraint –  $L_{\odot} = L_{\text{nuc}}$

$$\frac{L_{\text{pp-chain}}}{L_{\odot}} = 0.991^{+0.005}_{-0.004} \left[ {}^{+0.008}_{-0.013} \right] \iff \frac{L_{\text{CNO}}}{L_{\odot}} = 0.009^{+0.004}_{-0.005} \left[ {}^{+0.013}_{-0.008} \right]$$

Global analysis with more recent data needed, e.g. Borexino – see Ianni's talk

Bergstrom et al. 2016

# Neutrinos: theory vs experiment



# Solar core temperature → solar models

Strong T dependence of ν-fluxes

$$\Phi(^8\text{B}) \propto T_c^{25} \longrightarrow \text{SuperK+SNO} \longrightarrow \delta T_c/T_c \approx 0.1\%$$

$$\Phi(^7\text{Be}) \propto T_c^{10} \longrightarrow \text{Borexino} \longrightarrow \delta T_c/T_c \approx 0.4\%$$

But this determines only precision – actual  $T_c$  determination requires solar models

Problems start here – models depend on inputs

Solar composition

Radiative opacities

Nuclear reaction rates

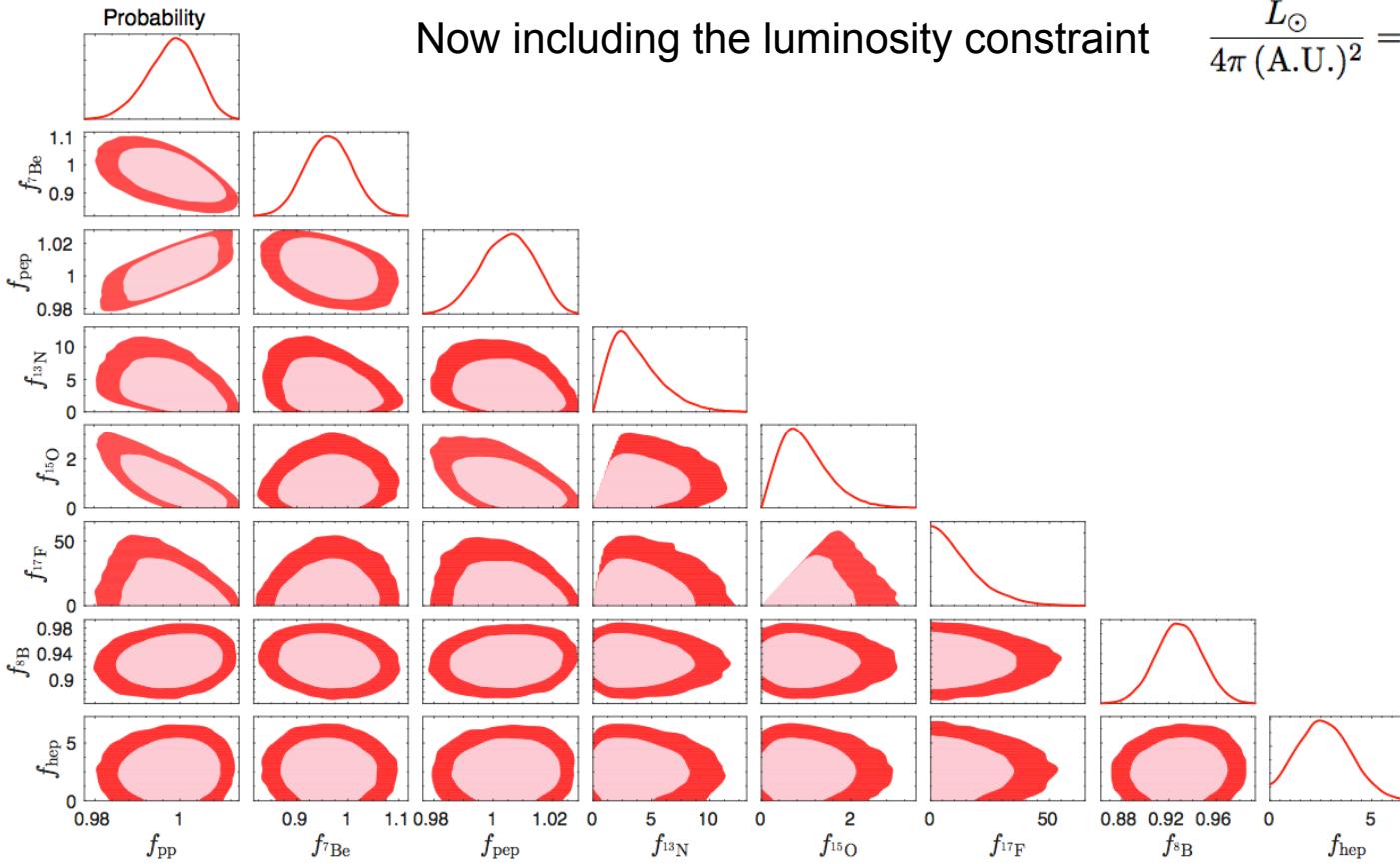
Chemical mixing (gravitational settling + other mixing processes)

Equation of state

# How does the Sun shine?

Now including the luminosity constraint

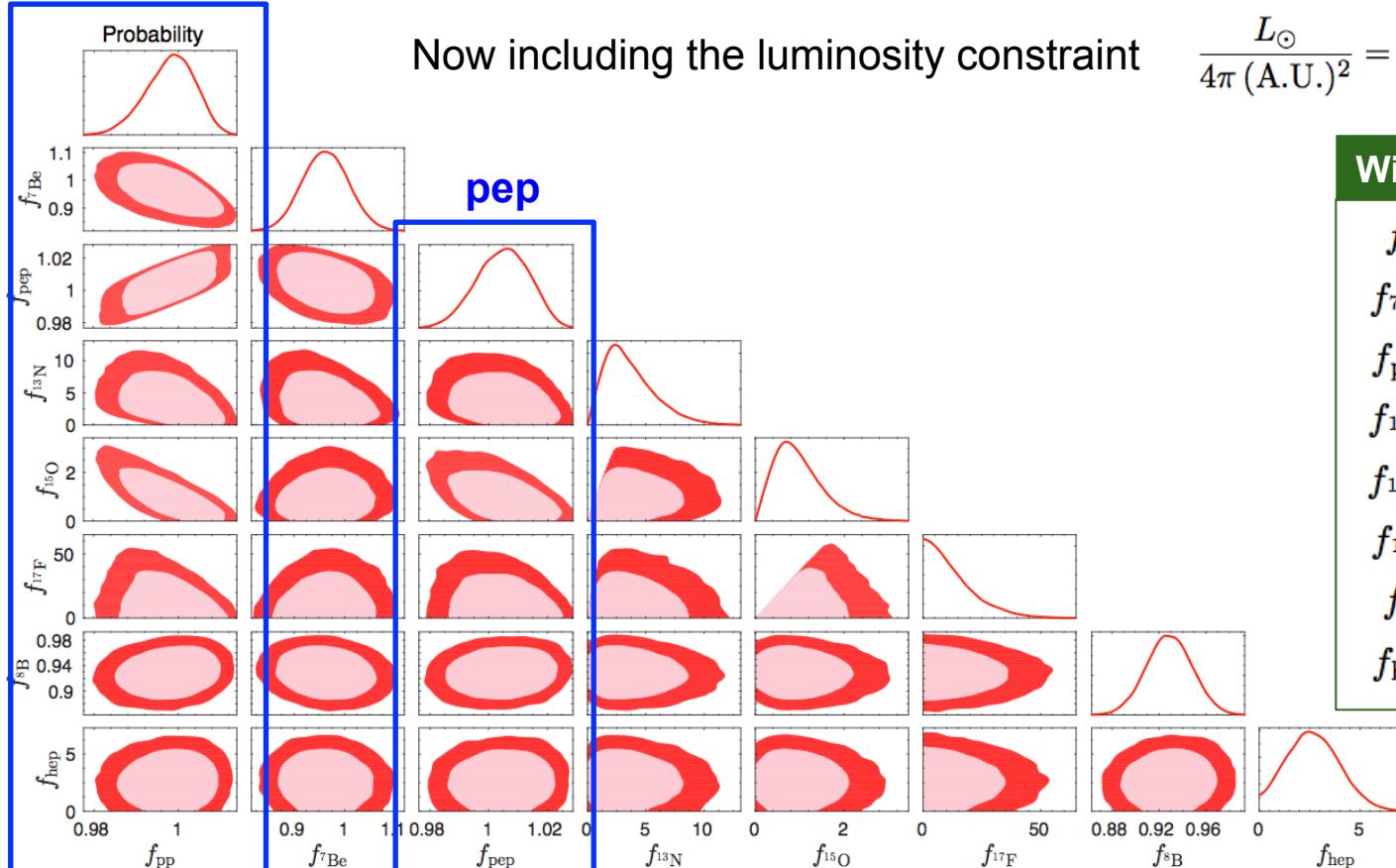
$$\frac{L_{\odot}}{4\pi \text{ (A.U.)}^2} = \sum_{i=1}^8 \alpha_i \Phi_i.$$



Bergstrom et al. 2016

# How does the Sun shine?

pp



Now including the luminosity constraint

$$\frac{L_{\odot}}{4\pi (\text{A.U.})^2} = \sum_{i=1}^8 \alpha_i \Phi_i.$$

**With luminosity constraint**

$$f_{\text{pp}} = 0.999^{+0.006}_{-0.005} [{}^{+0.012}_{-0.016}],$$

$$f_{^7\text{Be}} = 0.96^{+0.05}_{-0.04} [{}^{+0.12}_{-0.11}],$$

$$f_{\text{pep}} = 1.005 \pm 0.009 [{}^{+0.019}_{-0.024}],$$

$$f_{^{13}\text{N}} = 1.7^{+2.9}_{-1.0} [{}^{+8.4}_{-1.6}],$$

$$f_{^{15}\text{O}} = 0.6^{+0.6}_{-0.4} [{}^{+2.0}_{-0.6}],$$

$$f_{^{17}\text{F}} \leq 15 [46],$$

$$f_{^8\text{B}} = 0.92 \pm 0.02 [\pm 0.05],$$

$$f_{\text{hep}} = 2.4^{+1.5}_{-1.2} [\leq 5.9],$$

Bergstrom et al. 2016

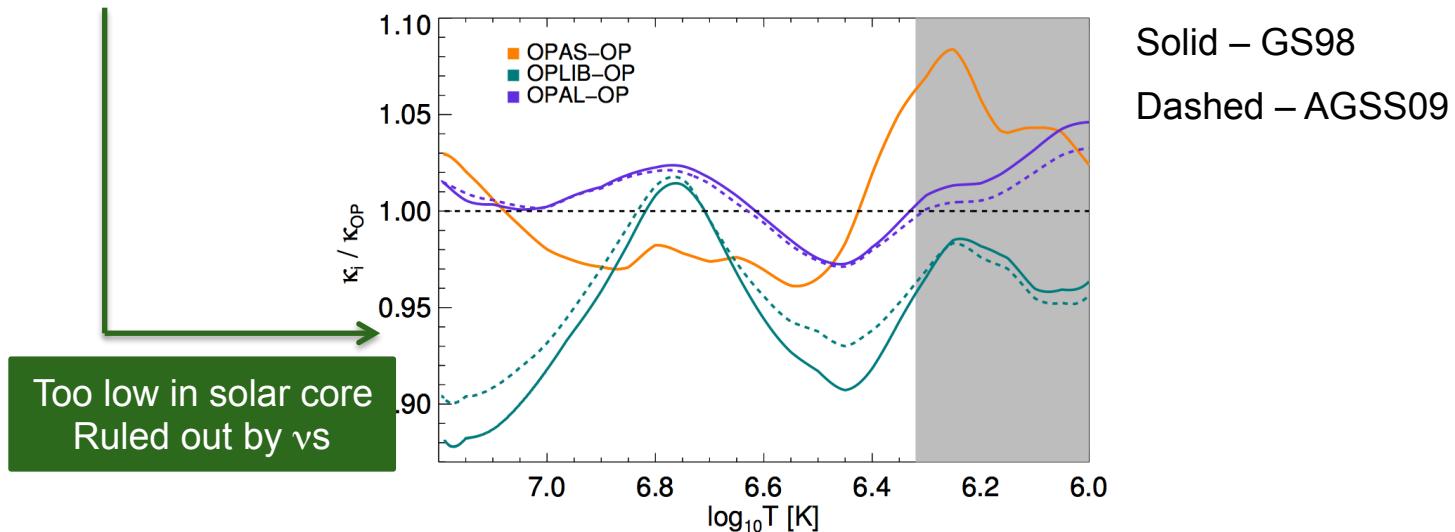
# Opacities – new calculations

## Old generation

- OPAL – Iglesias et al. 1996
- Opacity Project (OP) – Badnell et al. 2005

## New generation

- OPAS – Blancard et al. 2012 – now available Mondet et al. 2015 (only for AGSS09 composition)
- Los Alamos (OPLIB) – Colgan et al. 2016 – Most complete set from new generation



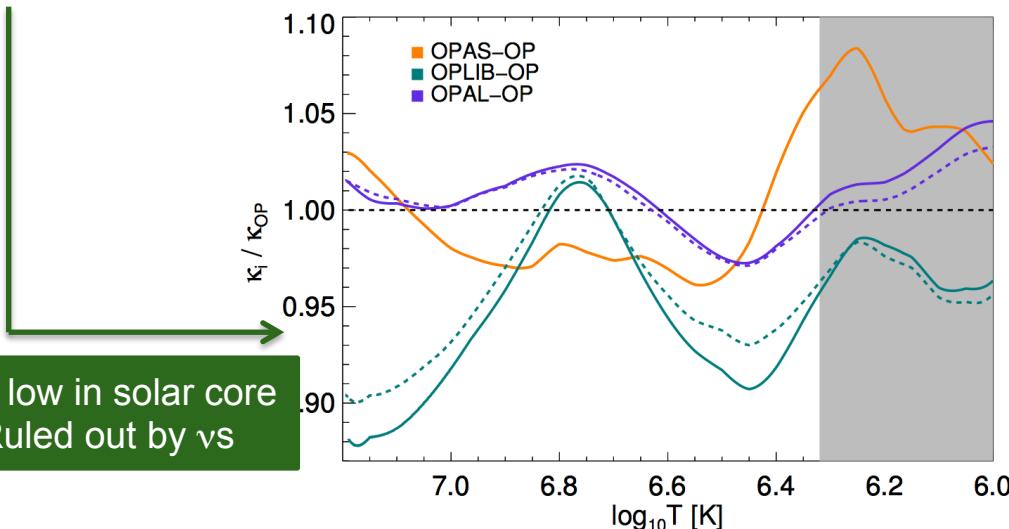
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Too low in solar core  
Ruled out by vs

Solid – GS98

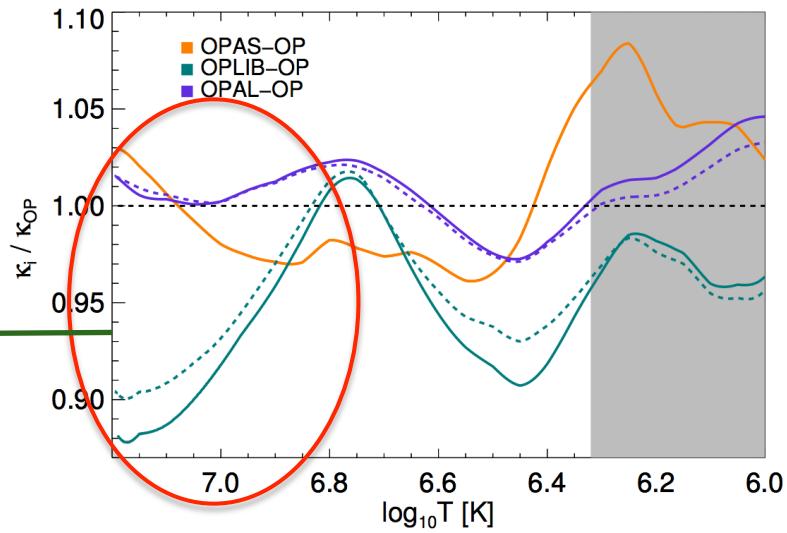
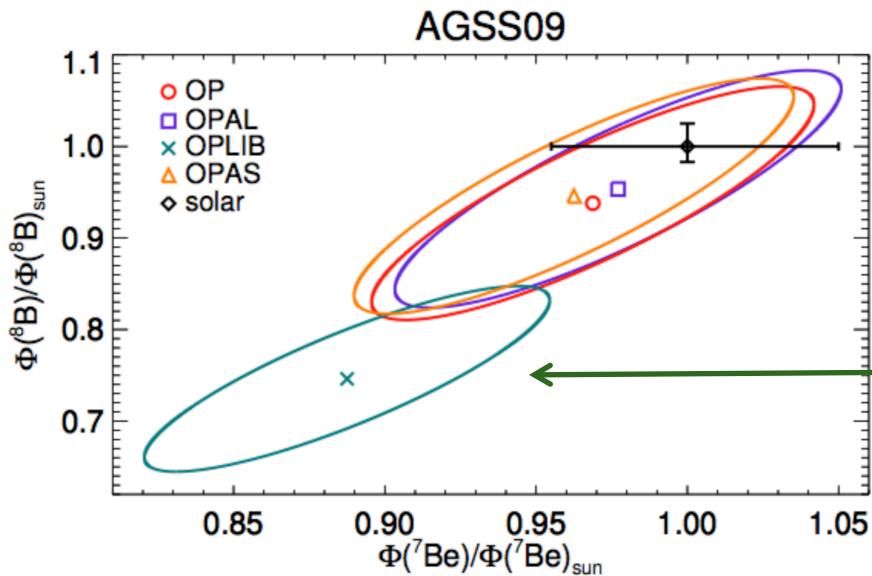
Dashed – AGSS09

**Not guaranteed that newer opacity models lead to higher opacity values**

**± 5% variations**

**Current situation unclear**

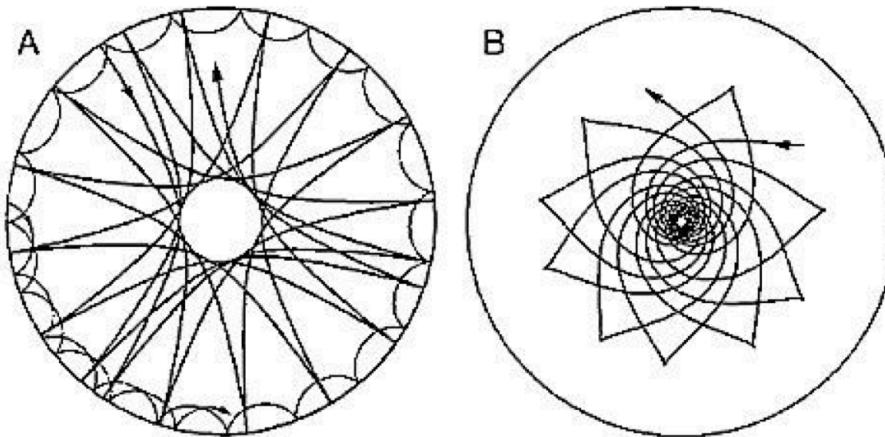
# SSM with new opacities



Solar vs rule out OPLIB opacities for low Z models

# g-modes detection (finally?)

g-modes probe inner regions – but strongly damped in the surface – tiny amplitudes & high background



direct searches for g-modes have failed (despite claims in Garcia et al. 2007)

Fossat et al. 2017 use new method: long term modulations in p-mode spectrum

Claim detections of more than 200 g-modes of angular degree  $l = 1, 2$

# g-modes detection (finally?)

Two important claims in Fossat et al. 2017

- 1) Asymptotic period spacings for  $\ell = 1, 2$

$$\Pi_\ell = \frac{2\pi^2}{\sqrt{\ell(\ell+1)}} \left[ \int_0^{R_{CZ}} N \frac{dr}{r} \right]^{-1}$$
$$N = g \left( \frac{1}{\Gamma_1} \frac{d \log p}{dr} - \frac{d \log \rho}{dr} \right)$$

Fossat et al.  $P_1 = 1443.1 \pm 0.5$ s -  $P_2 = 832.8 \pm 0.7$ s

GS98 SSMs:  $P_1 = 1525 - 1540$  s -  $P_2 = 880 - 890$  s

AGSS09 SSMs:  $P_1 = 1535 - 1560$  s -  $P_2 = 886 - 900$  s

- 2) Rotational splitting --> solar core rotation  $\sim x3$  faster than intermediate regions  
Maybe some impact for chemical mixing in the core – but in direction of lowering  $v$ -fluxes

# g-modes detection (finally?)

Two important claims in Fossat et al. 2017

- 1) Asymptotic period spacings for  $l=1, 2$

$$2\pi^2 \left[ \int_{R_{CZ}}^{R_{CZ}+dr} \right]^{-1} \left( 1 - d \log n - d \log a \right)$$

## From Appourchaux et al. 2010 review

and data-analysis perspectives – to give unambiguous detections of individual g modes.  
The review ends by concluding that, at the time of writing, there is indeed a consensus amongst the authors that *there is currently no undisputed detection of solar g modes.*

- 2)

Maybe some impact for chemical mixing in the core – but in direction of lowering  $\nu$ -fluxes

# SSM – B16 models

Flux	B16-GS98	B16-AGSS09met	Solar <sup>a</sup>	Chg.
$\Phi(pp)$	$5.98(1 \pm 0.006)$	$6.03(1 \pm 0.005)$	$5.97^{(1+0.006)}_{(1-0.005)}$	0.0
$\Phi(pep)$	$1.44(1 \pm 0.01)$	$1.46(1 \pm 0.009)$	$1.45^{(1+0.009)}_{(1-0.009)}$	0.0
$\Phi(hep)$	$7.98(1 \pm 0.30)$	$8.25(1 \pm 0.30)$	$19^{(1+0.63)}_{(1-0.47)}$	-0.7
$\Phi(^7Be)$	$4.93(1 \pm 0.06)$	$4.50(1 \pm 0.06)$	$4.80^{(1+0.050)}_{(1-0.046)}$	-1.4
$\phi(^8B)$	$5.46(1 \pm 0.12)$	$4.50(1 \pm 0.12)$	$5.16^{(1+0.025)}_{(1-0.017)}$	-2.2
$\phi(^{13}N)$	$2.78(1 \pm 0.15)$	$2.04(1 \pm 0.14)$	$\leq 13.7$	-6.1
$\phi(^{15}O)$	$2.05(1 \pm 0.17)$	$1.44(1 \pm 0.16)$	$\leq 2.8$	-8.1
$\phi(^{17}F)$	$5.29(1 \pm 0.20)$	$3.26(1 \pm 0.18)$	$\leq 85$	-4.2

New SSMs - changes in some nuclear rates  
(Vinyoles et al. 2017)

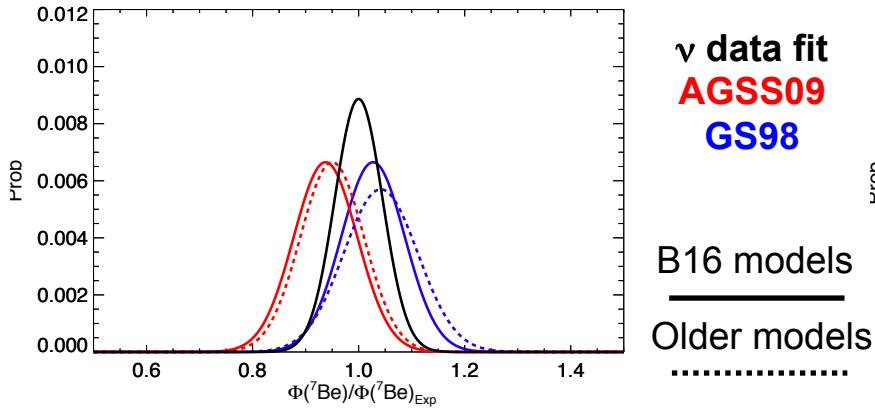
	$S(0)$	Uncert. %	$\Delta S(0)/S(0)$
$S_{11}$	$4.03 \cdot 10^{-25}$	1	+0.5%
$S_{17}$	$2.13 \cdot 10^{-5}$	4.7	+2.4%
$S_{114}$	$1.59 \cdot 10^{-3}$	7.5	-4.2%

Small changes  $^7\text{Be}$ - $^8\text{B}$  ( $S_{11}$ - $S_{17}$ )

Larger for  $^{13}\text{N}$ - $^{15}\text{O}$  ( $S_{11}$ - $S_{114}$ )

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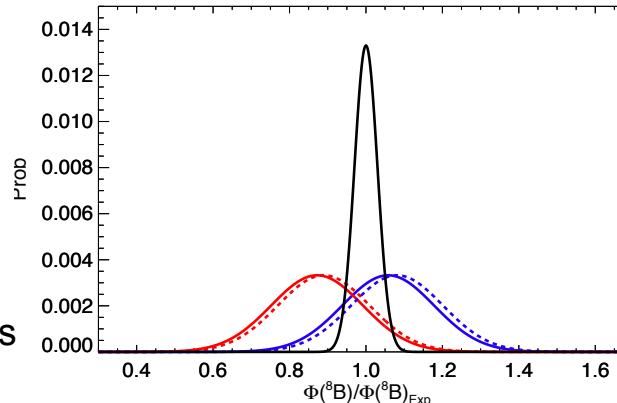


New SSMs - changes in some nuclear rates  
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Small changes  $^7\text{Be}-^8\text{B}$  ( $S_{11}-S_{17}$ )

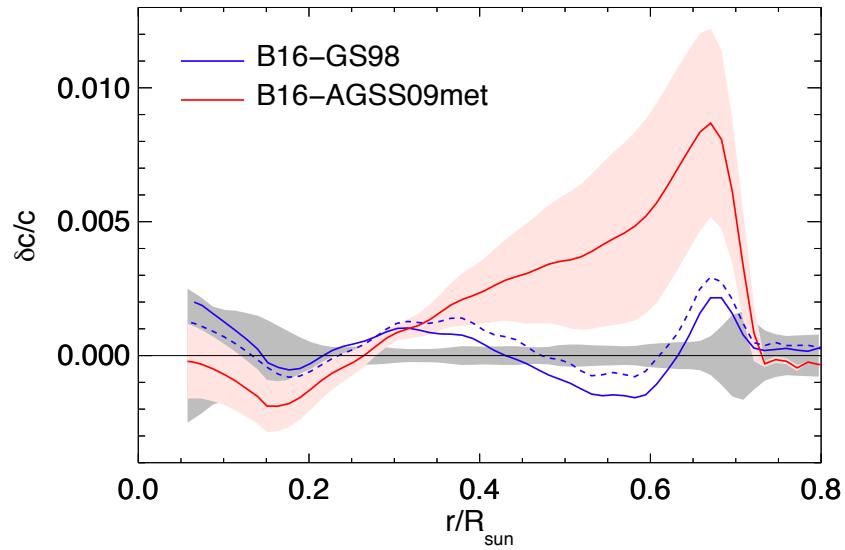
Larger for  $^{13}\text{N}-^{15}\text{O}$  ( $S_{11}-S_{114}$ )



Revision of global analysis including new Borexino data needed

# SSM – B16 models

## Small changes in helioseismic probes



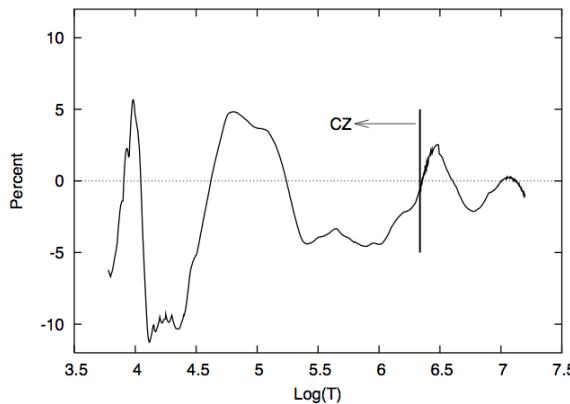
Qnt.	B16-GS98	B16-AGSS09met	Solar
$Y_S$	$0.2426 \pm 0.0059$	$0.2317 \pm 0.0059$	$0.2485 \pm 0.0035$
$R_{\text{CZ}}/R_{\odot}$	$0.7116 \pm 0.0048$	$0.7223 \pm 0.0053$	$0.713 \pm 0.001$
$\langle \delta c/c \rangle$	$0.0005 \pm 0.0004$	$0.0021 \pm 0.001$	—

# Opacities

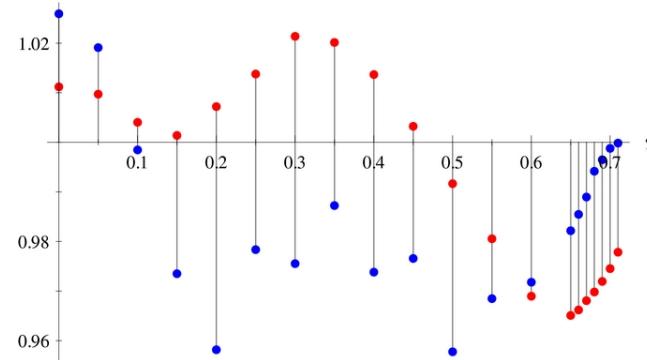
Helioseismic probes and pp  $\bar{\nu}$ s depend on “effective” opacity profiles: opacity models + composition details in F. Villante’s talk

Status of opacity models in 2014 @ “A special Borexino Event”

OP vs OPAL



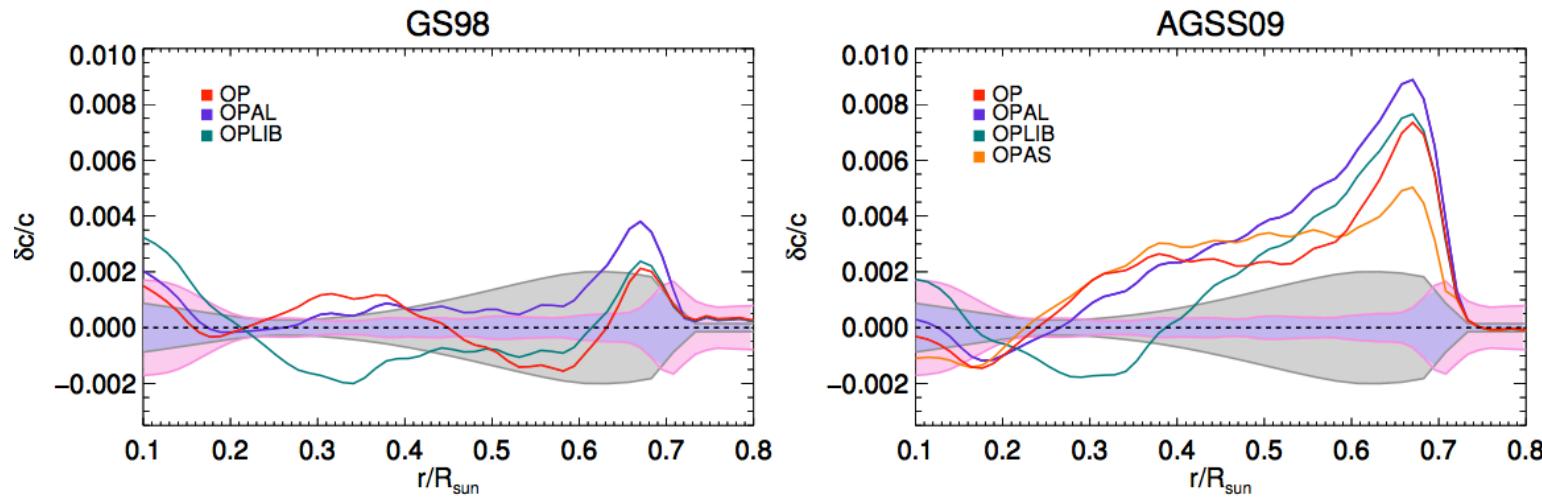
OPAS vs OP (blue)



Few percent differences in solar interiors

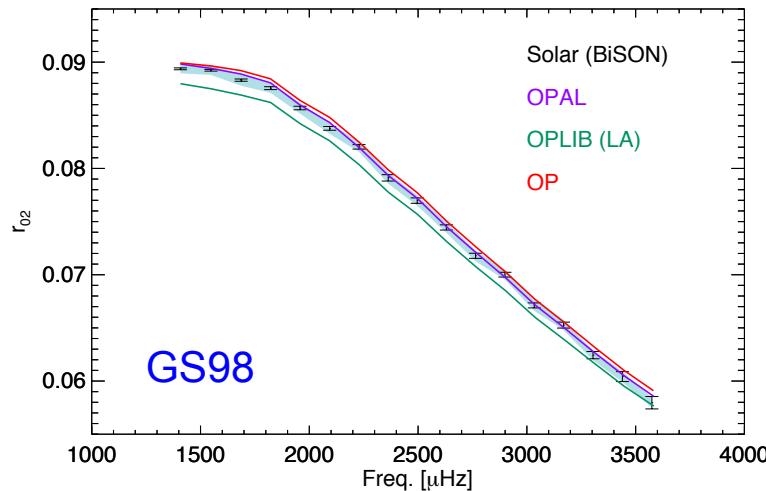
Only theoretical calculations available

# SSM with new opacities

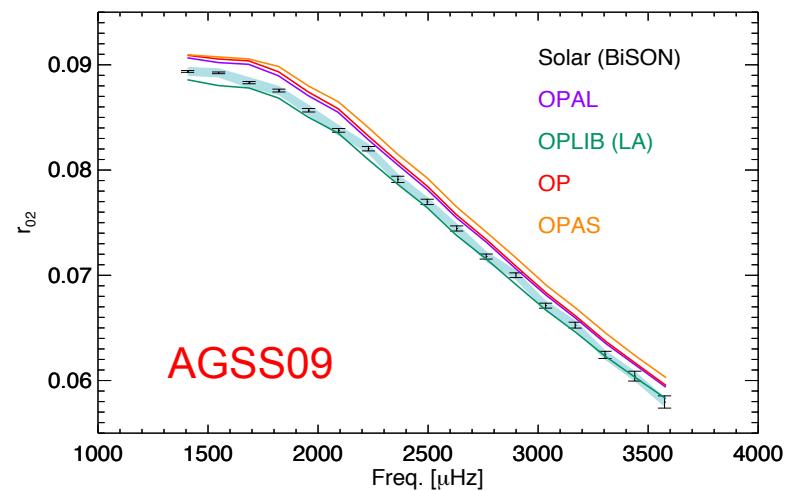


New opacities lead to some variations in sound speed profiles but nothing too dramatic

# SSM with new opacities



GS98

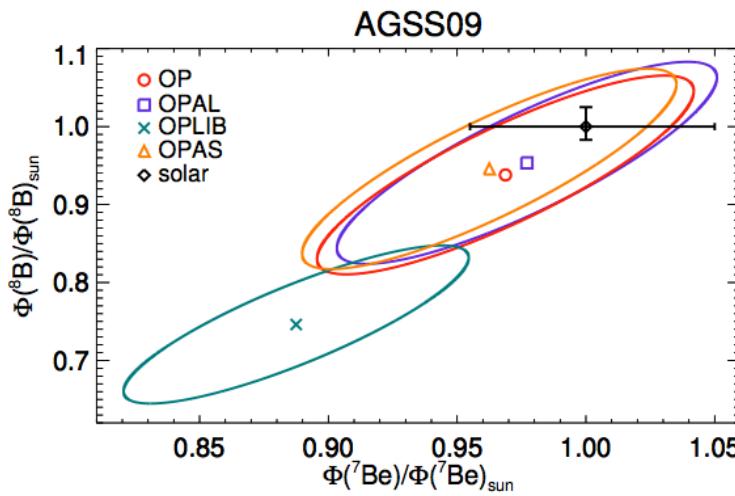
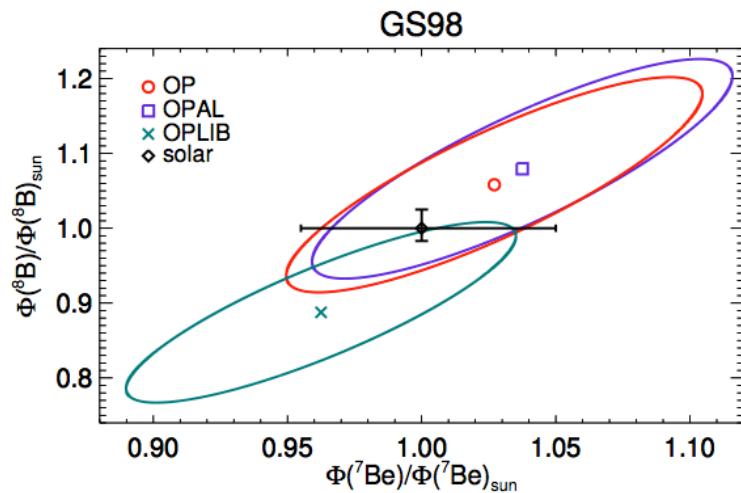


AGSS09

New OPLIB opacities lead to indecisive results for helioseismic probes

not all agree (disagree) with high(low) Z solar models

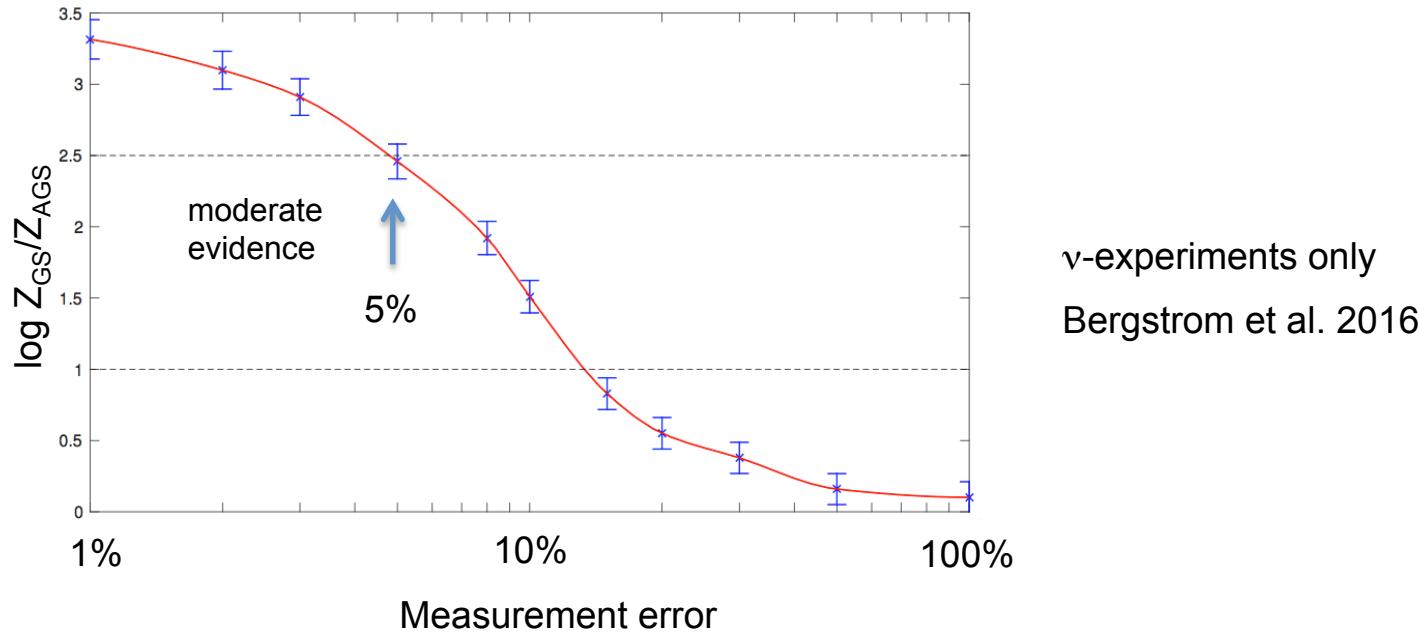
# SSM with new opacities



# SSM: the need for CN(O)

New opacity calculations do not alter state-of-the-art or complicate matters more

Most robust way to break the opacity  $\leftrightarrow$  composition degeneracy is through CNO  $n_s$



Discriminating power can improve if model information is added (Haxton et al. 2008)

# Extra slides

