

Fyzikálně-chemické základy nukleární magnetické rezonance

James Keeler: Understanding NMR Spectroscopy
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2002

NMR a elektromagnetické spektrum

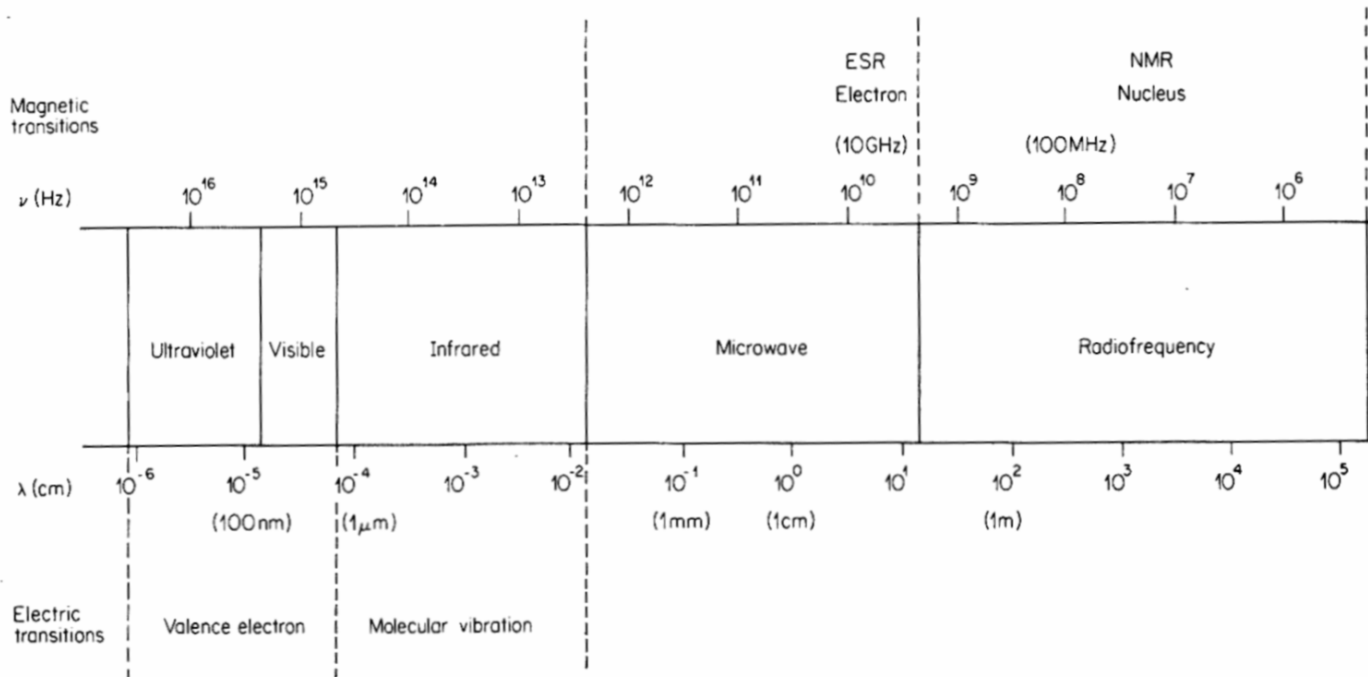
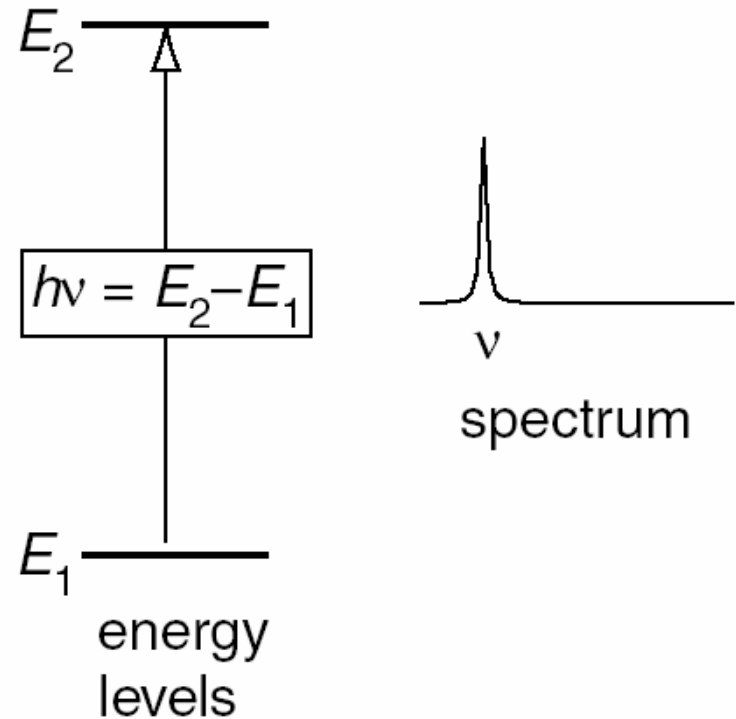


Figure 1.2 Spectral regions of the electromagnetic spectrum of interest in biological investigations

NMR a energiové hladiny

$$\Delta E = h\nu$$



h - Planckova konstanta = $6,626 \cdot 10^{-34}$ J.s

NMR a energiové hladiny

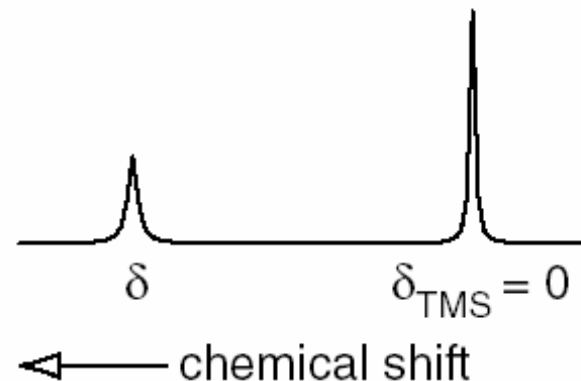
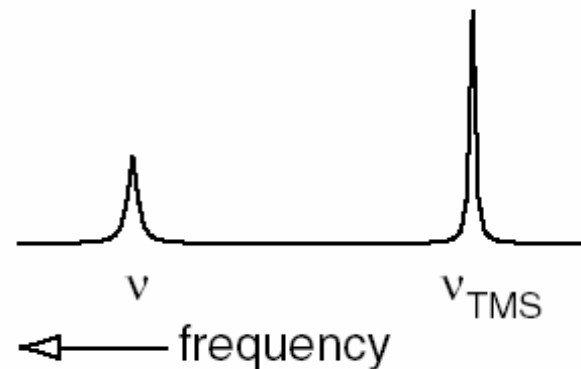
Chemický posun

$$\delta = \frac{\nu - \nu_{\text{TMS}}}{\nu_{\text{TMS}}}$$

$$\delta_{\text{ppm}} = 10^6 \times \frac{\nu - \nu_{\text{TMS}}}{\nu_{\text{TMS}}}$$

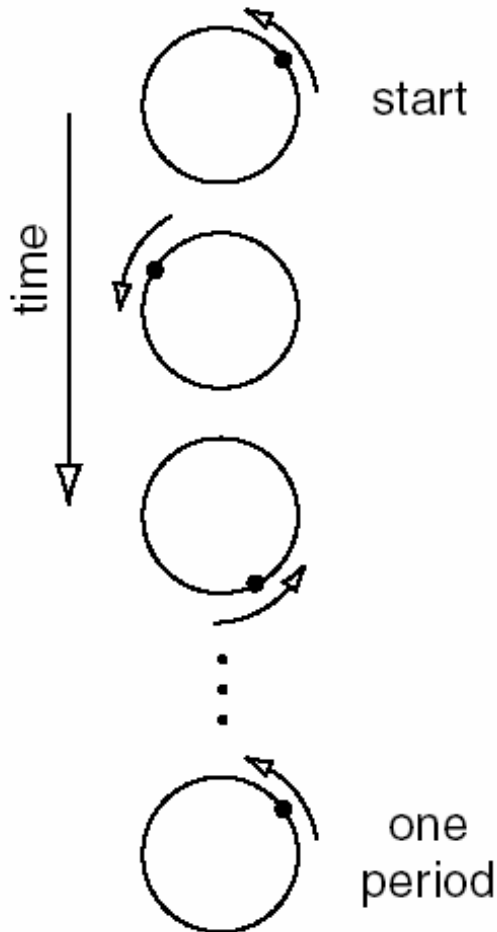
$$\begin{aligned}(\delta_2 - \delta_1) &= 10^6 \times \frac{\nu_2 - \nu_{\text{TMS}}}{\nu_{\text{TMS}}} - 10^6 \times \frac{\nu_1 - \nu_{\text{TMS}}}{\nu_{\text{TMS}}} \\ &= 10^6 \times \frac{\nu_2 - \nu_1}{\nu_{\text{TMS}}}\end{aligned}$$

$$(\nu_2 - \nu_1) = 10^{-6} \times \nu_{\text{TMS}} \times (\delta_2 - \delta_1)$$



NMR a energiové hladiny

Rotační frekvence a energie



$$\nu = \frac{1}{\tau}$$

$$\omega = \frac{2\pi}{\tau}$$

$$\nu = \frac{\omega}{2\pi} \quad \text{or} \quad \omega = 2\pi\nu$$

$$E = h \frac{\omega}{2\pi}$$
$$= \hbar\omega$$

NMR a energetické hladiny

Nukleární spin a spinové stavy

$I = k * \frac{1}{2}$, k je celé číslo 0, 1, 2 spinové kvantové číslo

$m = -I, -I+1, -I+2, \dots, I-2, I-1, I$ magnetické kvantové číslo

pro $I = \frac{1}{2}$ $m = \frac{1}{2}$ α stav
 $m = -\frac{1}{2}$ β stav

Pro 2 spiny s $I = \frac{1}{2}$ existují 4 možné kombinace stavů $\alpha\alpha$, $\alpha\beta$, $\beta\alpha$ a $\beta\beta$

Pro 3 spiny s $I = \frac{1}{2}$ existuje 8 možností kombinace stavů

$\alpha\alpha\alpha$, $\alpha\alpha\beta$, $\alpha\beta\alpha$, $\beta\alpha\alpha$, $\alpha\beta\beta$, $\beta\alpha\beta$, $\beta\beta\alpha$, a $\beta\beta\beta$



NMR a energiové hladiny

$$E_m = m * \nu_{0,1}$$

pořadové číslo spinu

ν_0 - Larmorova frekvence

$$E_\alpha = + 1/2 * \nu_{0,1} \quad E_\beta = - 1/2 * \nu_{0,1}$$

$$\nu_{0,1} = -\frac{1}{2\pi} \gamma_1 (1 + \delta_1) B_0$$

$$\omega_0 = -\gamma (1 + \delta) B_0$$

γ - magnetogyrická konstanta (poměr) [rad. s⁻¹.T⁻¹]

pro ¹H $\gamma = +2,67 \times 10^8 \text{ rad. s}^{-1}.\text{T}^{-1} = 42\,494\,369 \text{ s}^{-1}.\text{T}^{-1}$

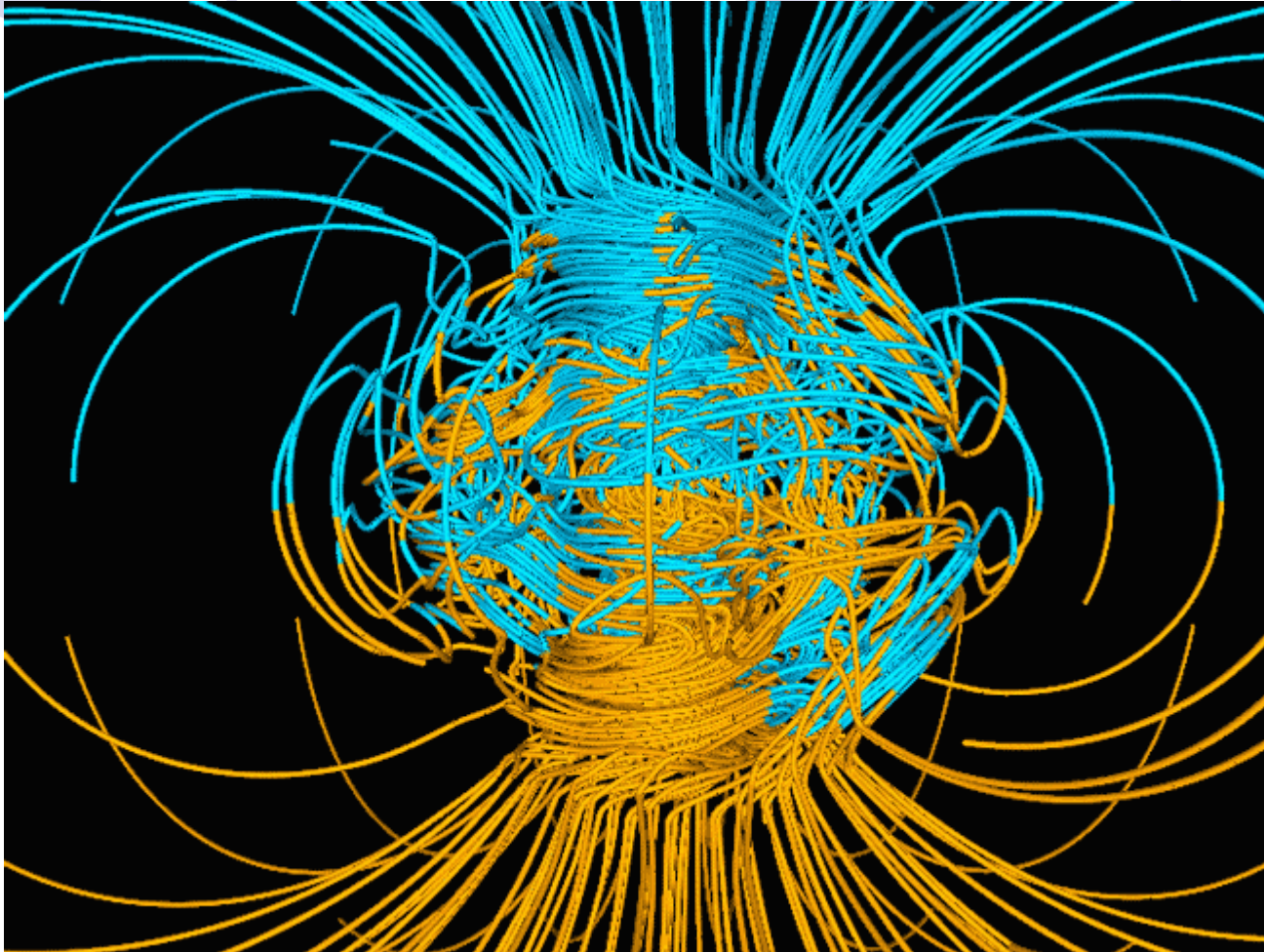
Při $B_0 = 4.7 \text{ T}$

$$\nu_0 = -200 \times 10^6 \text{ Hz}$$

a

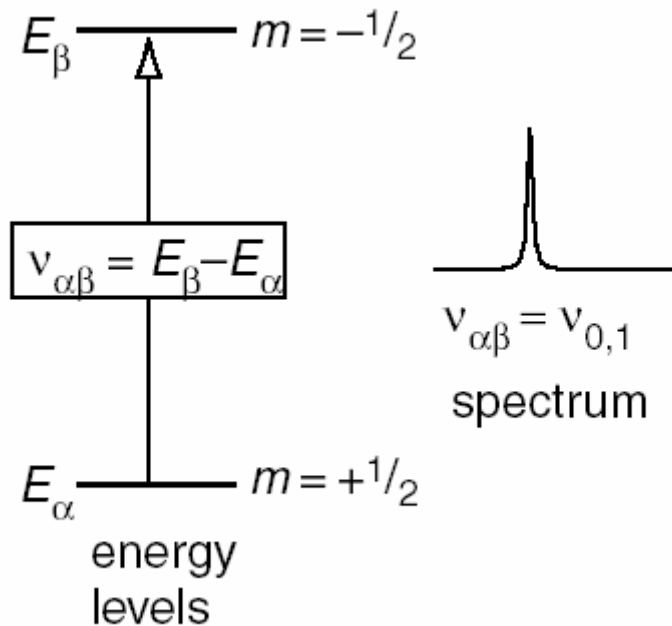
$$\omega_0 = -1,225 \times 10^9 \text{ rad.s}^{-1}$$

Zemské magnetické pole $B_0 = 40 \text{ mT}$, $\nu_0 = -1.7 \text{ MHz}$



NMR a energiové hladiny

Spektrum

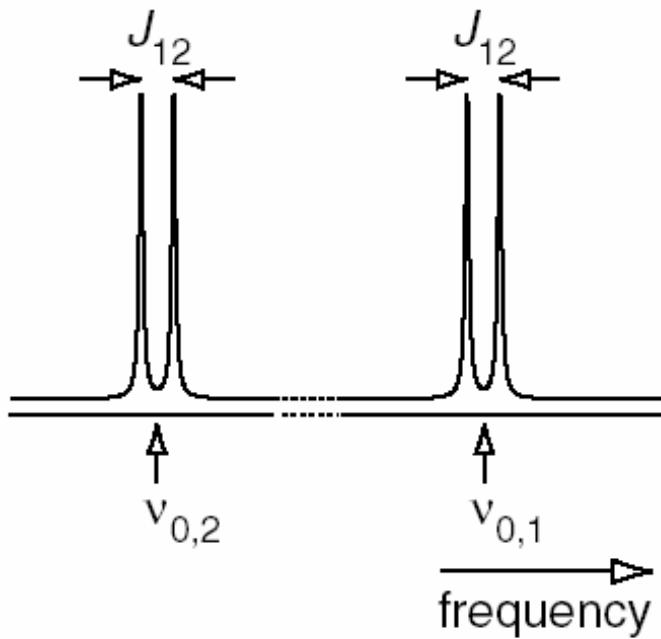


$$\Delta m = m(\text{initial state}) - m(\text{final state}) = \pm 1.$$

$$\begin{aligned} \nu_{\alpha\beta} &= E_\beta - E_\alpha \\ &= -\frac{1}{2}\nu_{0,1} - \left(+\frac{1}{2}\nu_{0,1}\right) \\ &= -\nu_{0,1}. \end{aligned}$$

NMR a energiové hladiny

Dva spiny



$$\nu_{0,1} = -\frac{1}{2\pi} \gamma_1 (1 + \delta_1) B_0$$

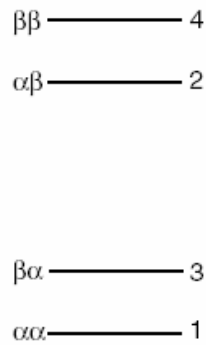
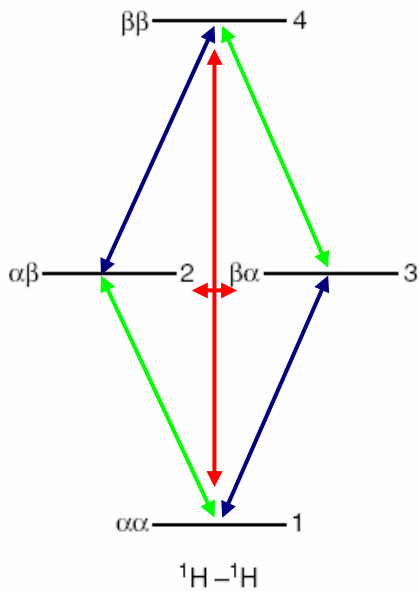
$$\nu_{0,2} = -\frac{1}{2\pi} \gamma_2 (1 + \delta_2) B_0.$$

| number | spin states | energy |
|--------|----------------|--|
| 1 | $\alpha\alpha$ | $+\frac{1}{2}\nu_{0,1} + \frac{1}{2}\nu_{0,2} + \frac{1}{4}J_{12}$ |
| 2 | $\alpha\beta$ | $+\frac{1}{2}\nu_{0,1} - \frac{1}{2}\nu_{0,2} - \frac{1}{4}J_{12}$ |
| 3 | $\beta\alpha$ | $-\frac{1}{2}\nu_{0,1} + \frac{1}{2}\nu_{0,2} - \frac{1}{4}J_{12}$ |
| 4 | $\beta\beta$ | $-\frac{1}{2}\nu_{0,1} - \frac{1}{2}\nu_{0,2} + \frac{1}{4}J_{12}$ |

$$E_{m_1 m_2} = m_1 \nu_{0,1} + m_2 \nu_{0,2} + m_1 m_2 J_{12}$$

NMR a energiové hladiny

Dva spiny



$$M = m_1 + m_2.$$

| number | spin states | M |
|--------|----------------|-----|
| 1 | $\alpha\alpha$ | 1 |
| 2 | $\alpha\beta$ | 0 |
| 3 | $\beta\alpha$ | 0 |
| 4 | $\beta\beta$ | -1 |

Výběrové pravidlo $\Delta M = \pm 1$

Dovolené přechody mezi

1-3 a 2-4, 1-2 a 3-4,

NMR a energiové hladiny

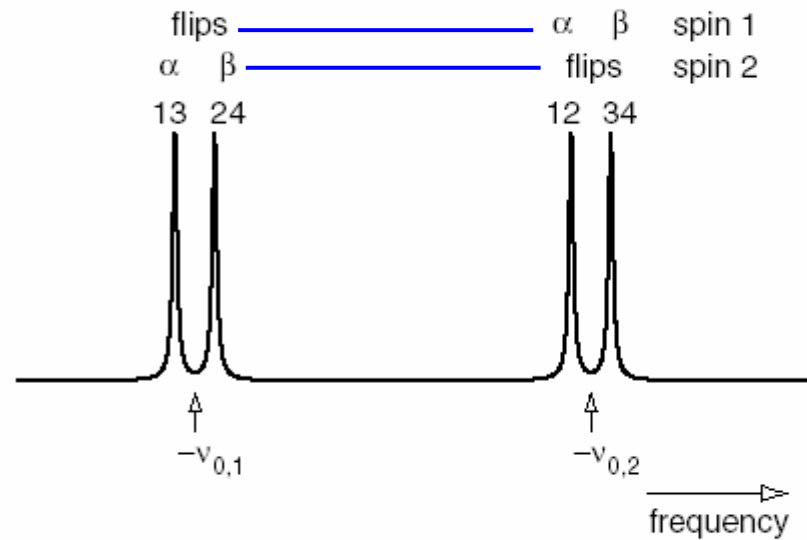
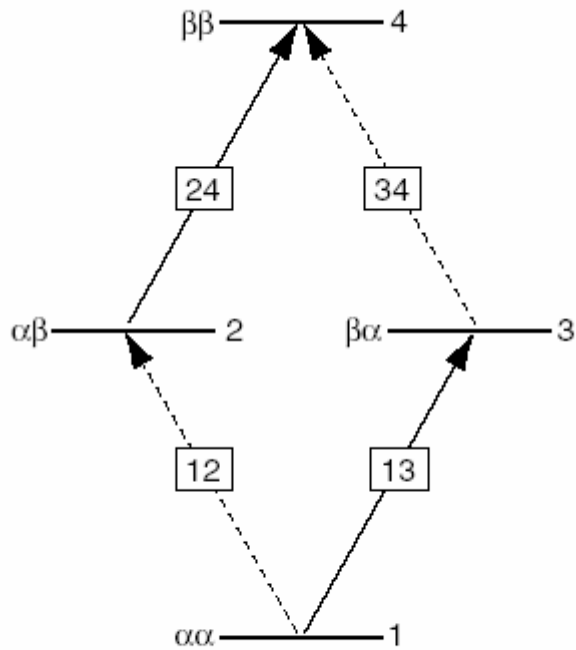
Dva spiny

$$\begin{aligned}\nu_{12} &= E_2 - E_1 \\ &= +\frac{1}{2}\nu_{0,1} - \frac{1}{2}\nu_{0,2} - \frac{1}{4}J_{12} - \left(\frac{1}{2}\nu_{0,1} + \frac{1}{2}\nu_{0,2} + \frac{1}{4}J_{12}\right) \\ &= -\nu_{0,2} - \frac{1}{2}J_{12}.\end{aligned}$$

| transition | spin states | frequency |
|-------------------|--|----------------------------------|
| $1 \rightarrow 2$ | $\alpha\alpha \rightarrow \alpha\beta$ | $-\nu_{0,2} - \frac{1}{2}J_{12}$ |
| $3 \rightarrow 4$ | $\beta\alpha \rightarrow \beta\beta$ | $-\nu_{0,2} + \frac{1}{2}J_{12}$ |
| $1 \rightarrow 3$ | $\alpha\alpha \rightarrow \beta\alpha$ | $-\nu_{0,1} - \frac{1}{2}J_{12}$ |
| $2 \rightarrow 4$ | $\alpha\beta \rightarrow \beta\beta$ | $-\nu_{0,1} + \frac{1}{2}J_{12}$ |

NMR a energiové hladiny

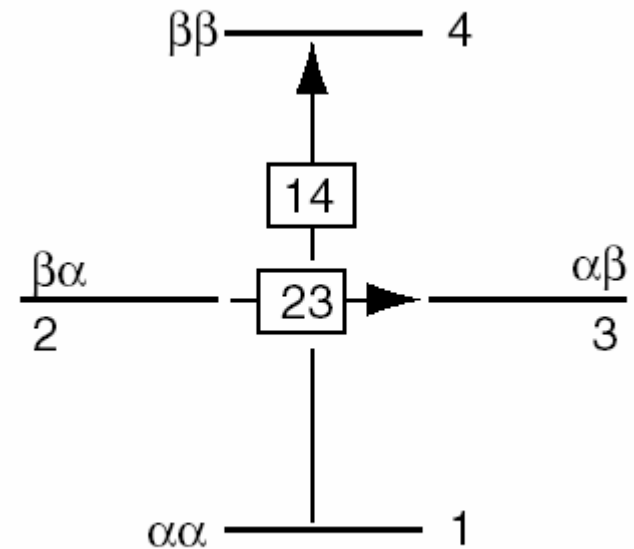
Dva spiny – jedno-kvantové přechody



NMR a energiové hladiny

Dva spiny – více-kvantové přechody

| number | spin states | M |
|--------|----------------|-----|
| 1 | $\alpha\alpha$ | 1 |
| 2 | $\alpha\beta$ | 0 |
| 3 | $\beta\alpha$ | 0 |
| 4 | $\beta\beta$ | -1 |



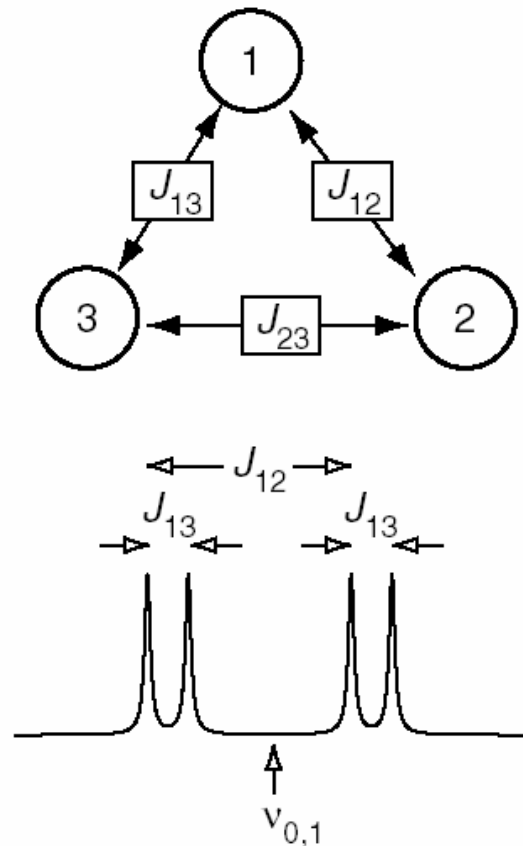
14 – dvou-kvantový přechod
23 – nul-kvantový přechod

NMR a energiové hladiny

Tři spiny

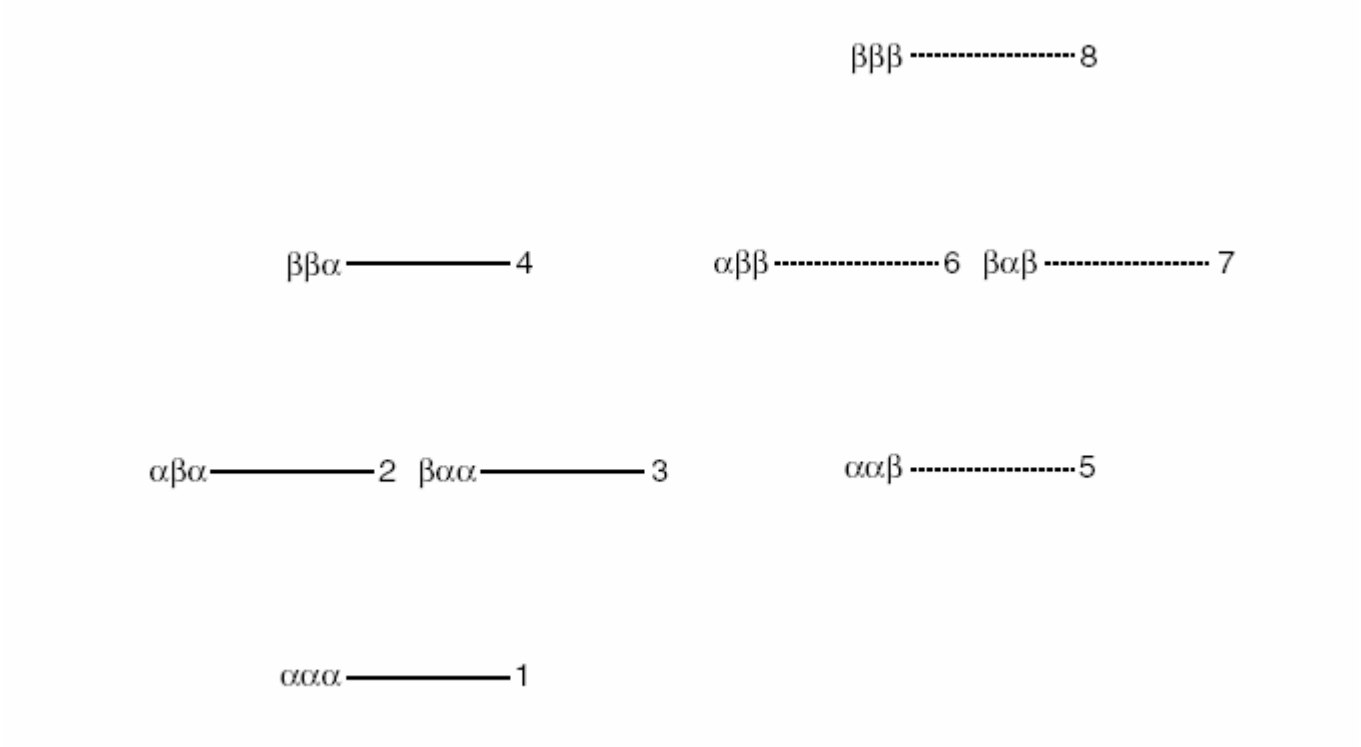
$$E_{m_1 m_2 m_3} = m_1 \nu_{0,1} + m_2 \nu_{0,2} + m_3 \nu_{0,3} + m_1 m_2 J_{12} + m_1 m_3 J_{13} + m_2 m_3 J_{23}$$

| number | spin states | M | energy |
|--------|----------------------|----------------|---|
| 1 | $\alpha\alpha\alpha$ | $\frac{3}{2}$ | $+\frac{1}{2}\nu_{0,1} + \frac{1}{2}\nu_{0,2} + \frac{1}{2}\nu_{0,3} + \frac{1}{4}J_{12} + \frac{1}{4}J_{13} + \frac{1}{4}J_{23}$ |
| 2 | $\alpha\beta\alpha$ | $\frac{1}{2}$ | $+\frac{1}{2}\nu_{0,1} - \frac{1}{2}\nu_{0,2} + \frac{1}{2}\nu_{0,3} - \frac{1}{4}J_{12} + \frac{1}{4}J_{13} - \frac{1}{4}J_{23}$ |
| 3 | $\beta\alpha\alpha$ | $\frac{1}{2}$ | $-\frac{1}{2}\nu_{0,1} + \frac{1}{2}\nu_{0,2} + \frac{1}{2}\nu_{0,3} - \frac{1}{4}J_{12} - \frac{1}{4}J_{13} + \frac{1}{4}J_{23}$ |
| 4 | $\beta\beta\alpha$ | $-\frac{1}{2}$ | $-\frac{1}{2}\nu_{0,1} - \frac{1}{2}\nu_{0,2} + \frac{1}{2}\nu_{0,3} + \frac{1}{4}J_{12} - \frac{1}{4}J_{13} - \frac{1}{4}J_{23}$ |
| 5 | $\alpha\alpha\beta$ | $\frac{1}{2}$ | $+\frac{1}{2}\nu_{0,1} + \frac{1}{2}\nu_{0,2} - \frac{1}{2}\nu_{0,3} + \frac{1}{4}J_{12} - \frac{1}{4}J_{13} - \frac{1}{4}J_{23}$ |
| 6 | $\alpha\beta\beta$ | $-\frac{1}{2}$ | $+\frac{1}{2}\nu_{0,1} - \frac{1}{2}\nu_{0,2} - \frac{1}{2}\nu_{0,3} - \frac{1}{4}J_{12} - \frac{1}{4}J_{13} + \frac{1}{4}J_{23}$ |
| 7 | $\beta\alpha\beta$ | $-\frac{1}{2}$ | $-\frac{1}{2}\nu_{0,1} + \frac{1}{2}\nu_{0,2} - \frac{1}{2}\nu_{0,3} - \frac{1}{4}J_{12} + \frac{1}{4}J_{13} - \frac{1}{4}J_{23}$ |
| 8 | $\beta\beta\beta$ | $-\frac{3}{2}$ | $-\frac{1}{2}\nu_{0,1} - \frac{1}{2}\nu_{0,2} - \frac{1}{2}\nu_{0,3} + \frac{1}{4}J_{12} + \frac{1}{4}J_{13} + \frac{1}{4}J_{23}$ |



NMR a energiové hladiny

Tři spiny



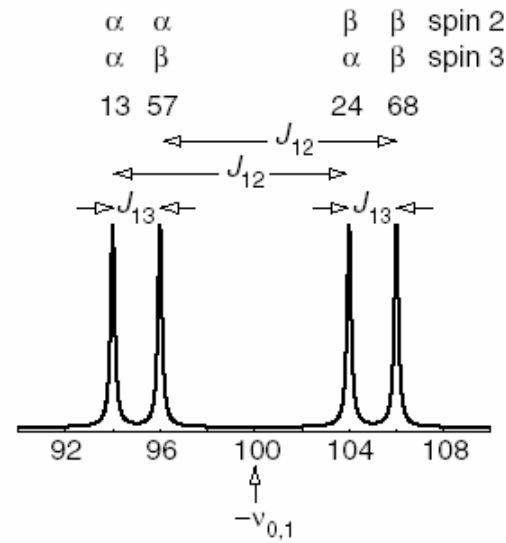
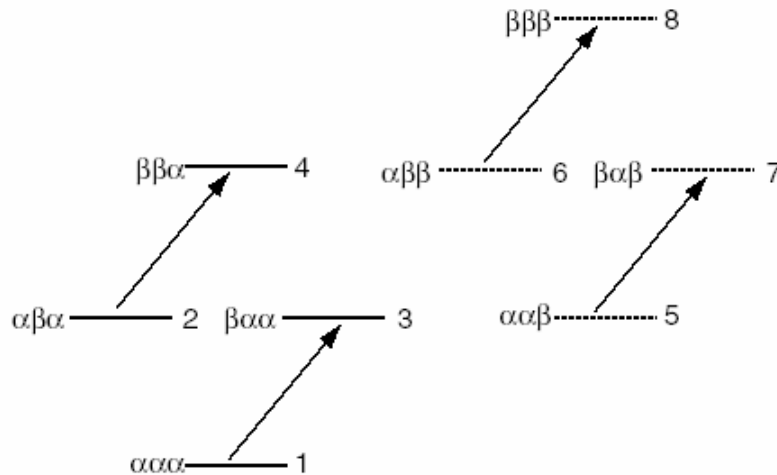
NMR a energiové hladiny

Tři spiny

| transition | state of spin 2 | state of spin 3 | frequency |
|------------|-----------------|-----------------|--|
| 1-3 | α | α | $-\nu_{0,1} - \frac{1}{2}J_{12} - \frac{1}{2}J_{13}$ |
| 2-4 | β | α | $-\nu_{0,1} + \frac{1}{2}J_{12} - \frac{1}{2}J_{13}$ |
| 5-7 | α | β | $-\nu_{0,1} - \frac{1}{2}J_{12} + \frac{1}{2}J_{13}$ |
| 6-8 | β | β | $-\nu_{0,1} + \frac{1}{2}J_{12} + \frac{1}{2}J_{13}$ |

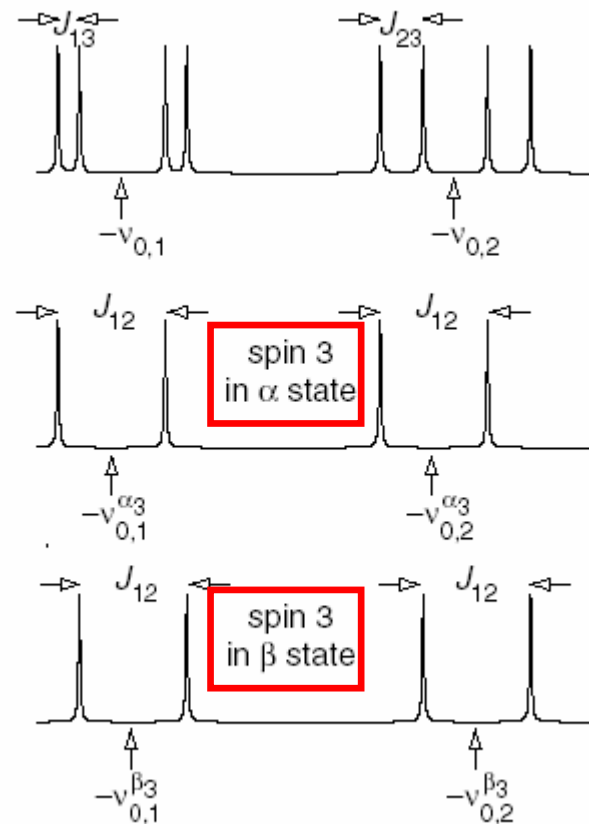
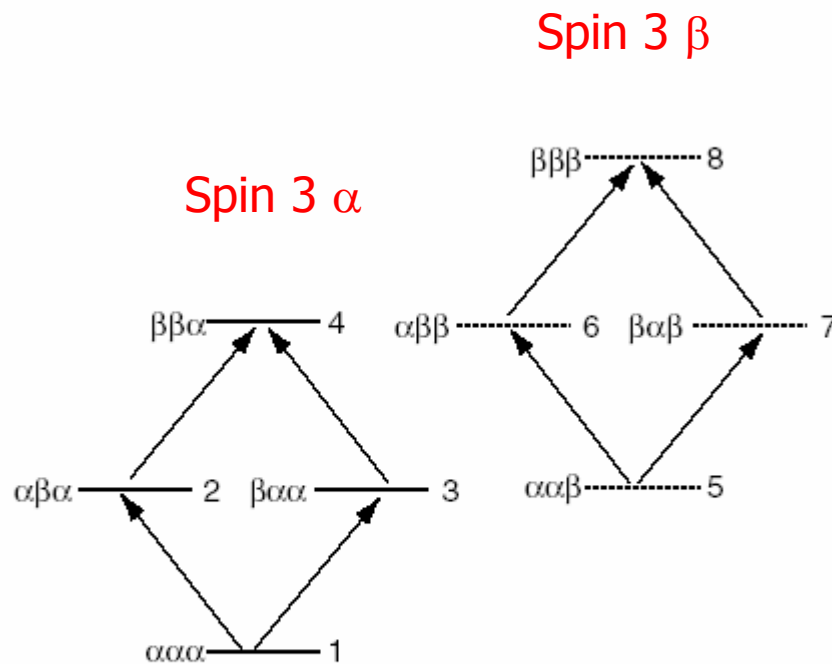
Výběrové pravidlo $\Delta M = \pm 1$

ale jen jeden spin může změnit svůj stav



NMR a energiové hladiny

Tři spiny - subspektra



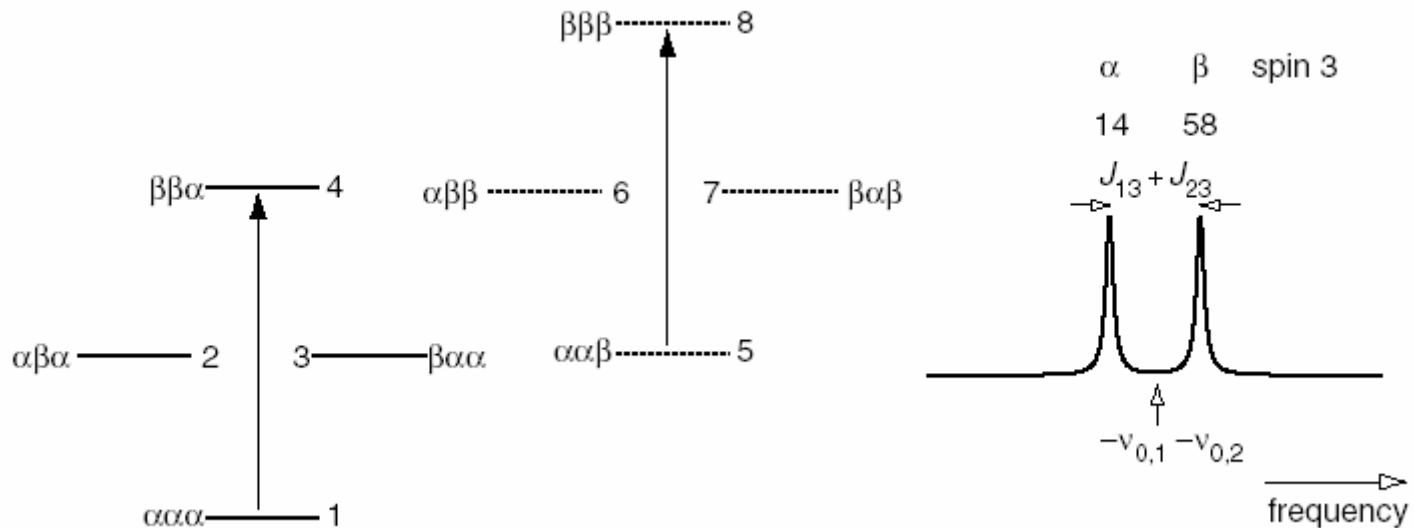
$$\nu_{0,1}^{\alpha 3} = \nu_{0,1} + \frac{1}{2}J_{13} \quad \nu_{0,2}^{\alpha 3} = \nu_{0,2} + \frac{1}{2}J_{23}.$$

$$\nu_{0,1}^{\beta 3} = \nu_{0,1} - \frac{1}{2}J_{13} \quad \nu_{0,2}^{\beta 3} = \nu_{0,2} - \frac{1}{2}J_{23}.$$

NMR a energiové hladiny

Tři spiny – více-kvantové přechody

| transition | initial state | final state | frequency |
|------------|----------------------|--------------------|--|
| 1-4 | $\alpha\alpha\alpha$ | $\beta\beta\alpha$ | $-\nu_{0,1} - \nu_{0,2} - \frac{1}{2}J_{13} - \frac{1}{2}J_{23}$ |
| 5-8 | $\alpha\alpha\beta$ | $\beta\beta\beta$ | $-\nu_{0,1} - \nu_{0,2} + \frac{1}{2}J_{13} + \frac{1}{2}J_{23}$ |
| 1-7 | $\alpha\alpha\alpha$ | $\beta\alpha\beta$ | $-\nu_{0,1} - \nu_{0,3} - \frac{1}{2}J_{12} - \frac{1}{2}J_{23}$ |
| 2-8 | $\alpha\beta\alpha$ | $\beta\beta\beta$ | $-\nu_{0,1} - \nu_{0,3} + \frac{1}{2}J_{12} + \frac{1}{2}J_{23}$ |
| 1-6 | $\alpha\alpha\alpha$ | $\alpha\beta\beta$ | $-\nu_{0,2} - \nu_{0,3} - \frac{1}{2}J_{12} - \frac{1}{2}J_{13}$ |
| 3-8 | $\beta\alpha\alpha$ | $\beta\beta\beta$ | $-\nu_{0,2} - \nu_{0,3} + \frac{1}{2}J_{12} + \frac{1}{2}J_{13}$ |



NMR a energiové hladiny

Dva spiny – silná interakce

| transition | frequency | intensity |
|------------|---|----------------------|
| 1-2 | $\frac{1}{2}D - \frac{1}{2}\Sigma - \frac{1}{2}J_{12}$ | $(1 + \sin 2\theta)$ |
| 3-4 | $\frac{1}{2}D - \frac{1}{2}\Sigma + \frac{1}{2}J_{12}$ | $(1 - \sin 2\theta)$ |
| 1-3 | $-\frac{1}{2}D - \frac{1}{2}\Sigma - \frac{1}{2}J_{12}$ | $(1 - \sin 2\theta)$ |
| 2-4 | $-\frac{1}{2}D - \frac{1}{2}\Sigma + \frac{1}{2}J_{12}$ | $(1 + \sin 2\theta)$ |

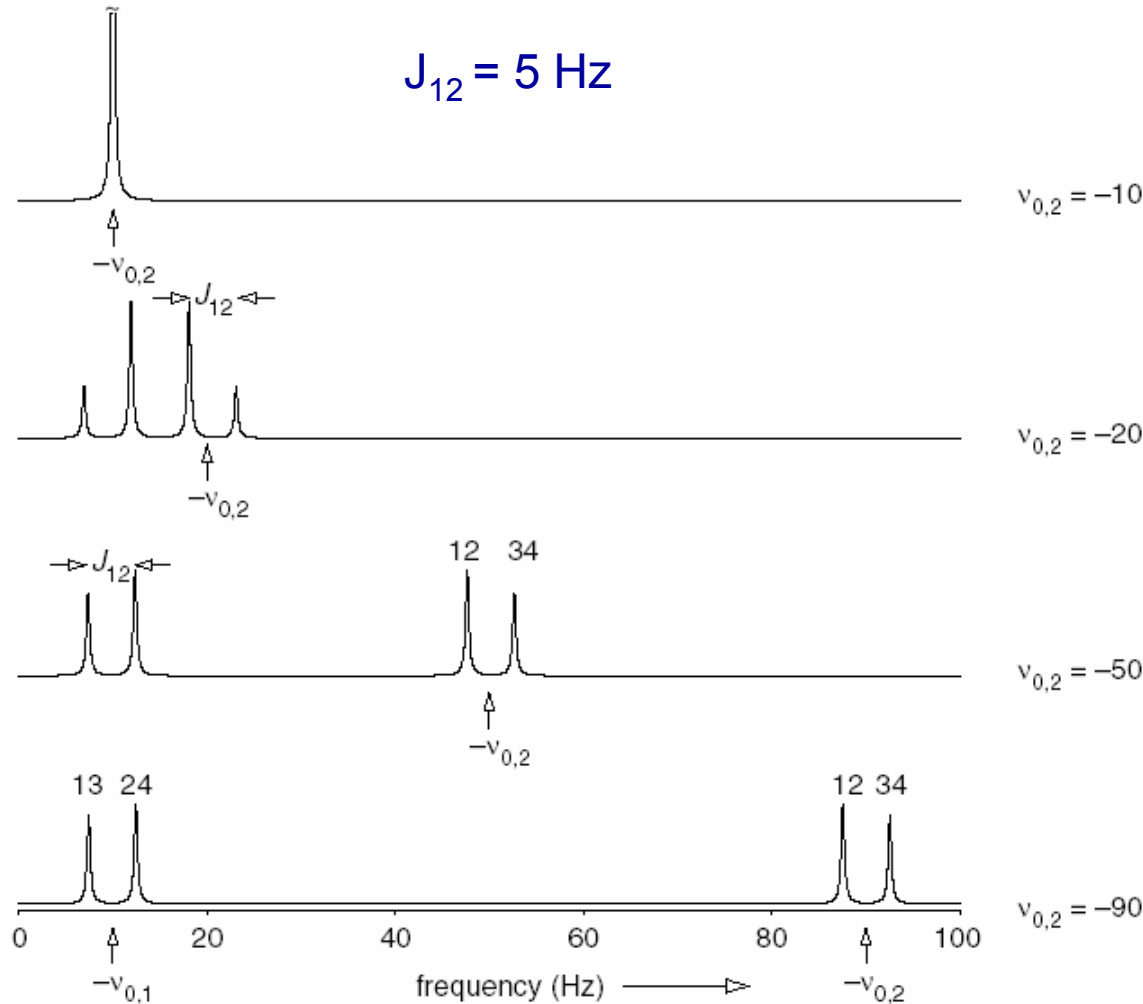
$$\Sigma = \nu_{0,1} + \nu_{0,2} \quad D^2 = (\nu_{0,1} - \nu_{0,2})^2 + J_{12}^2. \quad \sin 2\theta = \frac{J_{12}}{D}.$$

$$D^2 = (\nu_{0,1} - \nu_{0,2})^2 + J_{12}^2 \\ \approx (\nu_{0,1} - \nu_{0,2})^2$$

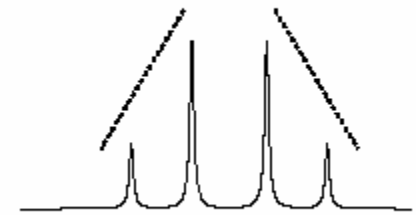
NMR a energiové hladiny

Dva spiny – silná interakce

$$\Delta\delta_{12} < 7 \cdot J_{12}$$



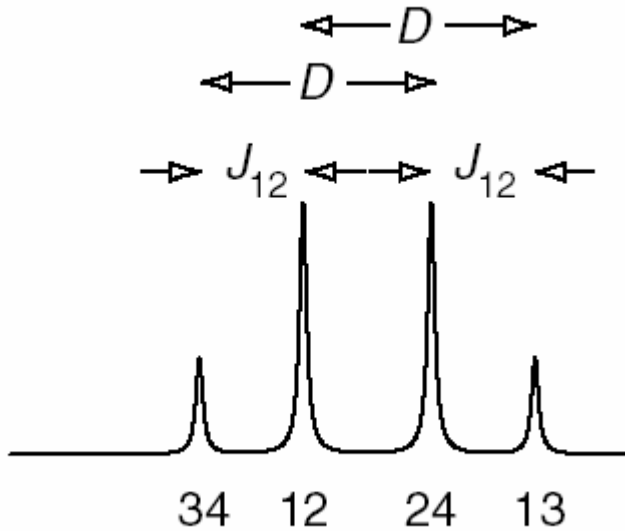
Střížkový efekt



NMR a energiové hladiny

Dva spiny – silná interakce

$$\Delta\delta_{12} < 7 \cdot J_{12}$$



$$\begin{aligned} \nu_{34} - \nu_{24} &= \left(\frac{1}{2}D - \frac{1}{2}\Sigma + \frac{1}{2}J_{12}\right) - \left(-\frac{1}{2}D - \frac{1}{2}\Sigma + \frac{1}{2}J_{12}\right) \\ &= D. \end{aligned}$$

$$D^2 = (\nu_{0,1} - \nu_{0,2})^2 + J_{12}^2$$

$$\begin{aligned} \nu_{12} + \nu_{24} &= \left(\frac{1}{2}D - \frac{1}{2}\Sigma - \frac{1}{2}J_{12}\right) + \left(-\frac{1}{2}D - \frac{1}{2}\Sigma + \frac{1}{2}J_{12}\right) \\ &= -\Sigma. \end{aligned}$$

$$\nu_{0,1} = \frac{1}{2}(\Sigma + (\nu_{0,1} - \nu_{0,2})) \quad \nu_{0,2} = \frac{1}{2}(\Sigma - (\nu_{0,1} - \nu_{0,2}))$$

NMR a energiové hladiny

Tři spiny – ABX systém - subspektra

