



# **Lasers - instrumentation**

Vítězslav Otruba

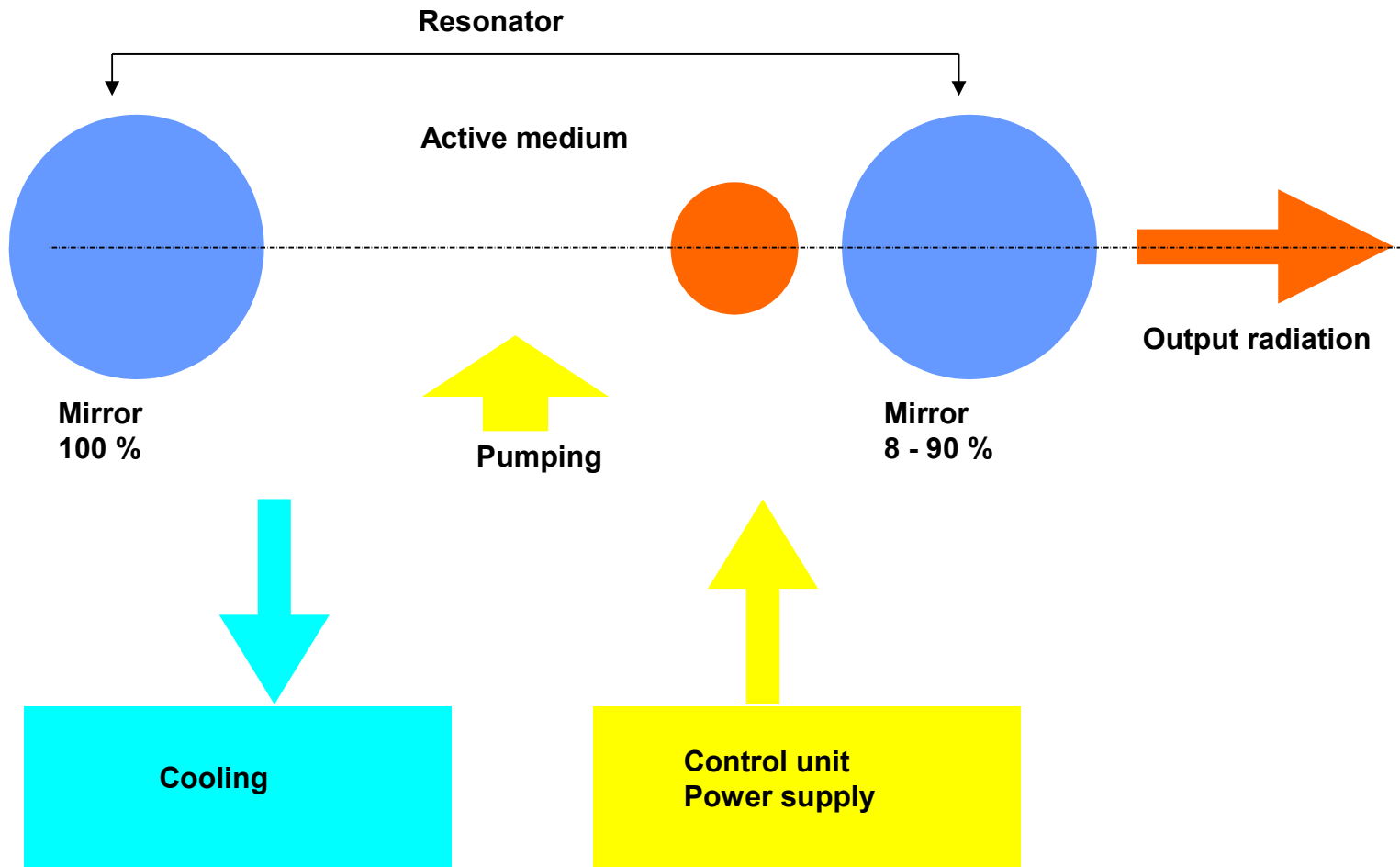
1

# Less common prefixes

prefix	symbol	exponent
atto	a	$10^{-18}$
femto	f	$10^{-15}$
piko	p	$10^{-12}$
nano	n	$10^{-9}$
giga	G	$10^9$
tera	T	$10^{12}$
peta	P	$10^{15}$
exa	E	$10^{18}$

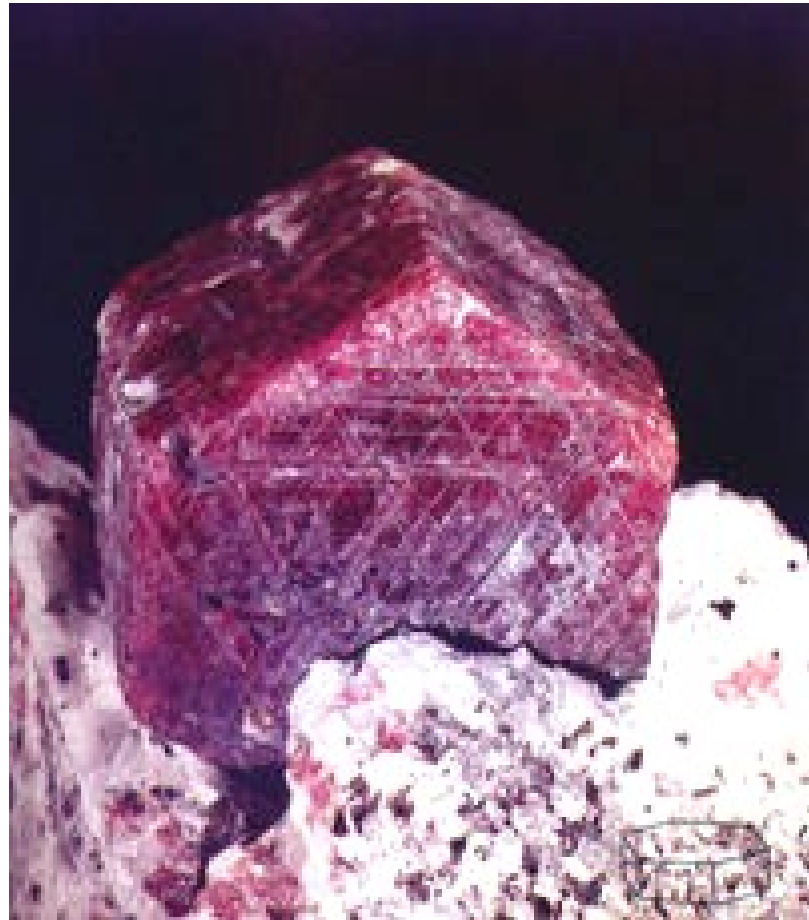


# Solid-state laser

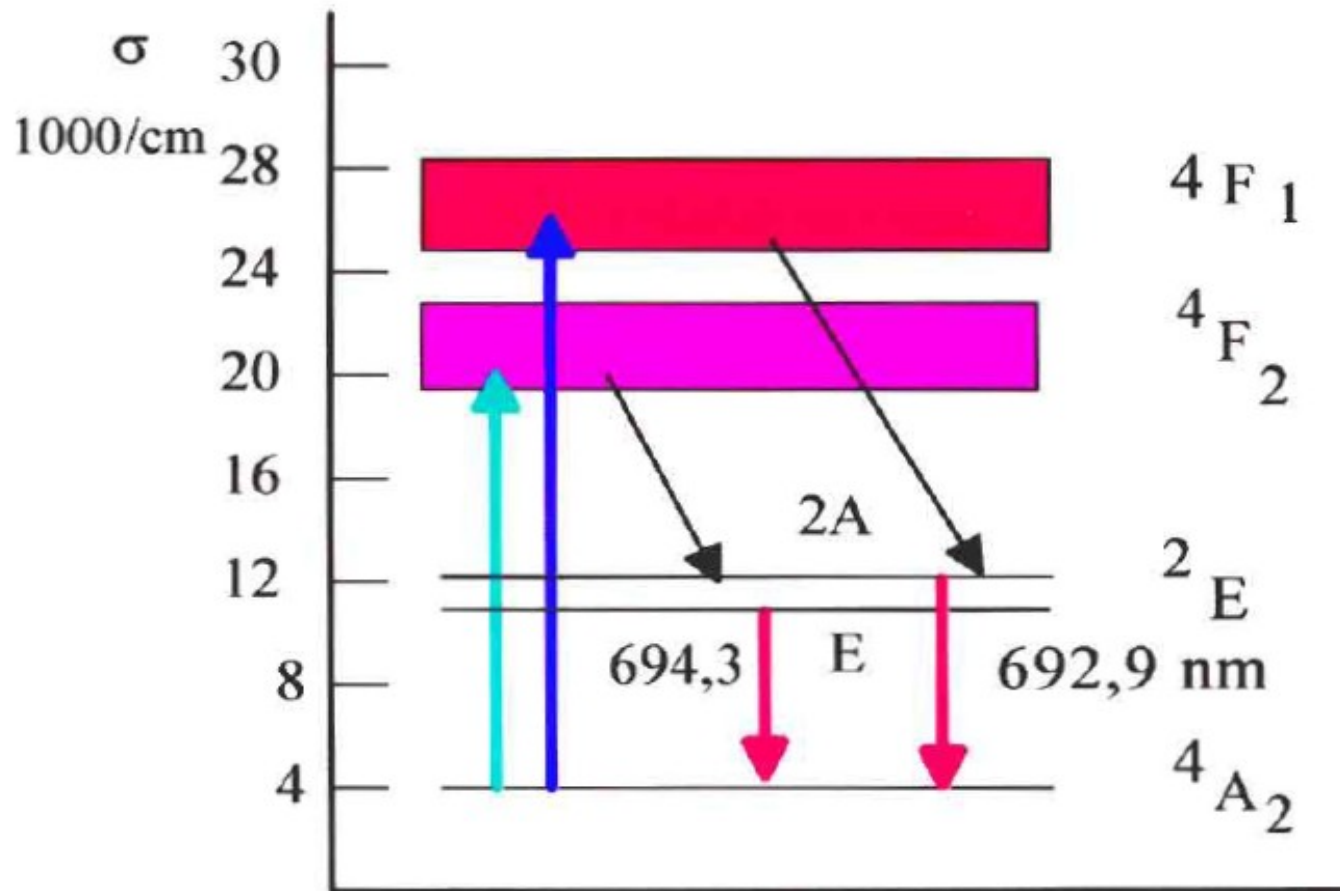


# Ruby laser ( $\text{Cr}^{3+}:\text{Al}_2\text{O}_3$ )

- The first laser designed by T. Maiman in 1960.
- Used in pulse mode, power in free running mode to 10J (1ms),
- in Q-switched mode to 5J (1 – 10 ns)
- three-level system

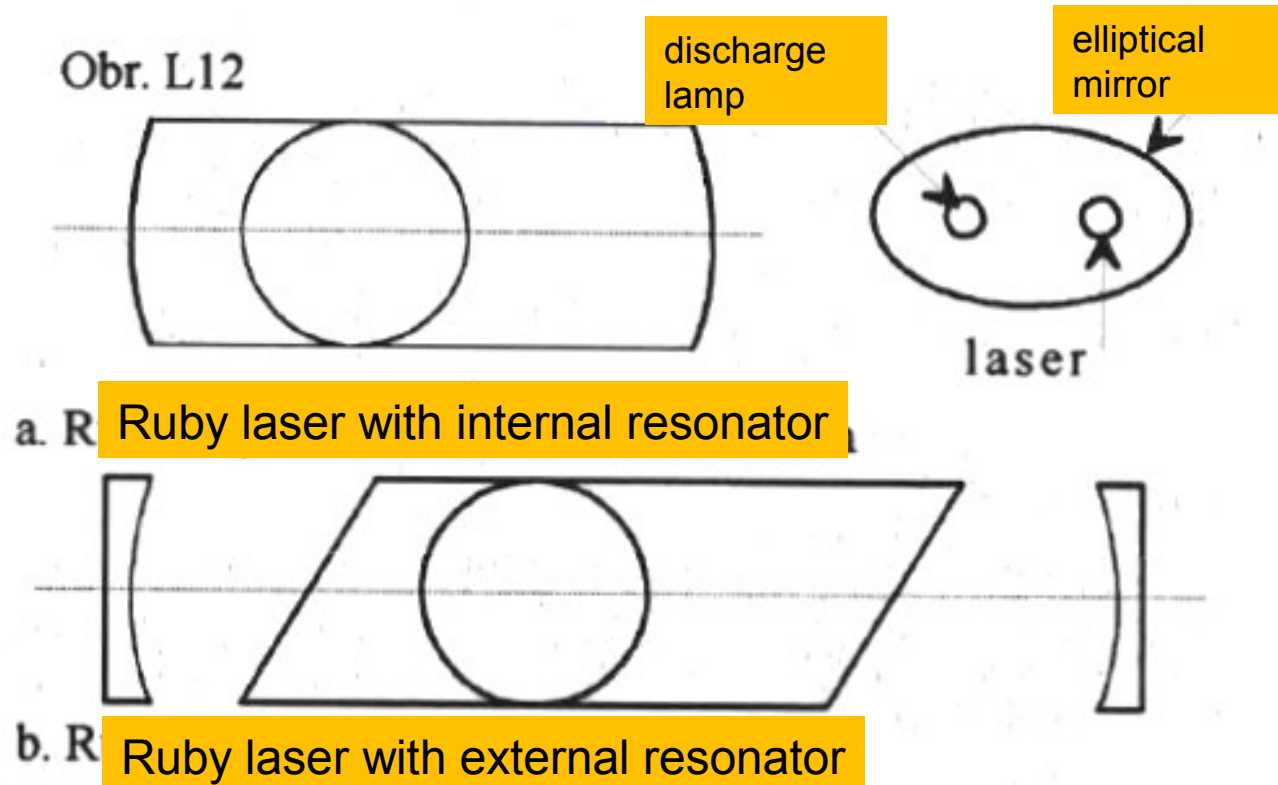


# diagram of chromium energy levels in ruby laser



# ruby laser

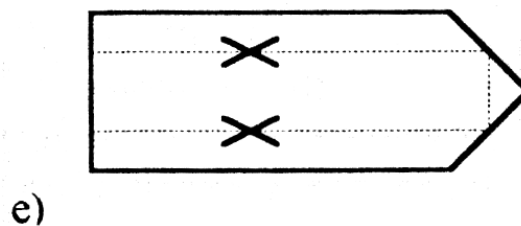
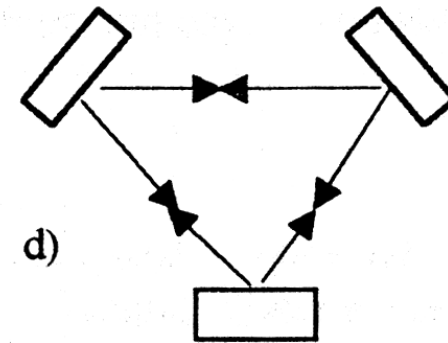
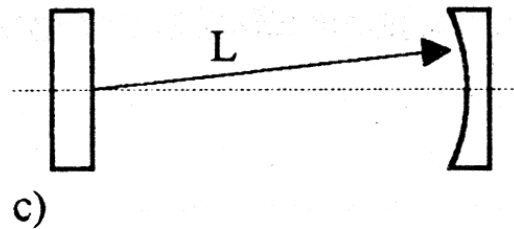
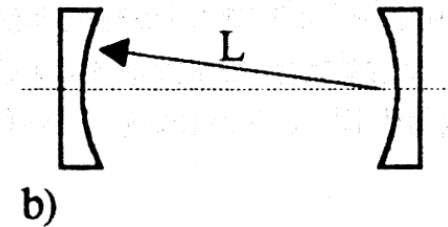
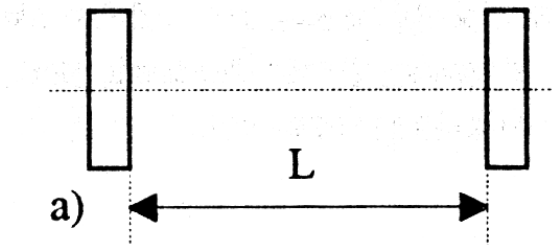
- arrangement of ruby laser



# Resonators

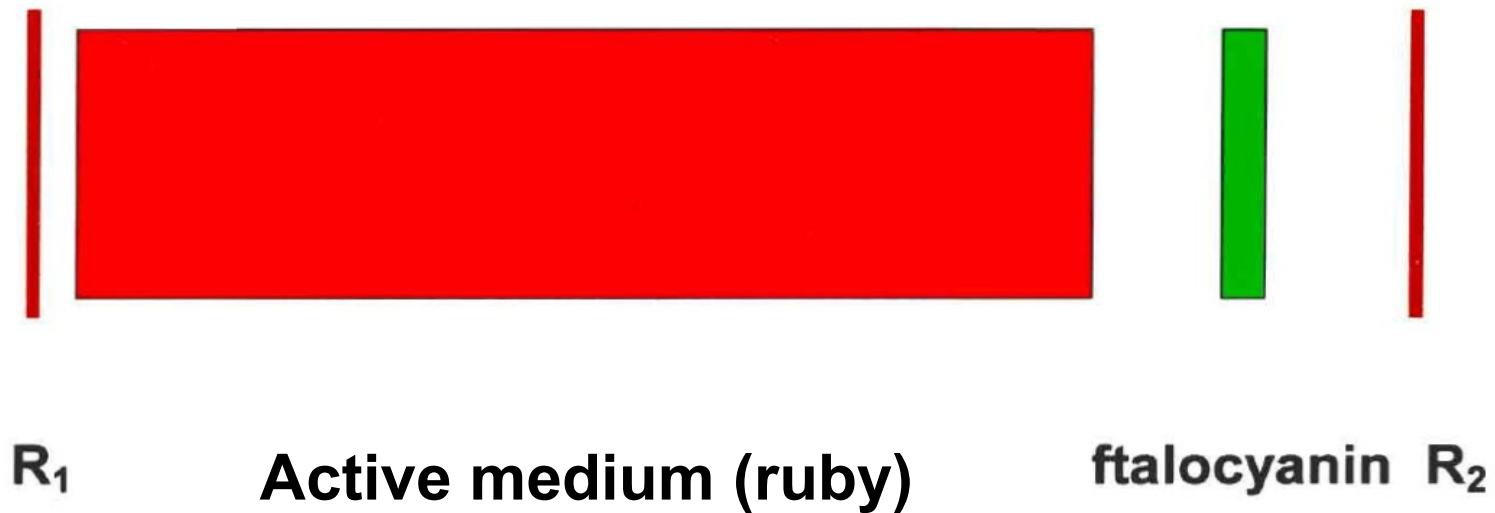
## ○ Fabry-Perot etalon

- ❖ plan parallel
- ❖ confocal
- ❖ hemispherical
- ❖ circular
- ❖ roof



# Passive Q-modulation

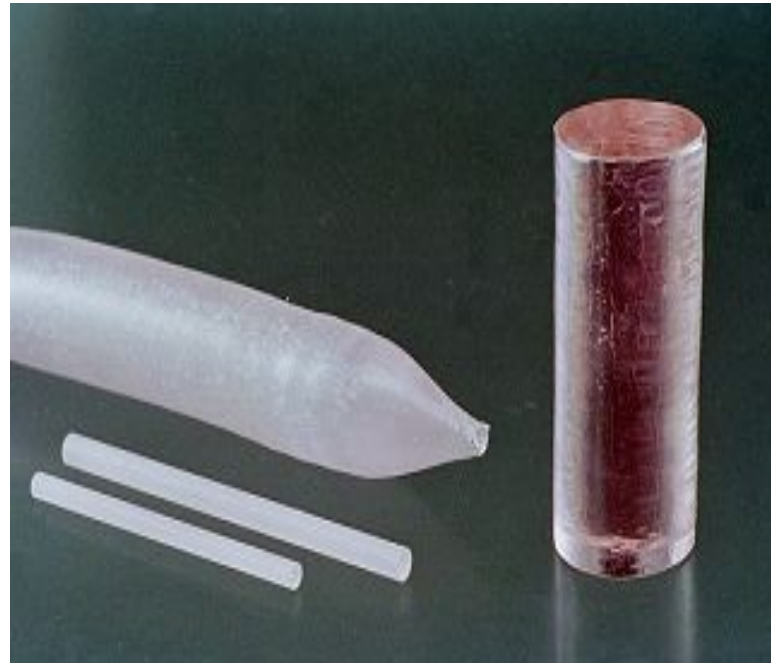
- Example of use of saturation absorber for generation of short (nanosecond) power pulses (GW) in ruby laser





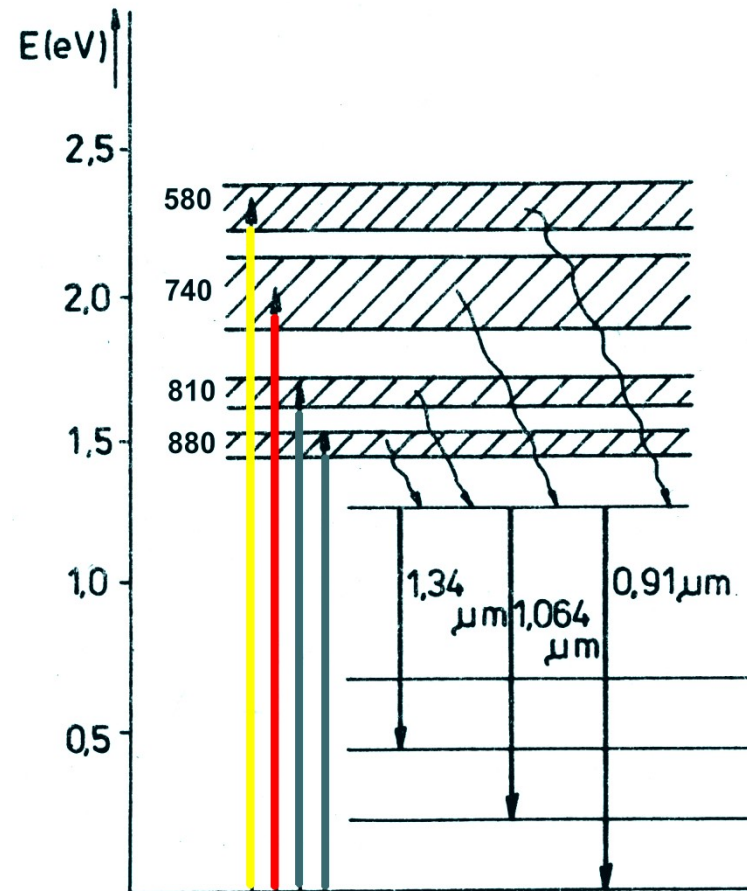
# Neodymium laser

- It is the most widespread solid state laser (cca 1% Nd v Y3Al5O12). It operates at 1,064 nm, in continuous mode power to 1 kW,
- pulse to 10 J, repetition rate up to kHz. In Q-switched pulse mode 1 – 10 ns,
- at synchronizing modes to 10 ps.

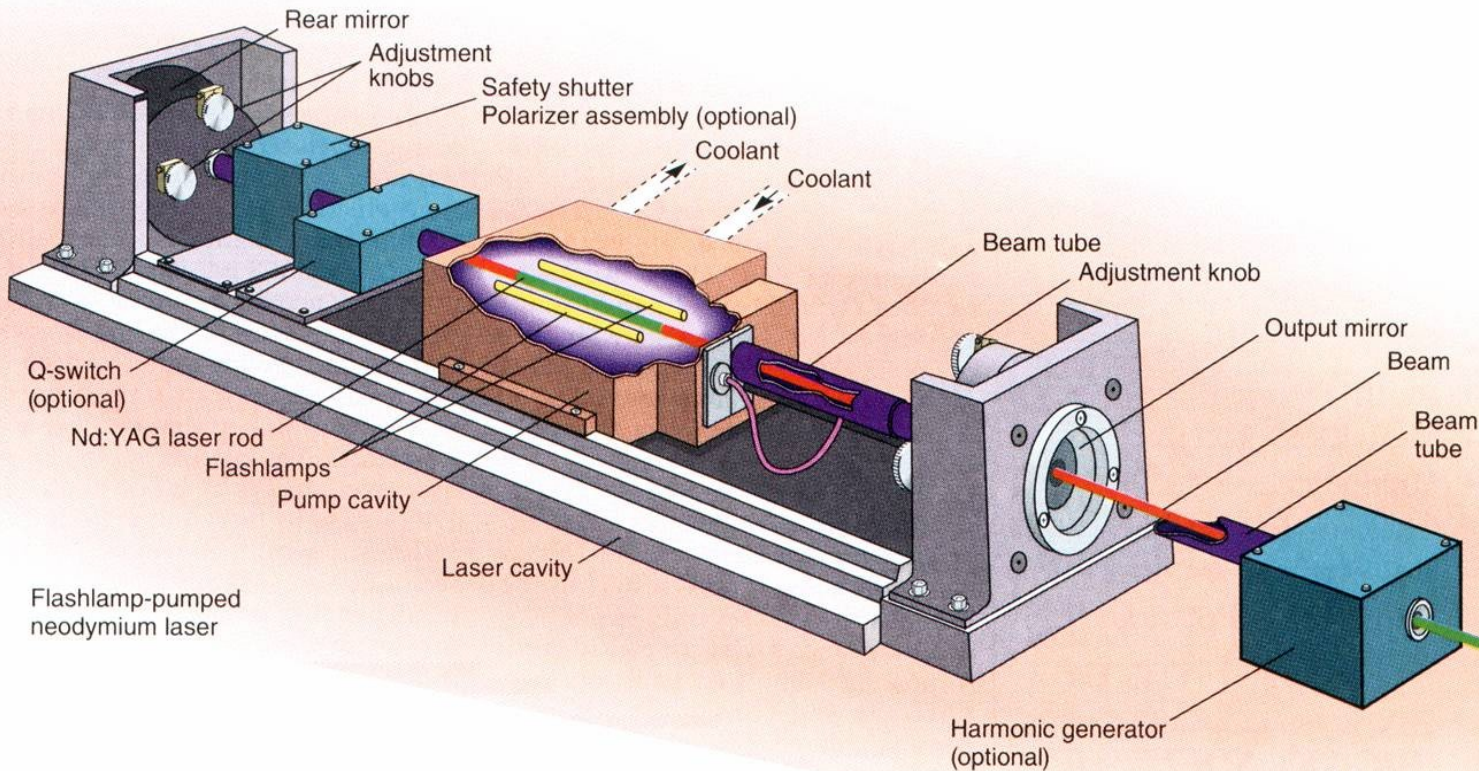


# Neodymium energy level diagram of Nd in: YAG laser

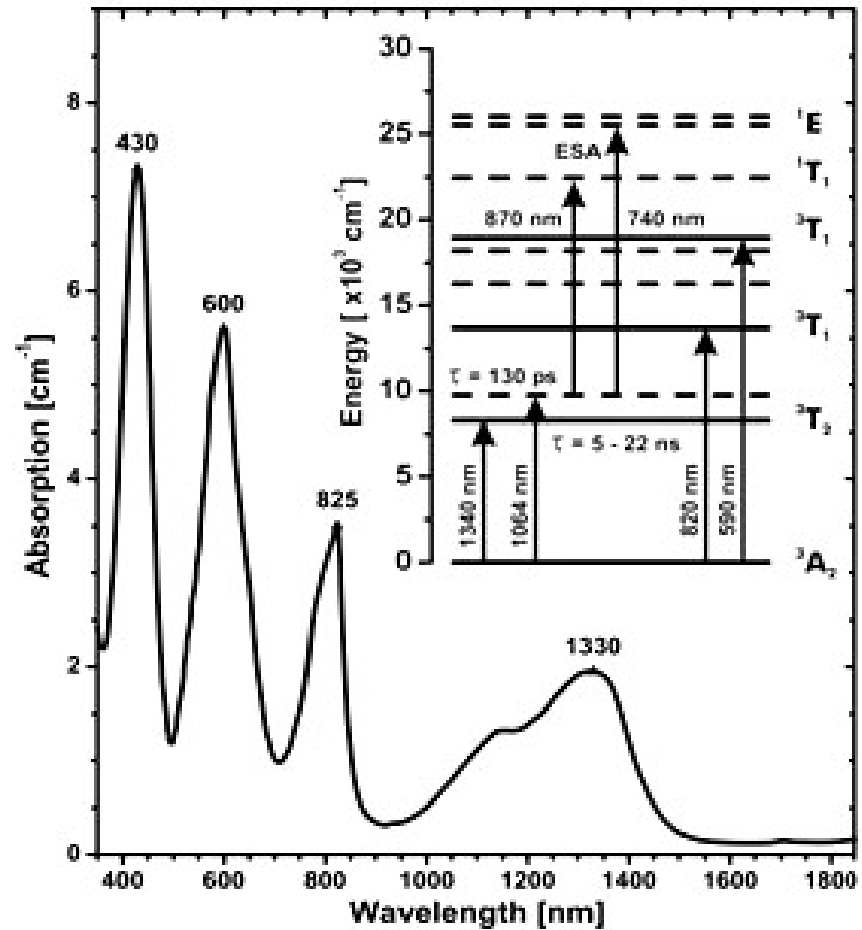
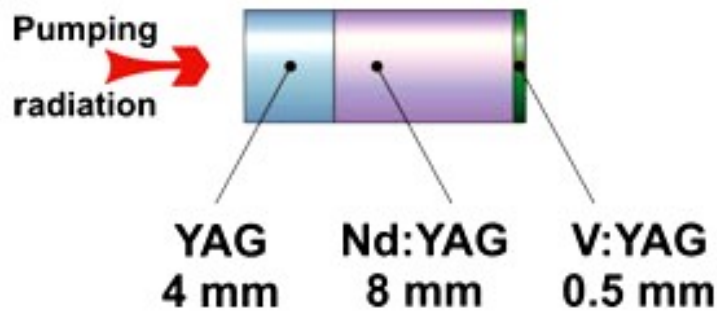
- Nd<sup>3+</sup> in yttrium-aluminum garnet (Y<sub>3</sub>Al<sub>5</sub>O<sub>12</sub>) replaces ions Y<sup>3+</sup>. Monocrystals are mechanically strong, thermally stable with minimal optical defects unlike neodymium glasses. Xenon lamps or laser or LED diodes are used for pumping.



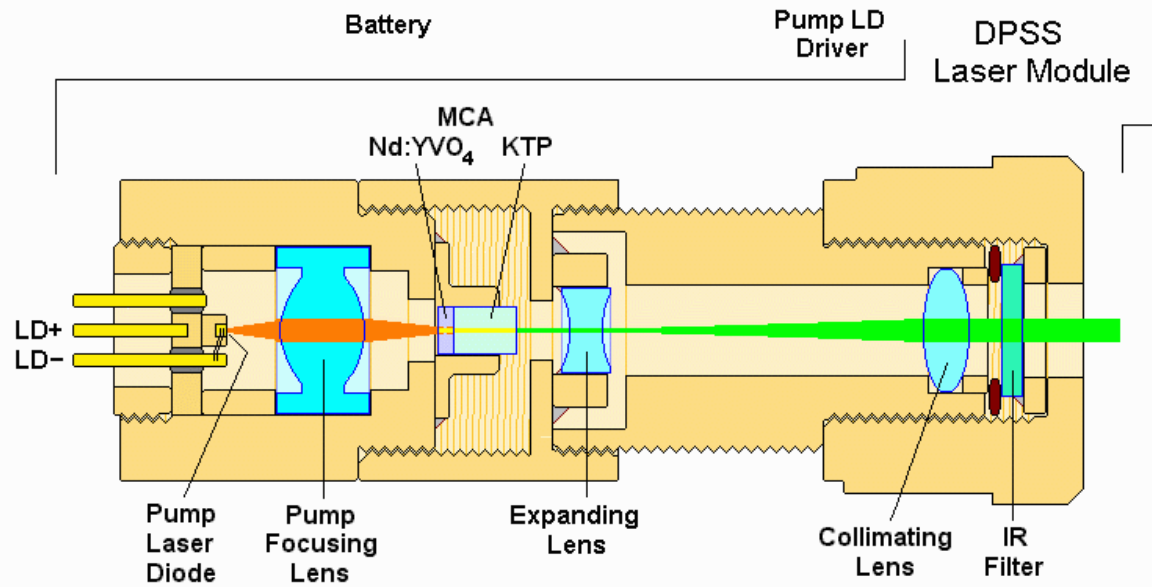
# Neodymium laser



# Passive Q-modulation V3+:YAG



# Green laser pointer

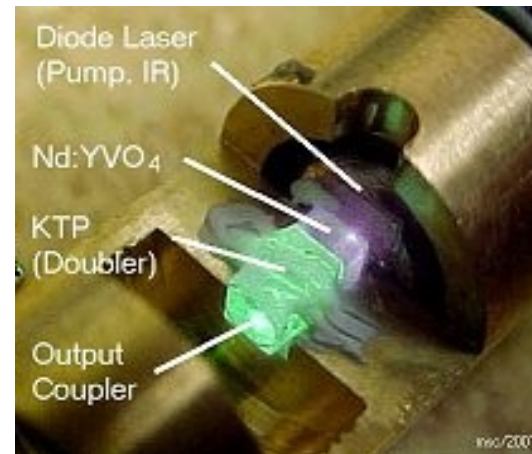
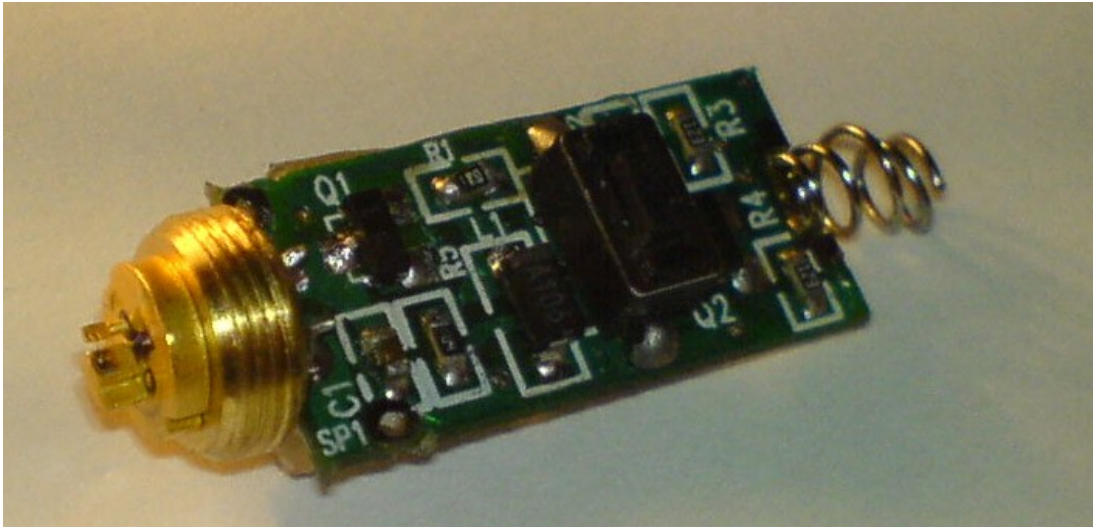


Beam Paths: 808 nm — 1064+532 nm — 532 nm —



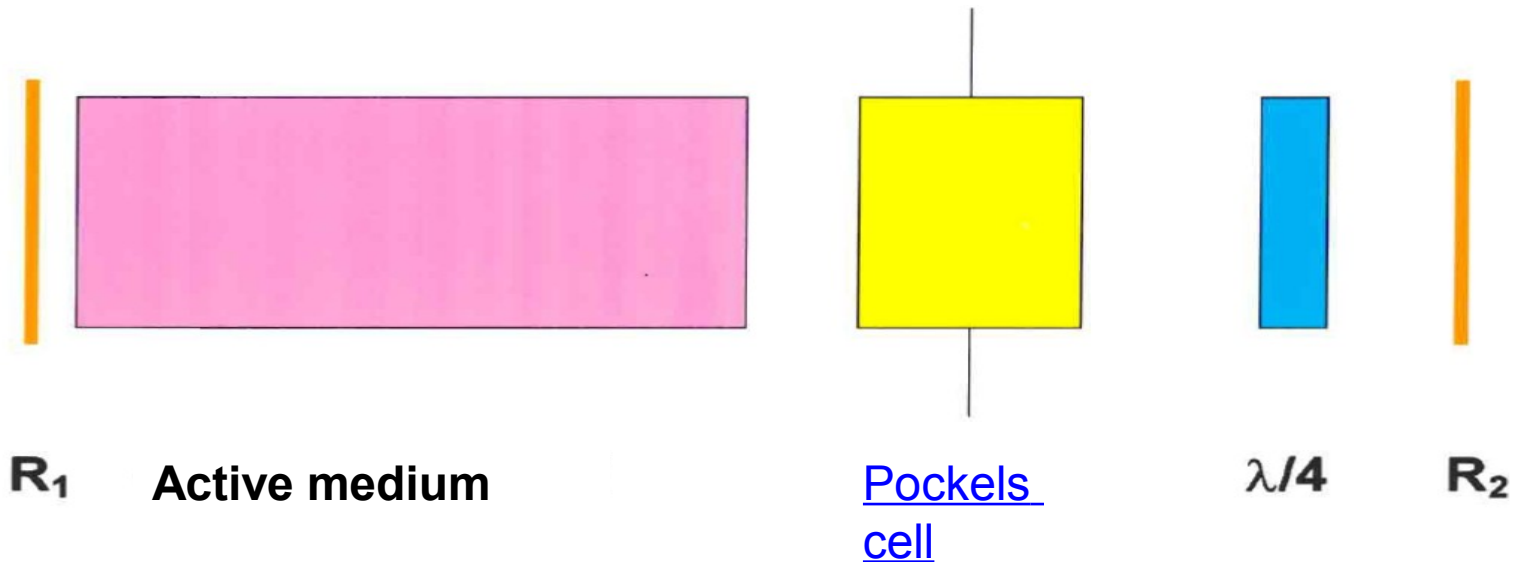


# Green laser pointer



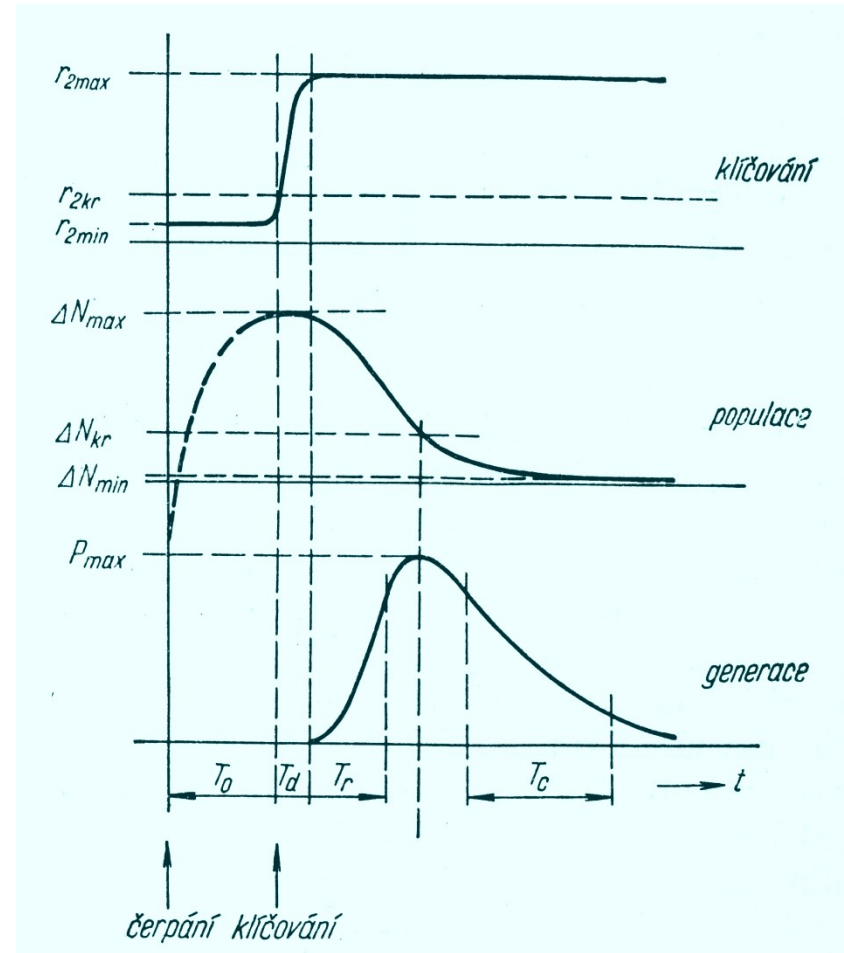
# Active Q-modulation

- In this case, the Q of resonator is modulated by optical shutters, such as a Kerr effect electro-optical modulator or an acoustooptical module.



# Timing of active Q-modulation

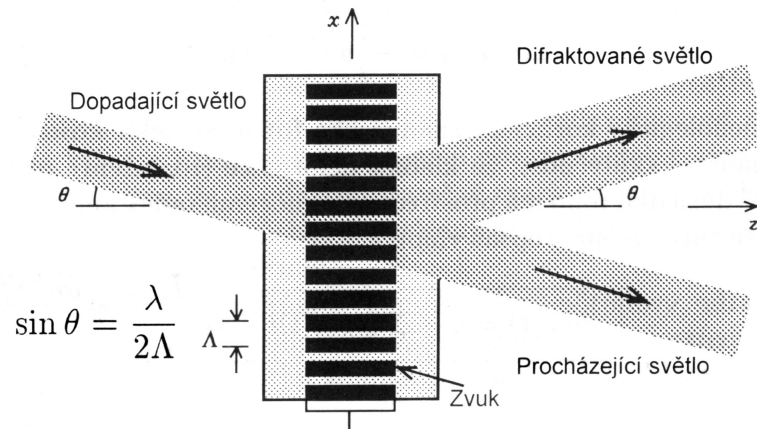
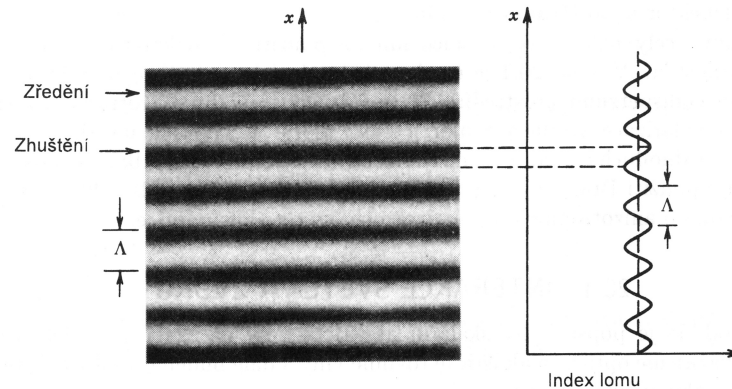
- Typical timing for Nd:YAG laser:
- Inverse population increase (T<sub>0</sub>) 150μs
- Trigger pulse (TD) 1 ns
- Generation of radiation 10ns





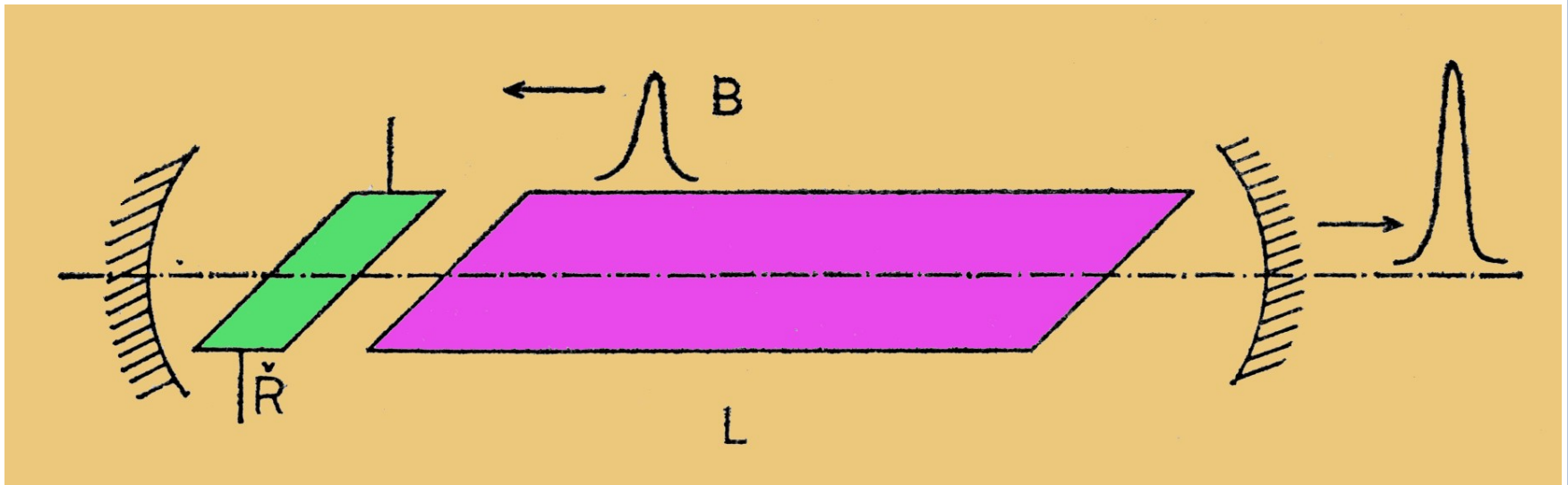
# Optoacoustic modulator

- If the sound propagates through the optical environment, the density and thus the refractive index change.
- The simplest is the Bragg diffraction: the acoustic plane wave causes a partial reflection of the radiation if the angle  $\Theta$  satisfies the Bragg condition (Bragg cell)



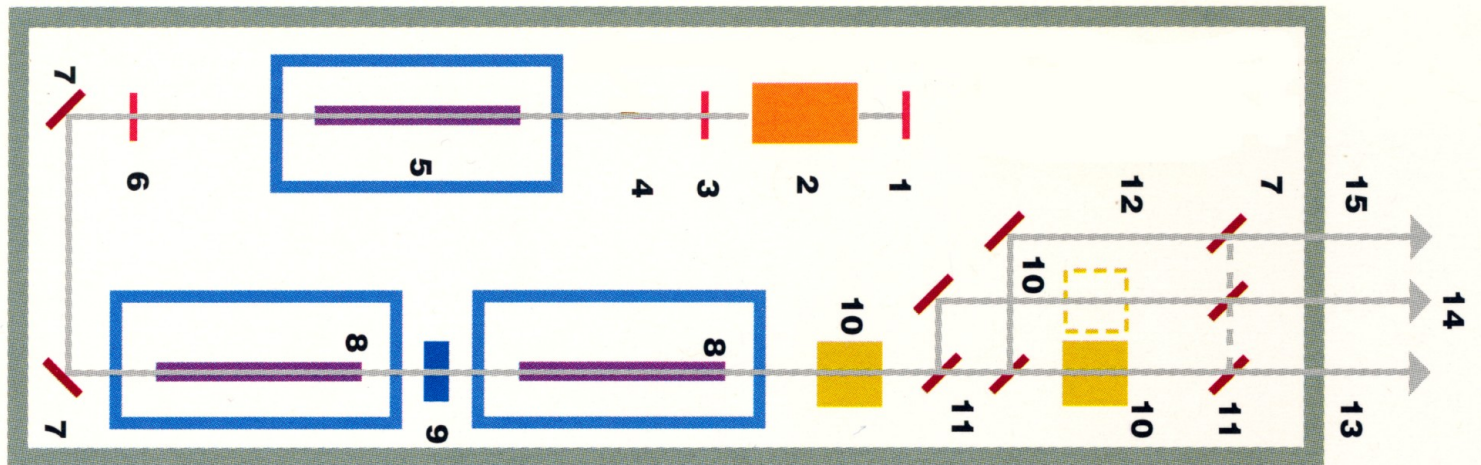
# Mode-locking

- Passive or active modulation of the resonator with frequencies  $f = c / 2L$  gives a sequence of very short pulses, the length of which is determined by the Fourier transform of the spectral line and the repetition rate by the time the photon cloud travels back and forth.



# Nd:YAG laser

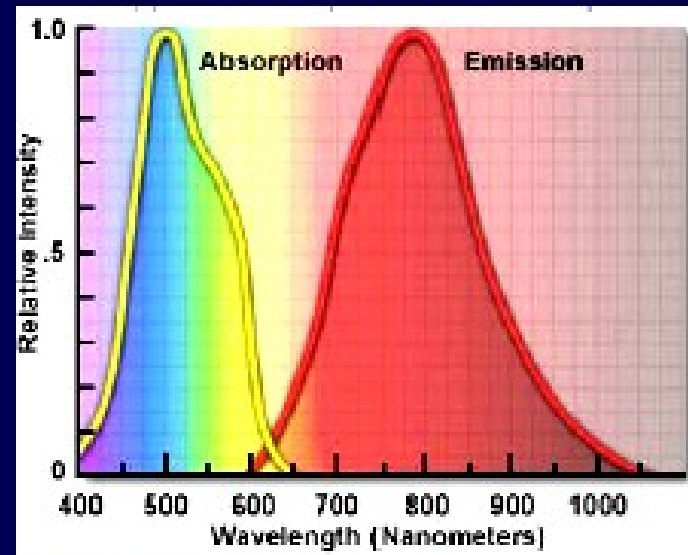
- Laser with Q-modulation (1-6), two-stage amplifier (8), birefringence compensator (9) and harmonic frequency generators (10), output 1064 nm (13), 532/355 nm (14), 266/1064 nm residual (15)



# Titanium Doped Sapphire $\text{Al}_2\text{O}_3:\text{Ti}^{3+}$



## Ti:sapphire electronic spectra



- Excitation at 532 nm from a frequency doubled Nd:YVO<sub>4</sub> laser



# Titanium Doped Sapphire $\text{Al}_2\text{O}_3:\text{Ti}^{3+}$

- $\text{Ti}_2\text{O}_3$  concentration 0.06-0.5 wt%

Hardness 9 Mohs

Thermal conductivity 0.11 cal/  
( $^{\circ}\text{C} \times \text{sec} \times \text{cm}$ )

- **Optical Properties**

Laser action 4-Level Vibronic

Fluorescence lifetime 3.2  $\mu\text{sec}$   
( $T = 300 \text{ K}$ )

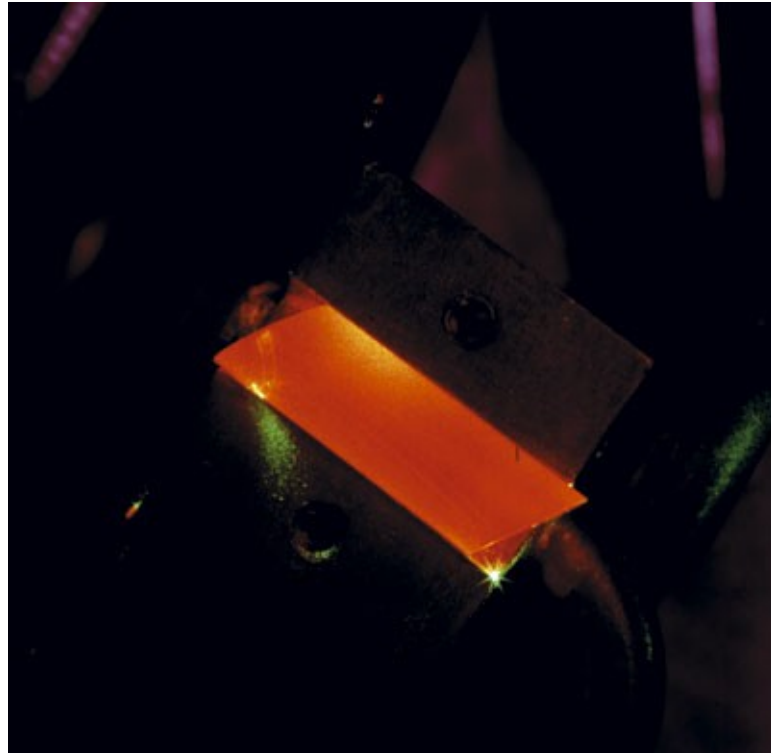
Tuning range 660-1050 nm

Absorption range 400-600 nm

Emission peak 795 nm

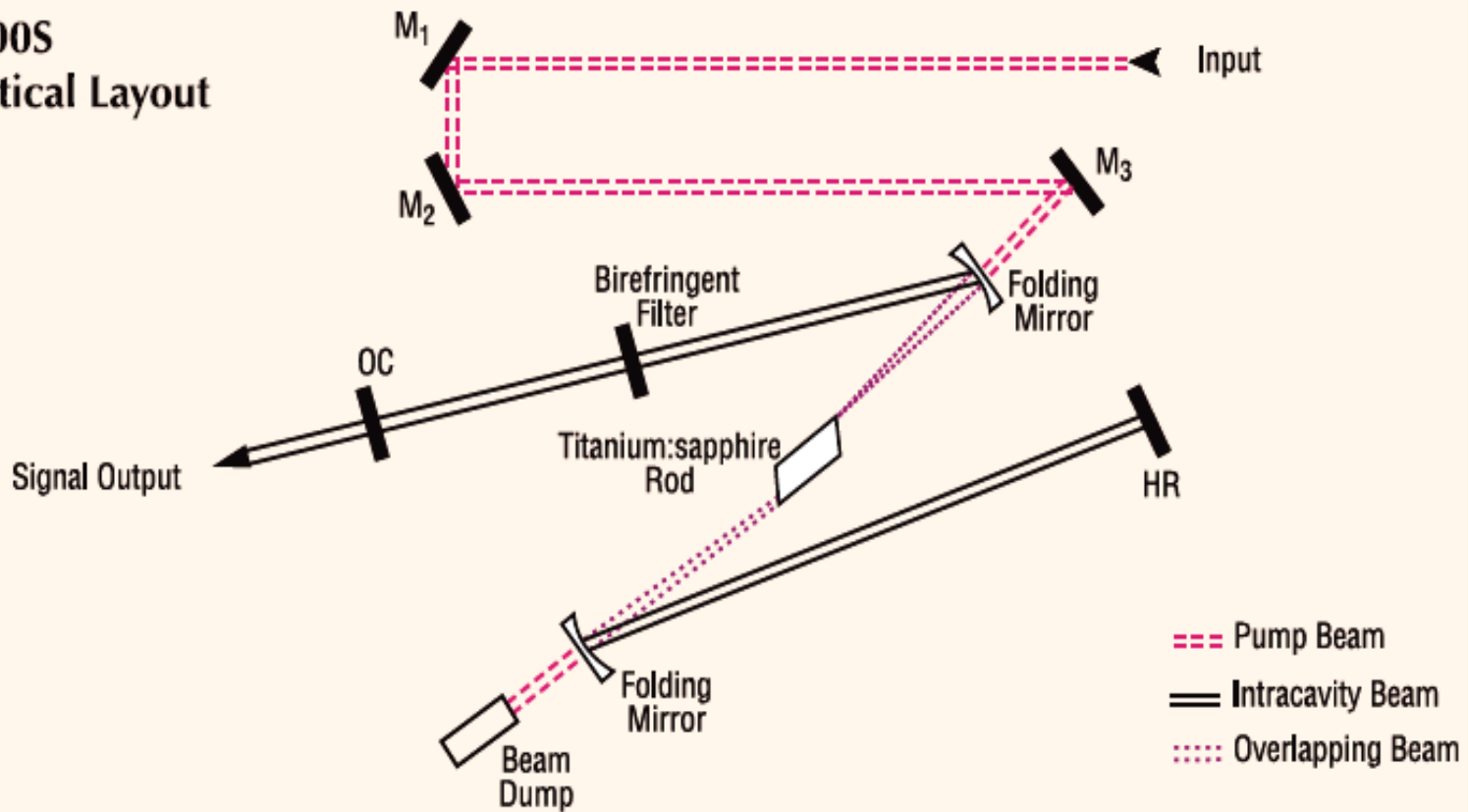
Absorption peak 488 nm

Refractive index 1.76 @ 800 nm



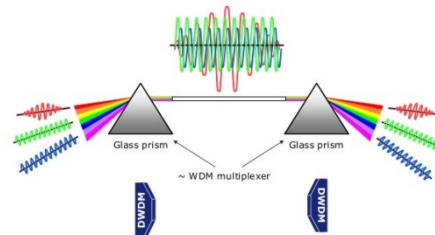
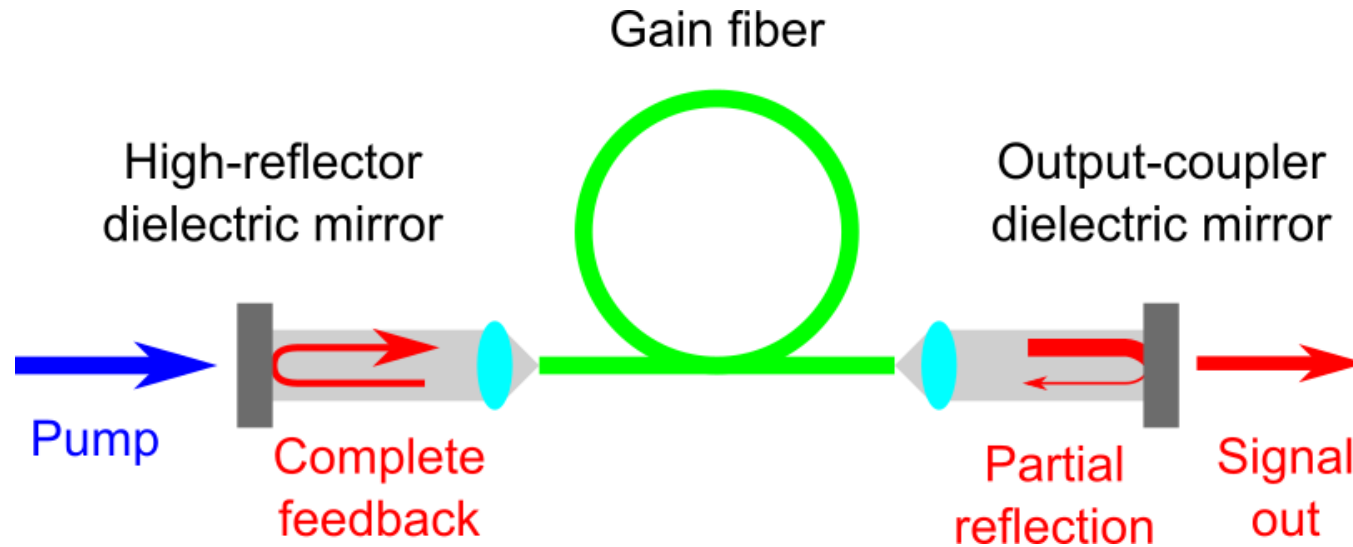
# Titanium Sapphire Laser

## 3900S Optical Layout



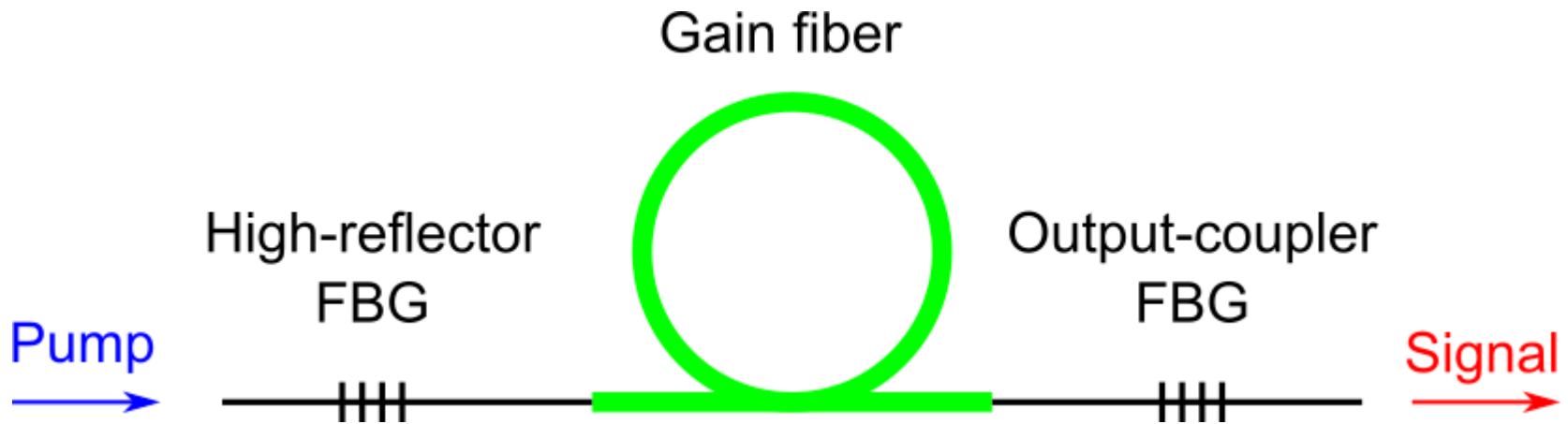
# Fiber lasers

- Arrangement with Linear Fabry Perot Resonator
- *WDM* (Wavelength Division Multiplex)



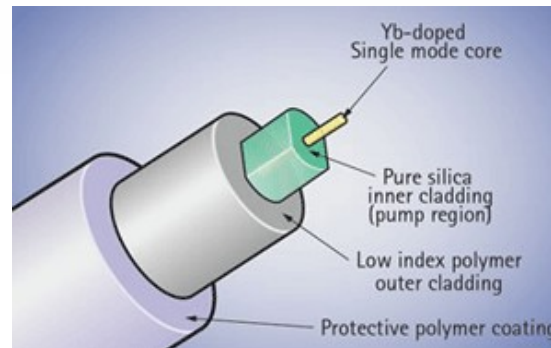
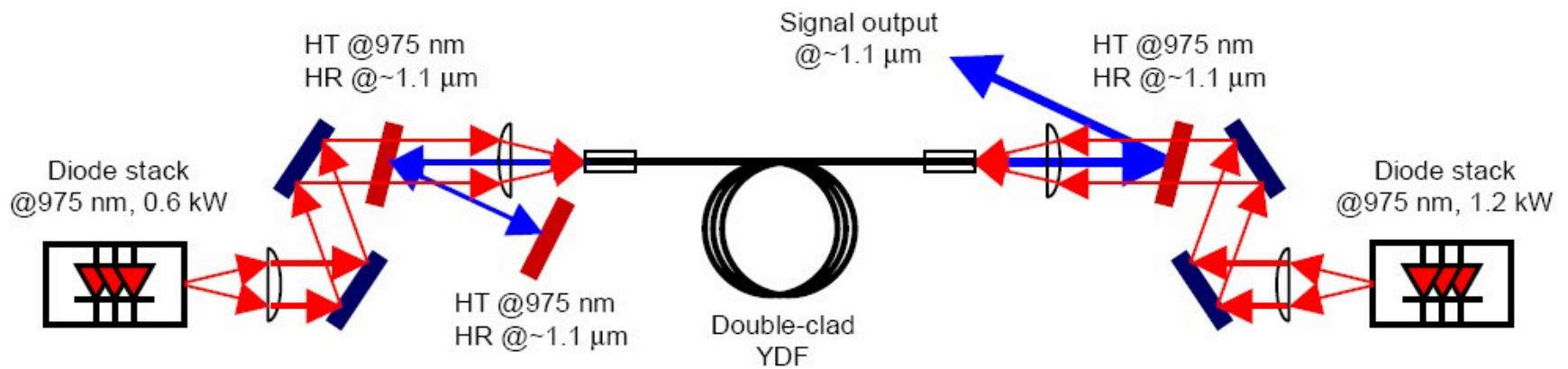
# Fiber lasers

- FBG – Fiber Bragg Grating

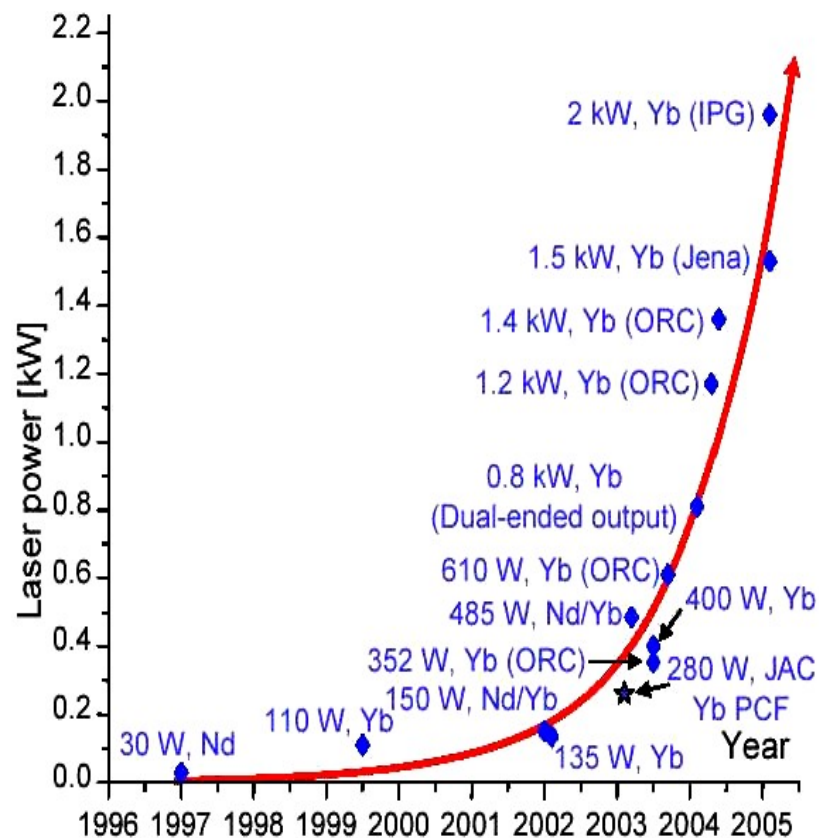
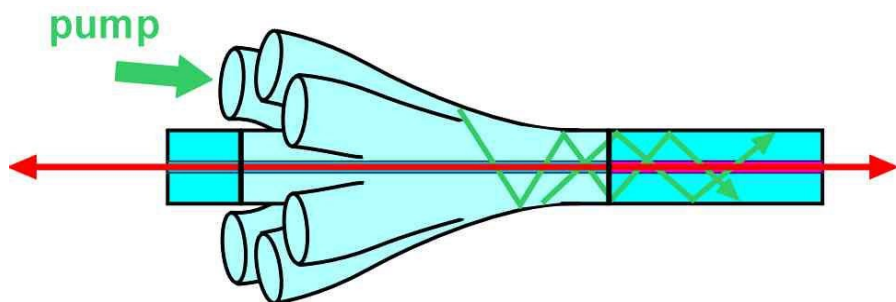
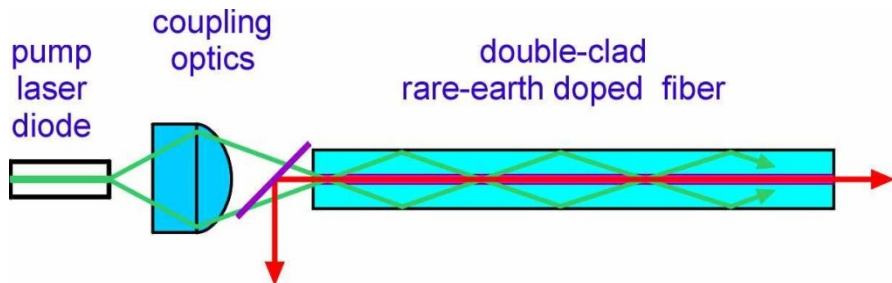
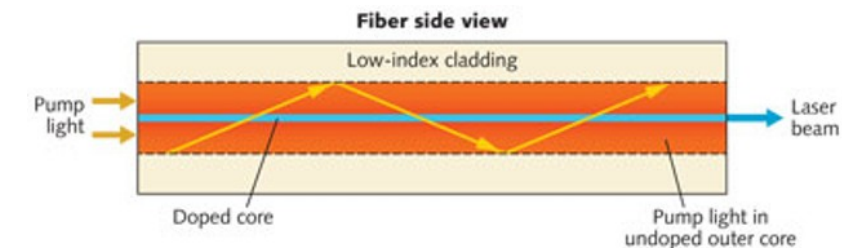
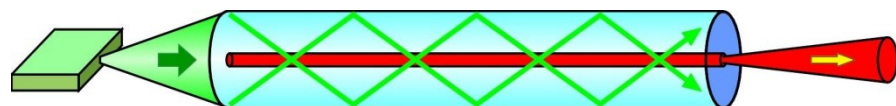




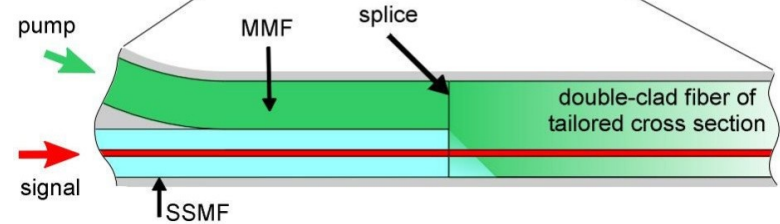
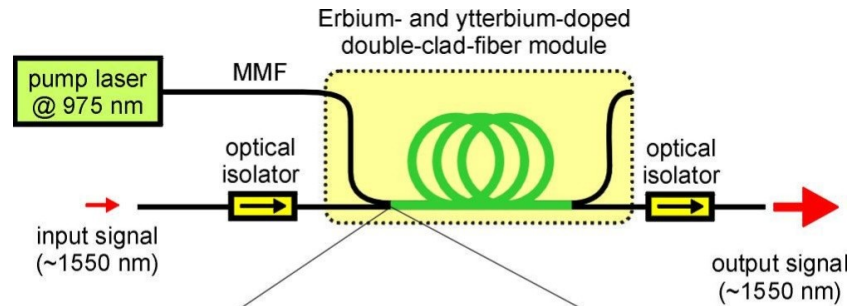
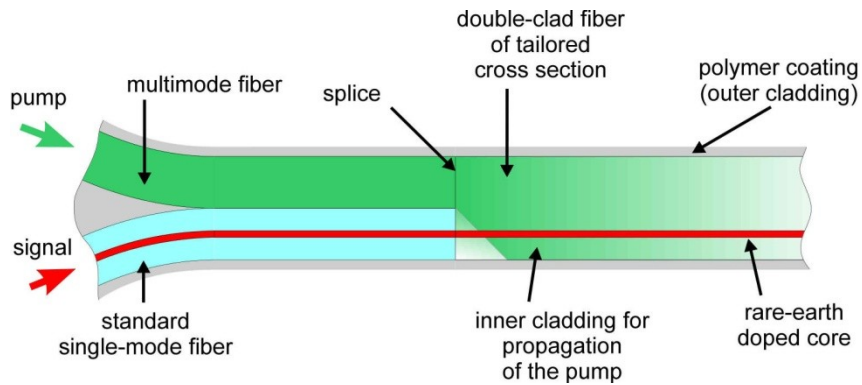
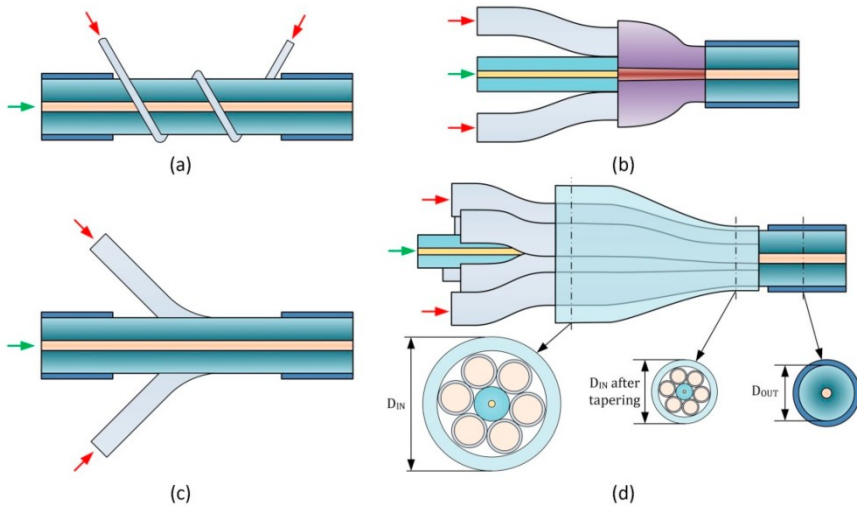
# Ytterbium doped fiber laser



# Optical excitation systems

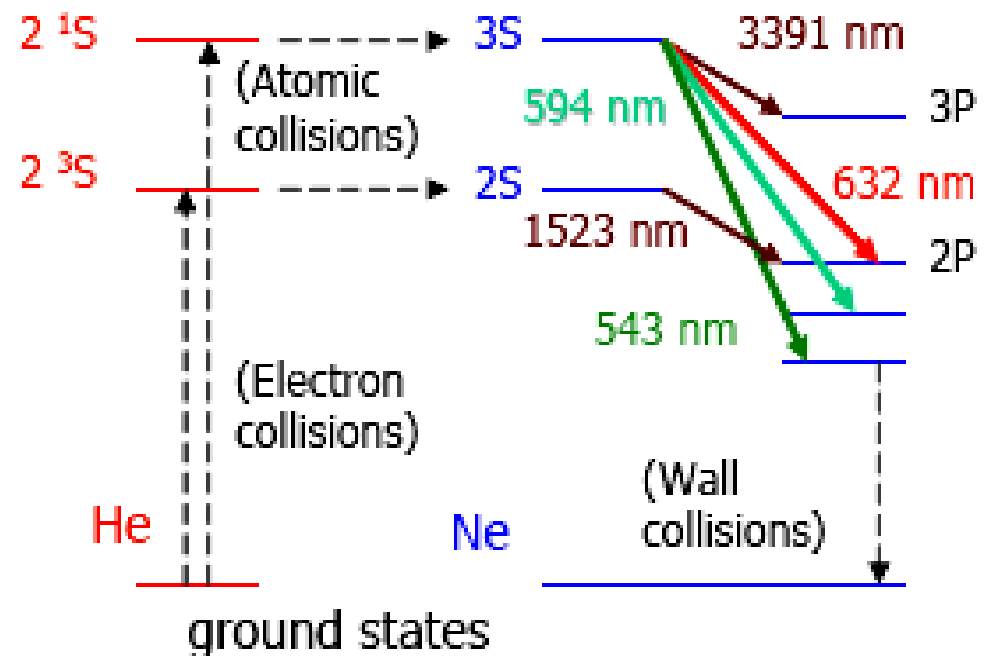


# Fiber laser amplifiers



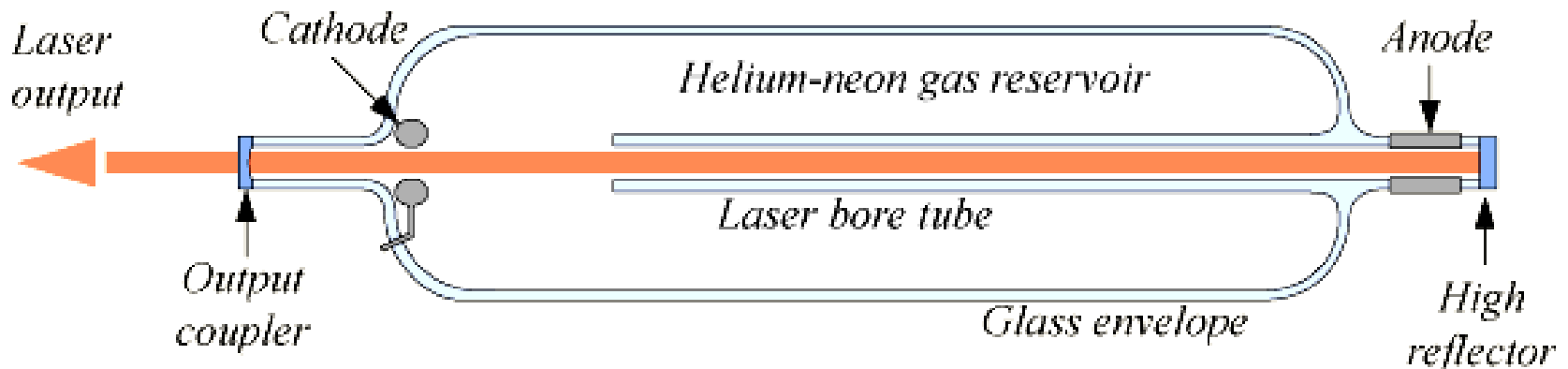
# Gas laser He-Ne

- Invented in 1960 as IR laser; red line used first in 1962
- Electric discharge in gas excites He to 2S levels
- Nearly parallel Ne levels exist
- Atomic collisions transfer excitation

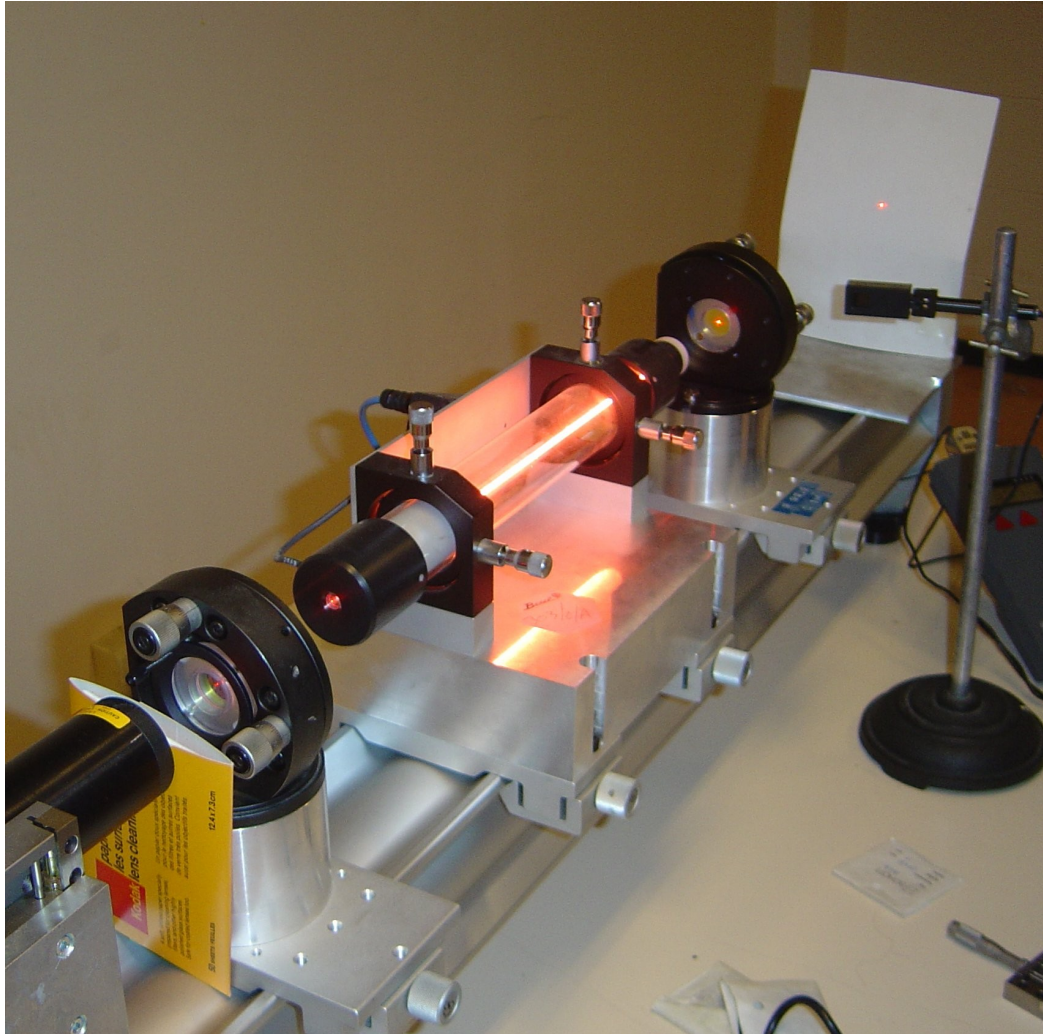


# Gas laser He-Ne

- Cheap and easy to manufacture – first lasers under \$100
- Gas tube has 85% He, 15% Ne

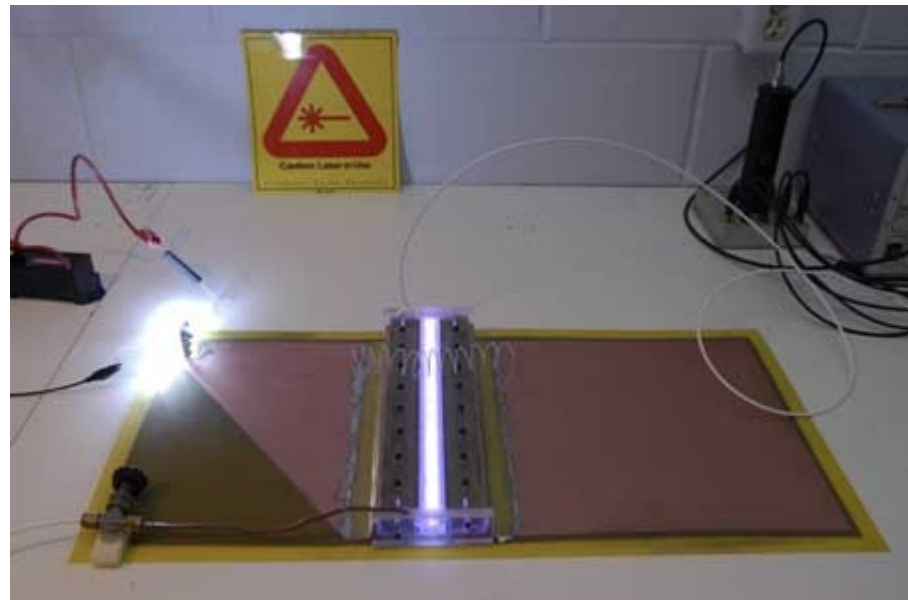
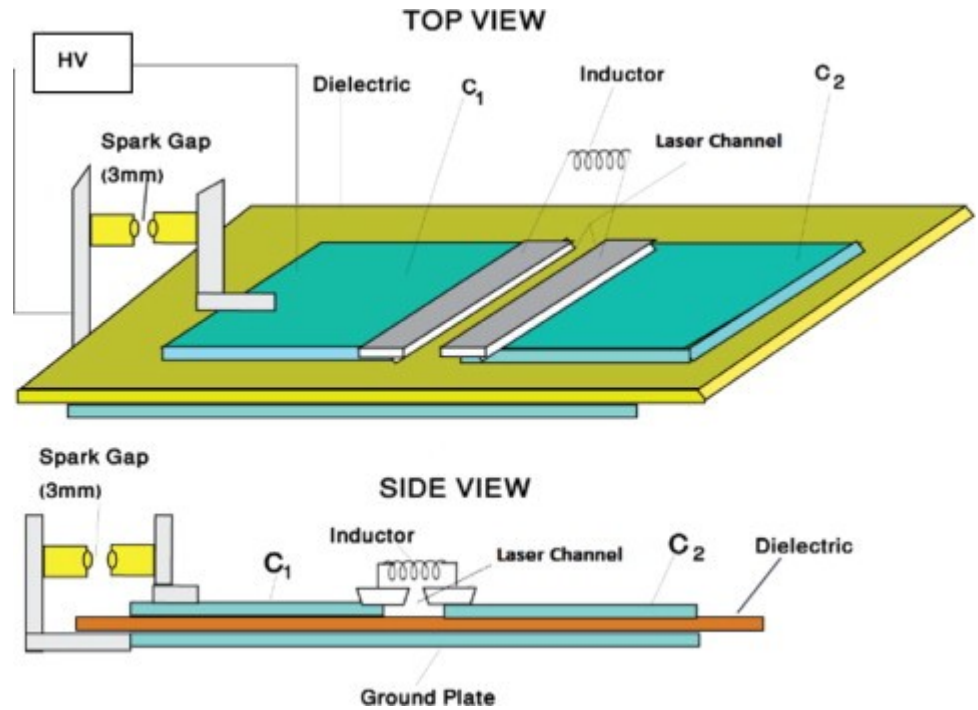
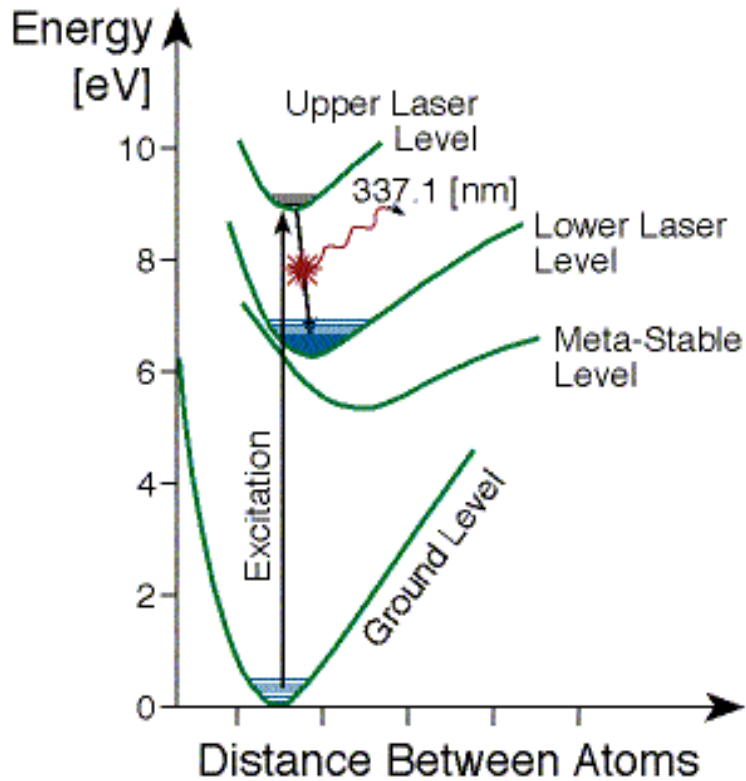


# He-Ne laser

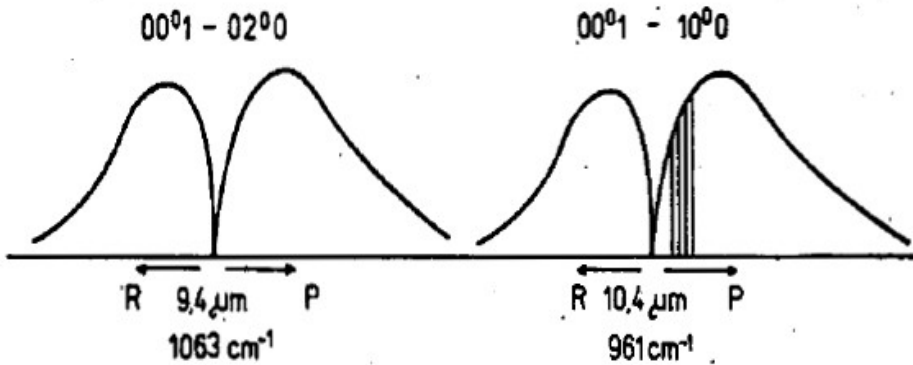




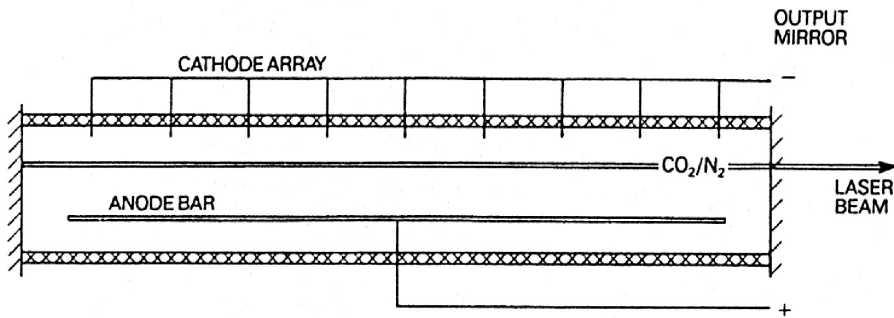
# Nitrogen laser



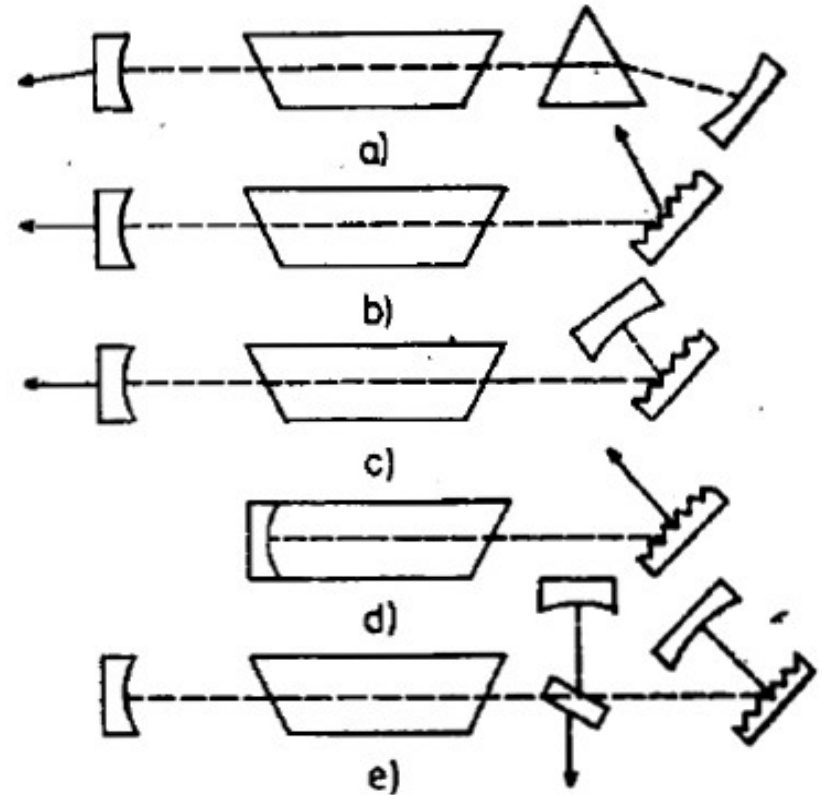
# CO2 laser



## Waveform intensity of CO2 laser lines



Electrode arrangement in a transverse excitation atmospheric (TEA) carbon dioxide laser

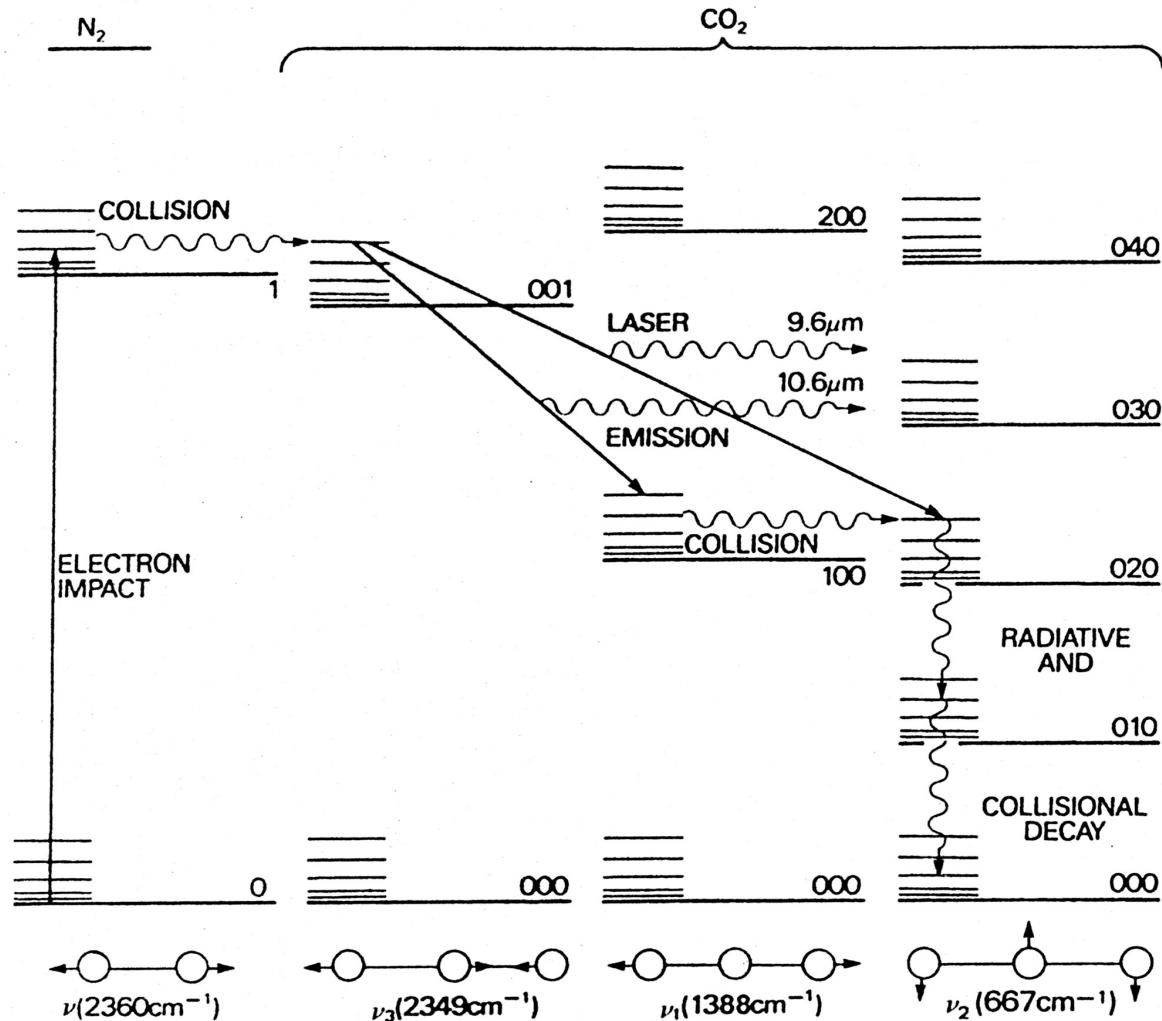


Selection of wavelengths and beam binding





# Energetics diagram of CO2 laser

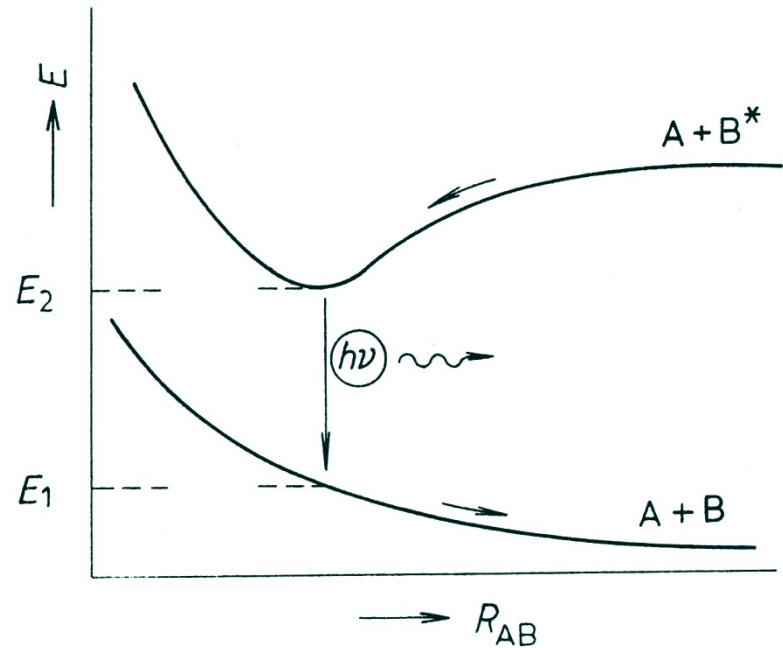


Energetics of the carbon dioxide laser. The rotational structure of the vibrational levels is shown only schematically



# Excimer

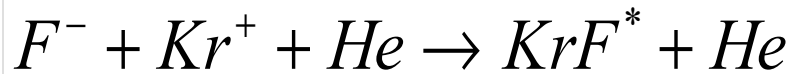
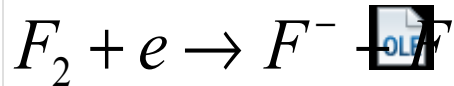
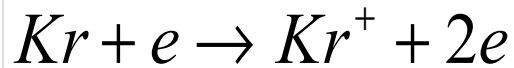
- Excimer - an unstable molecule formed for a limited time due to the action of an excited atom (molecule) with an atom (molecule) in its base ground state.
- Dissociation occurs within 10-14s after excimer transition to base ground state (photon emission).



Dependence of the potential energy  $E$  of an excimer-forming system of atoms (molecules)  $A$ ,  $B$  on their  $R_{AB}$  distance



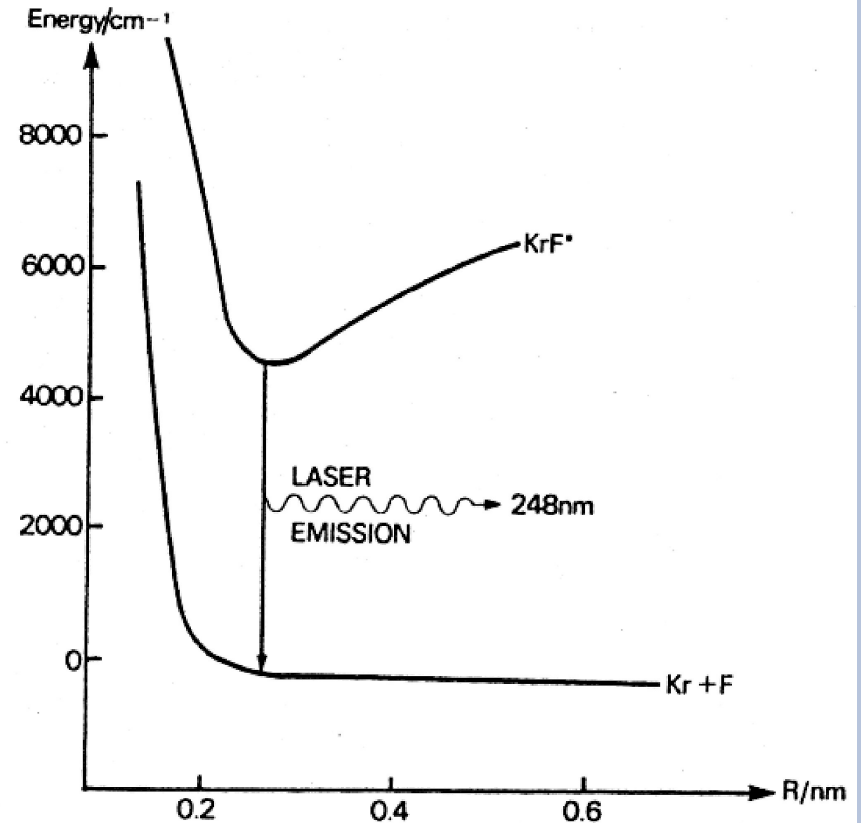
# KrF exciplex laser



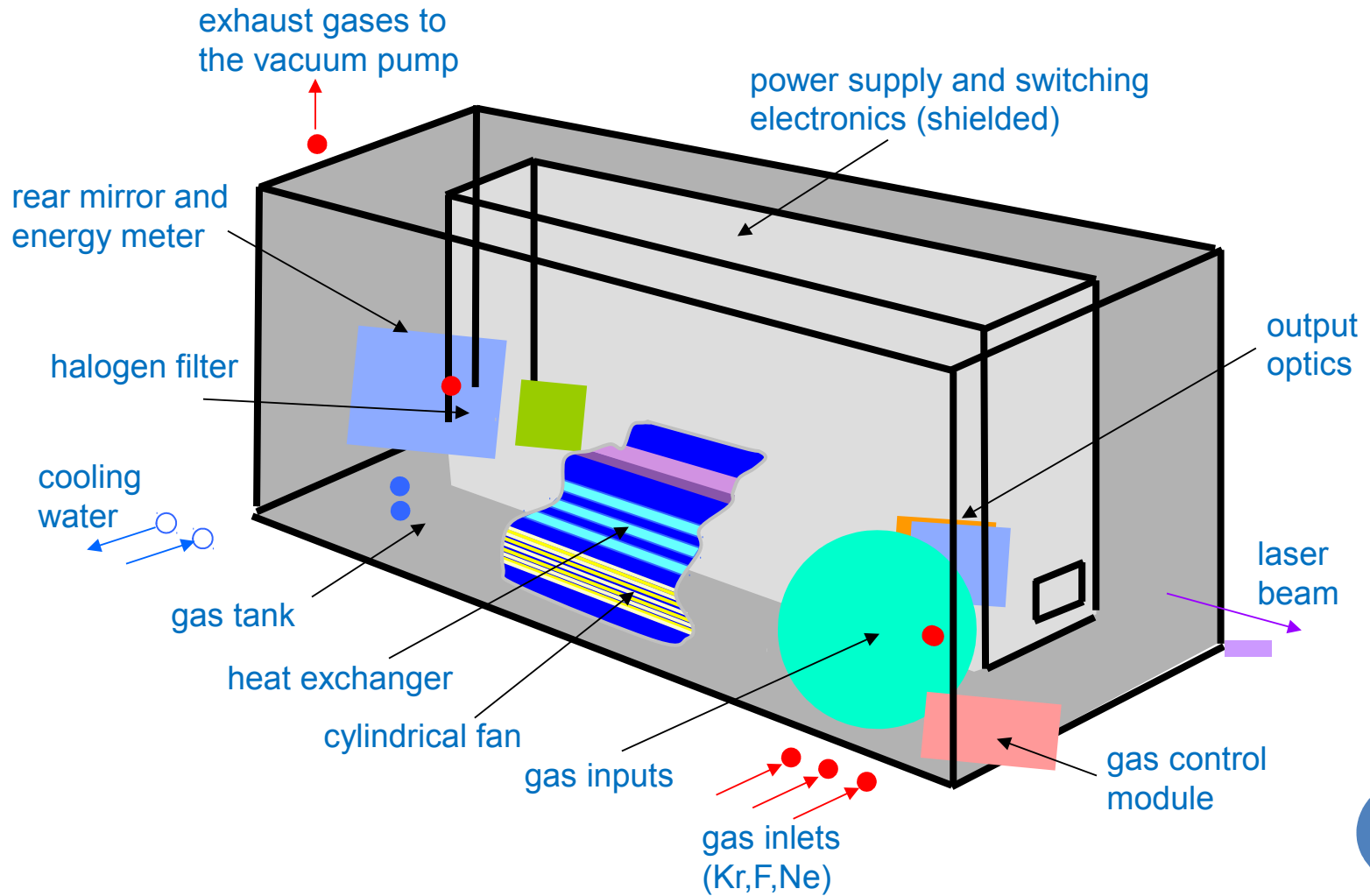
Exciplex - excited complex

Excimer - excited dimer

Realization : 1970 Basov, Xe2\*  
excited electrons

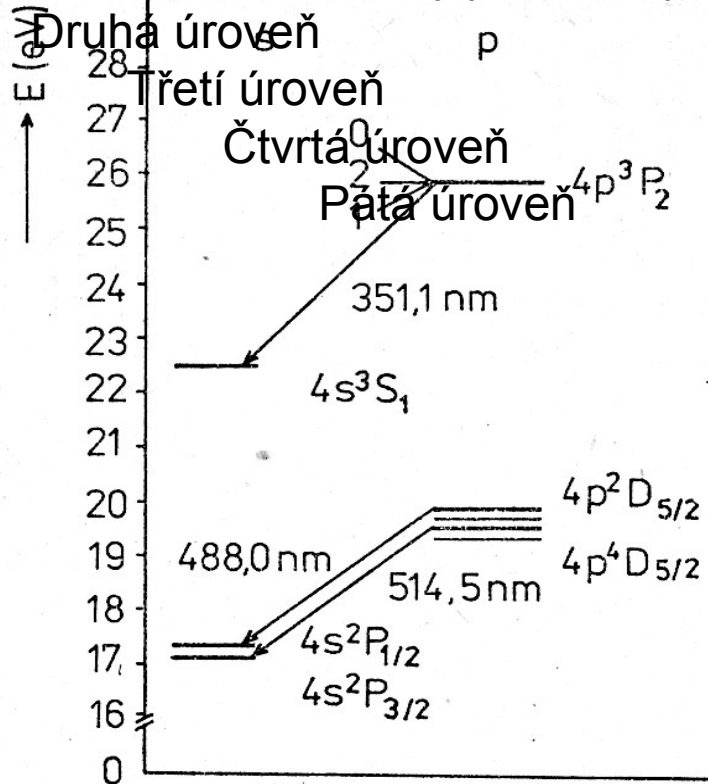


# Excimer laser

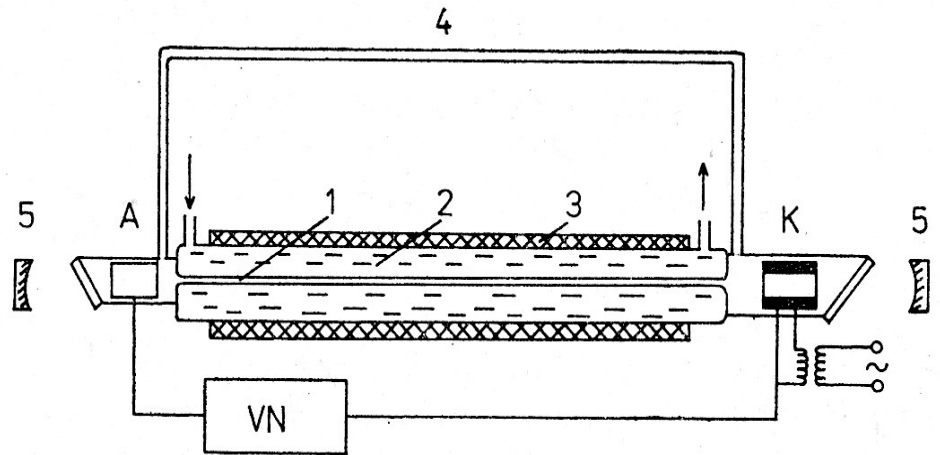


# Argon ion laser

Klepnutím lze upravit styly předlohy text



Scheme of transitions ArII and ArIII ions

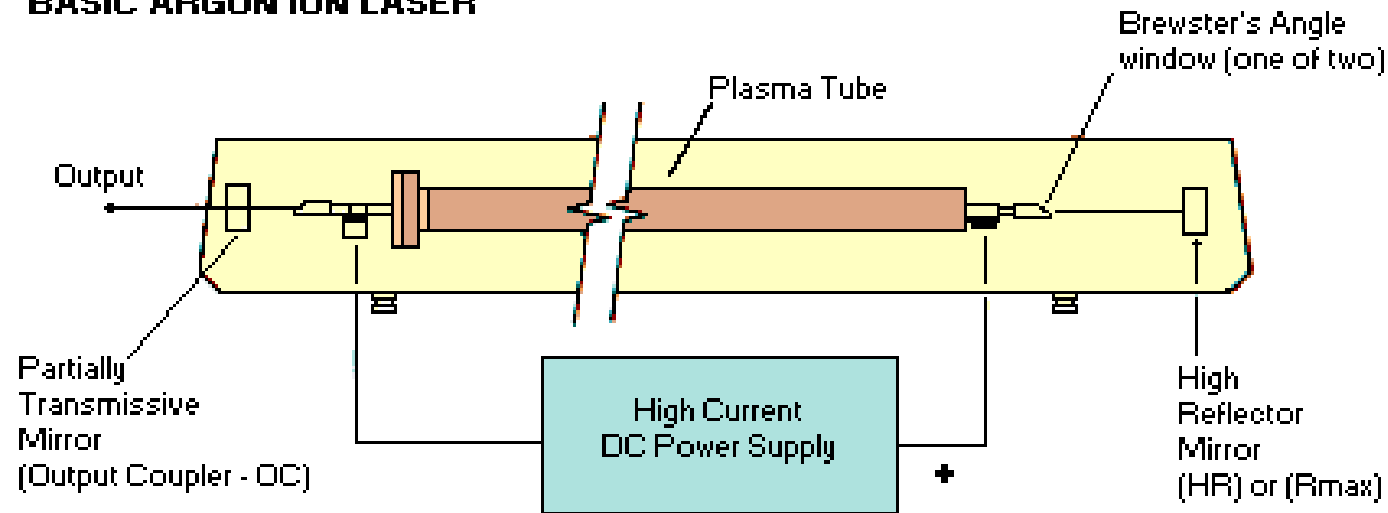


Argon laser scheme:  
 1-high current capillary  
 2-water cooling  
 3-solenoid  
 4-level capillary  
 5-mirrors  
 A-anode, K-hot cathode,  
 VN - high voltage



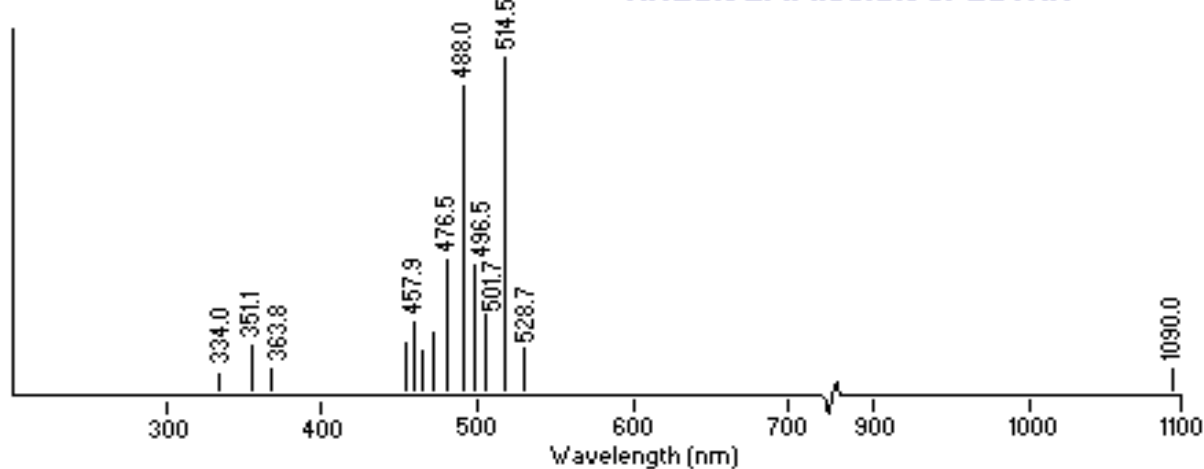
# Argon ion laser

## BASIC ARGON ION LASER



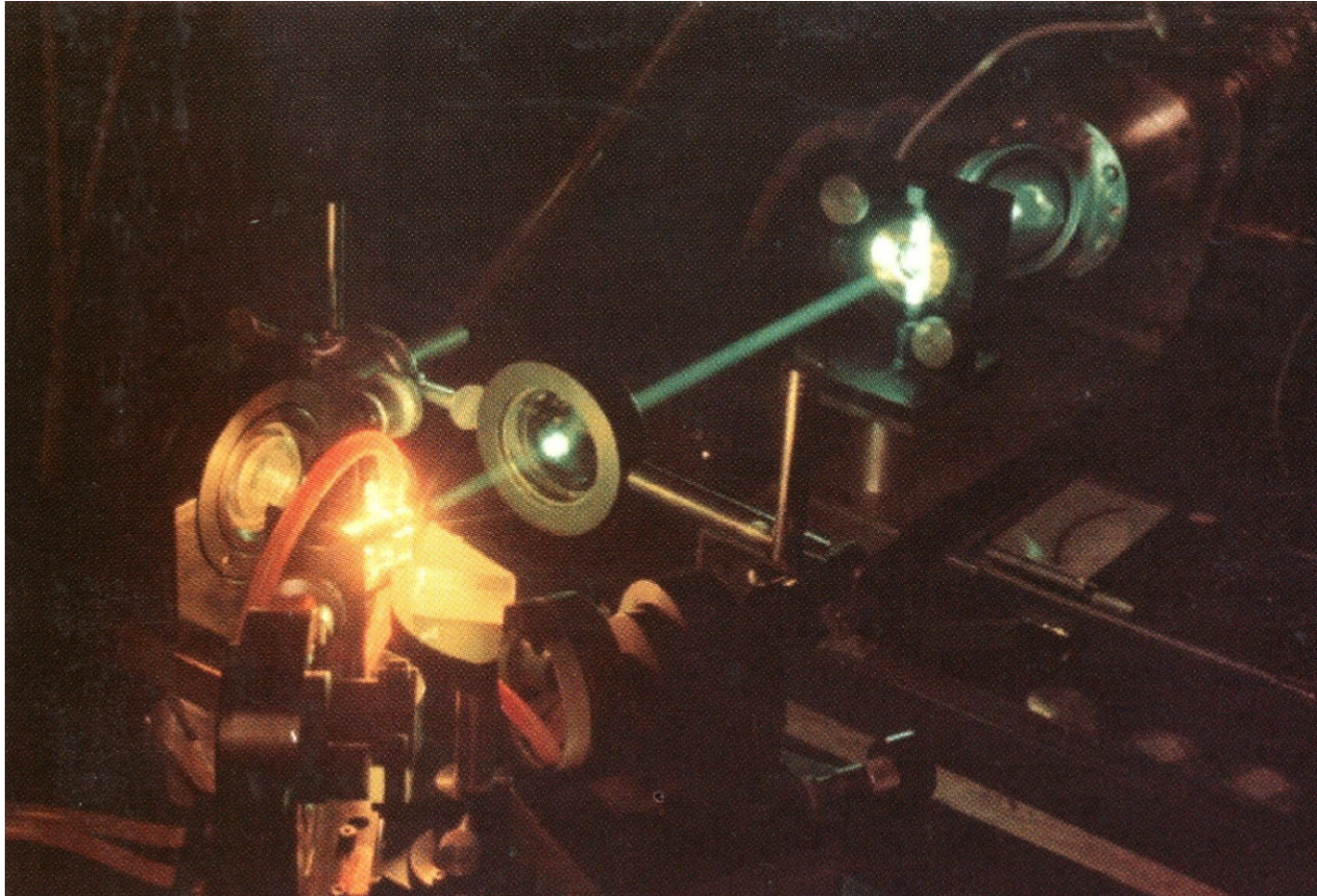
RELATIVE LASER OUTPUT POWER

## ARGON EMISSION SPECTRA





# Dye lasers

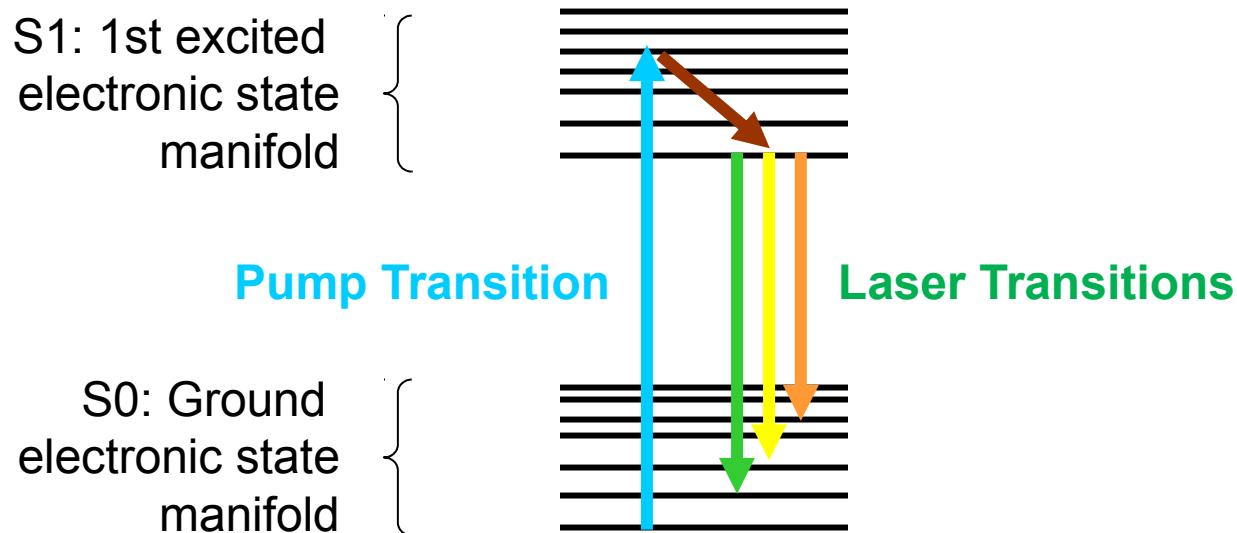


Dye lasers are an ideal four-level system, and a given dye will laser over a range of  $\sim 100$  nm.



# A dye's energy levels

- The lower laser level can be almost any level in the S0 manifold.



Dyes are so ideal that it's often difficult to stop them from lasing in all directions!

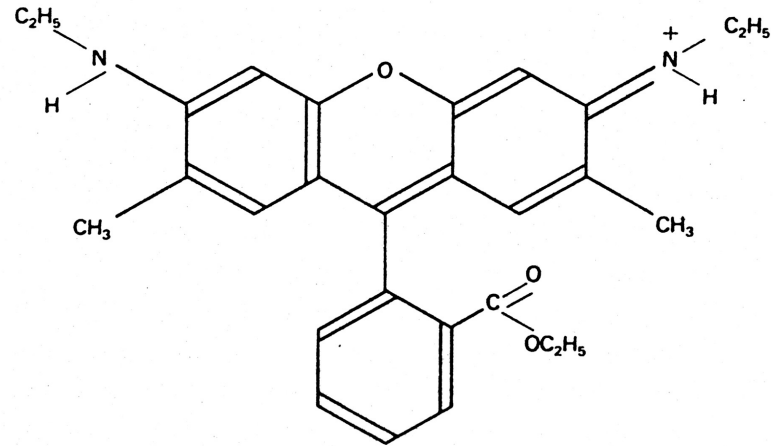




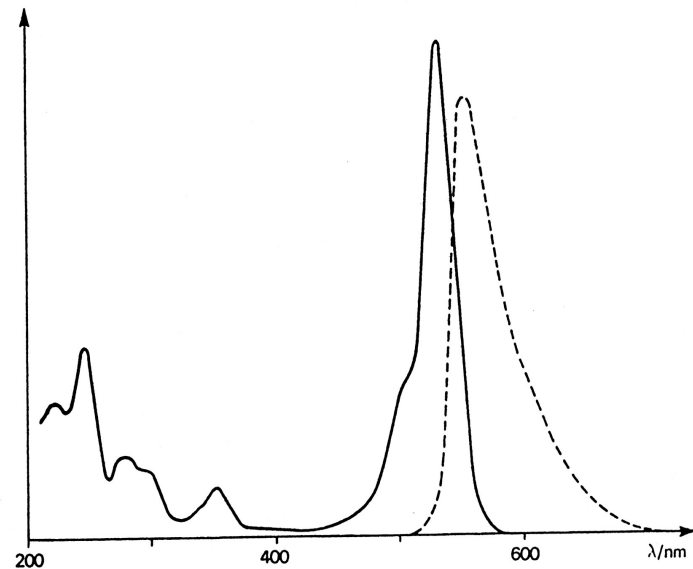
# Dye lasers

- They are characterized by high spectral bandwidth gain (10 -100 nm) and it follows:

1. Possibility of continuous change of the wavelength of the laser radiation in the range of sufficient gain bandwidth
2. Possibility to generate short pulses, up to 1 ps



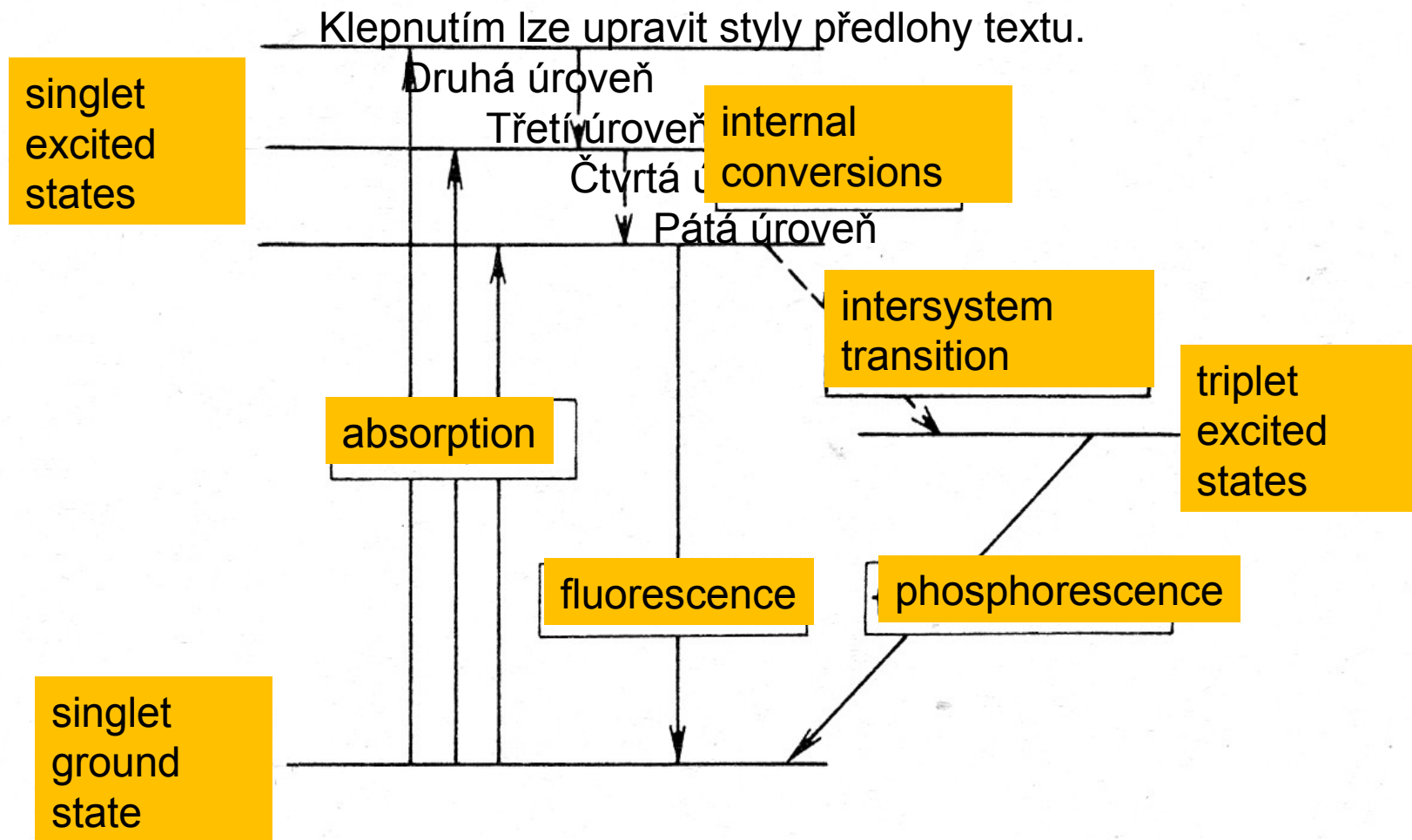
One of the resonance structures of the Rhodamine 6G cation



Solution spectra of Rhodamine 6G in ethanol. The solid curve shows the absorption, and the dotted curve the fluorescence at longer wavelength

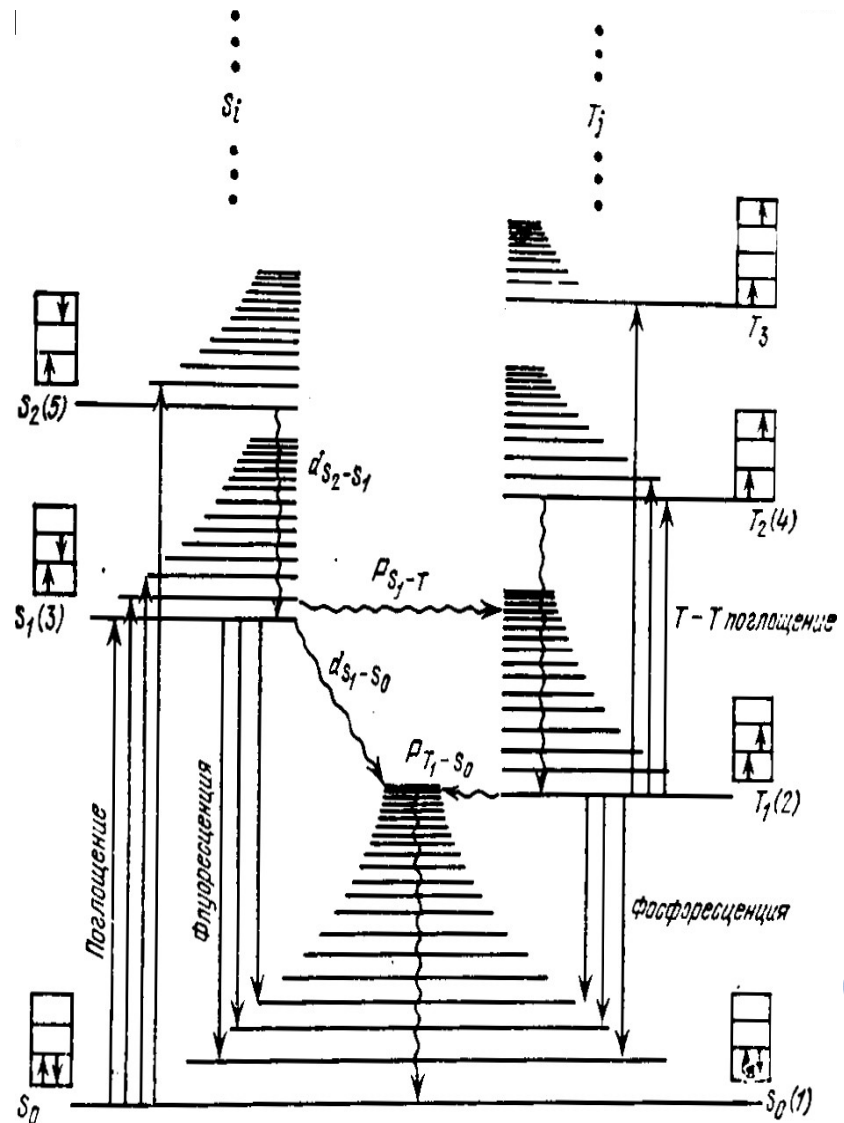


# Generalized Jablonski diagram of energy levels and dye transitions

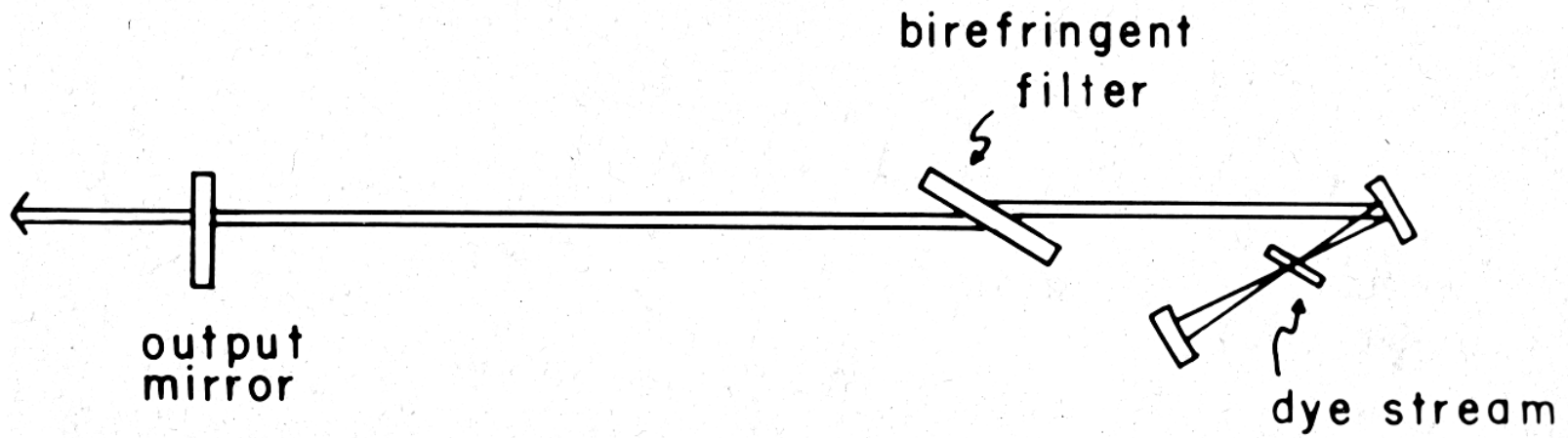


# Dye energy system

- Excitation by absorption of radiation by transition from ground to first singlet state
- Fluorescent transition to the ground state (population inversion possible)
- Non-radiative transition from S1 to metastable triplet T1 state (parasitic process)
- Absorption of fluorescent radiation by T1-T2 (T3) transition - quenches the fluorescence, decreases the gain of the active medium

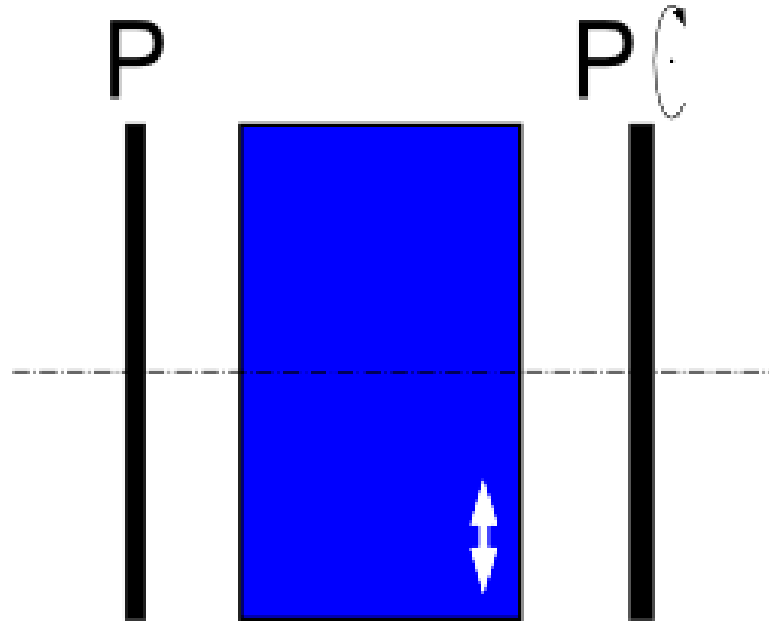


# Dye laser - principle



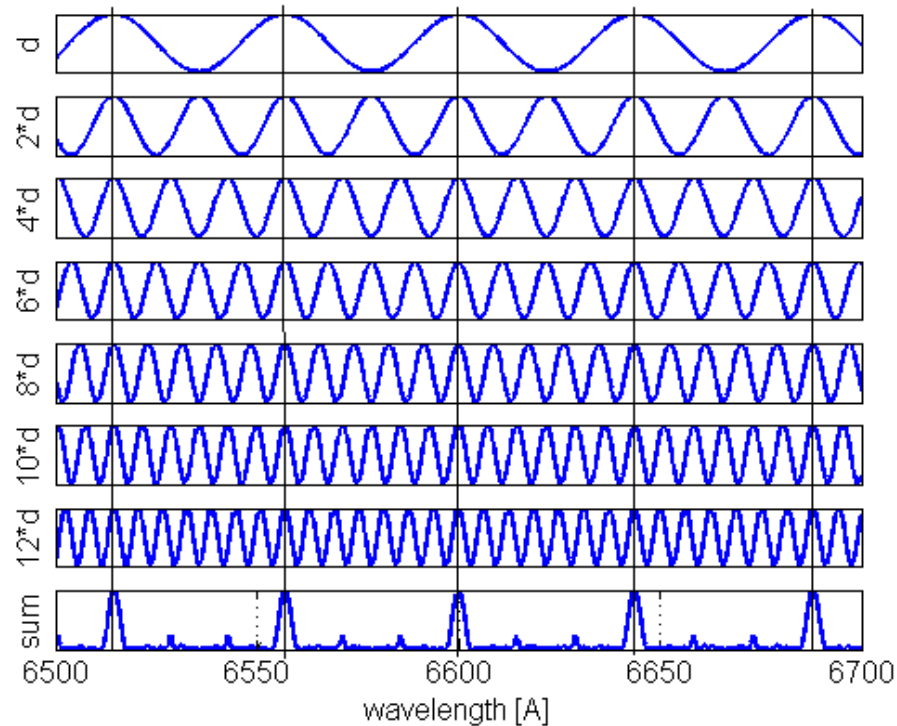
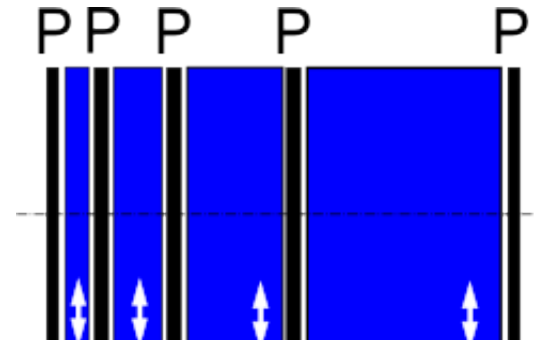
# Wood birefringent filters

- Wood birefringent filter consists of two polarizers and a crystalline quartz plate cutted parallel with crystal axis. The thickness of the plate depends on wavelenghts we want to separate. For already given example of Sodium doublet it gives the thickness of approximately 31.8 mm (depends also on operated temperature). These types of filters are very exact optical devices and it is necessary to hold very exact manufacture thickness tolerances.



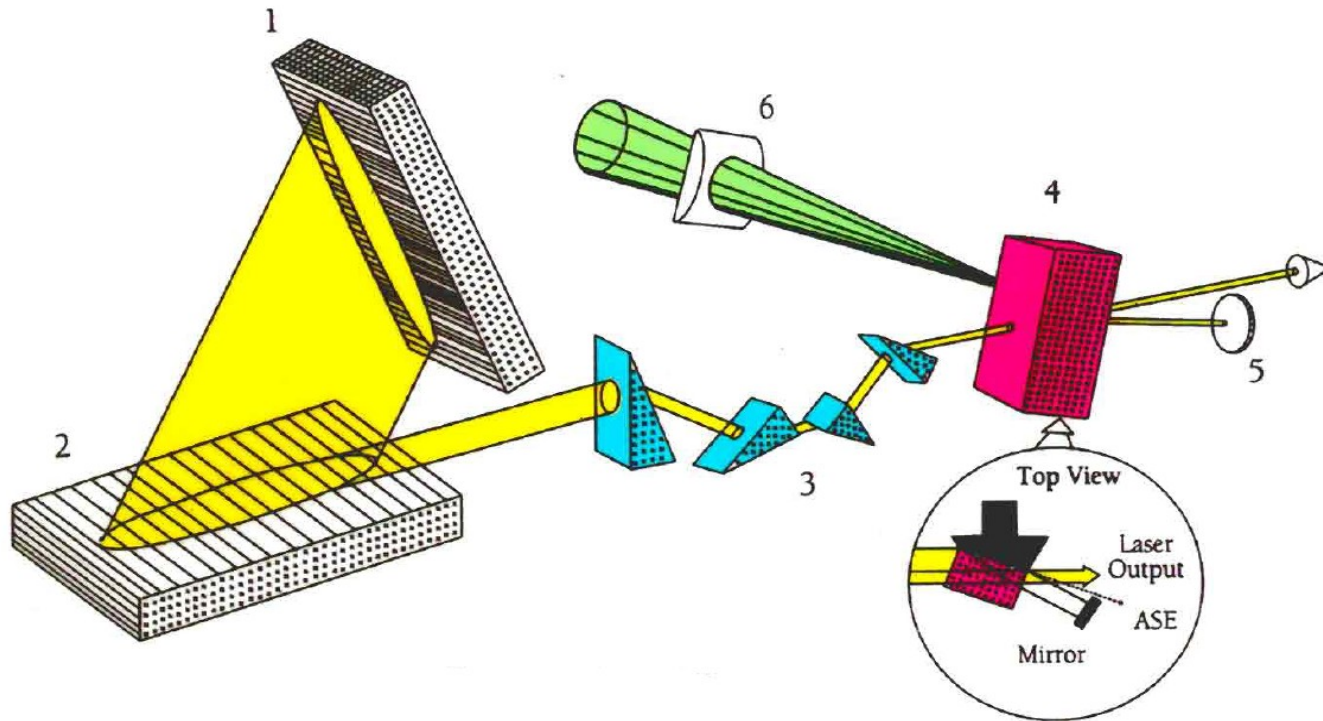
# Lyot birefringent filters

- This filter is in fact constructed from several Wood filters serially lined up. The thickness 'd' of the first plate is such that transmits requested wavelength and provides requested performance of the filter. Each next birefringent plate has a double width of the previous one. That provides two facts; firstly, the requested wavelength is transmitted and secondly, the unwanted transmitted wavelengths of a previous birefringent plate are filtered out. Such a cascade of birefringent plates sandwiched between polarizers provides high performance filter with a half-width in order of  $1/10$  nanometers

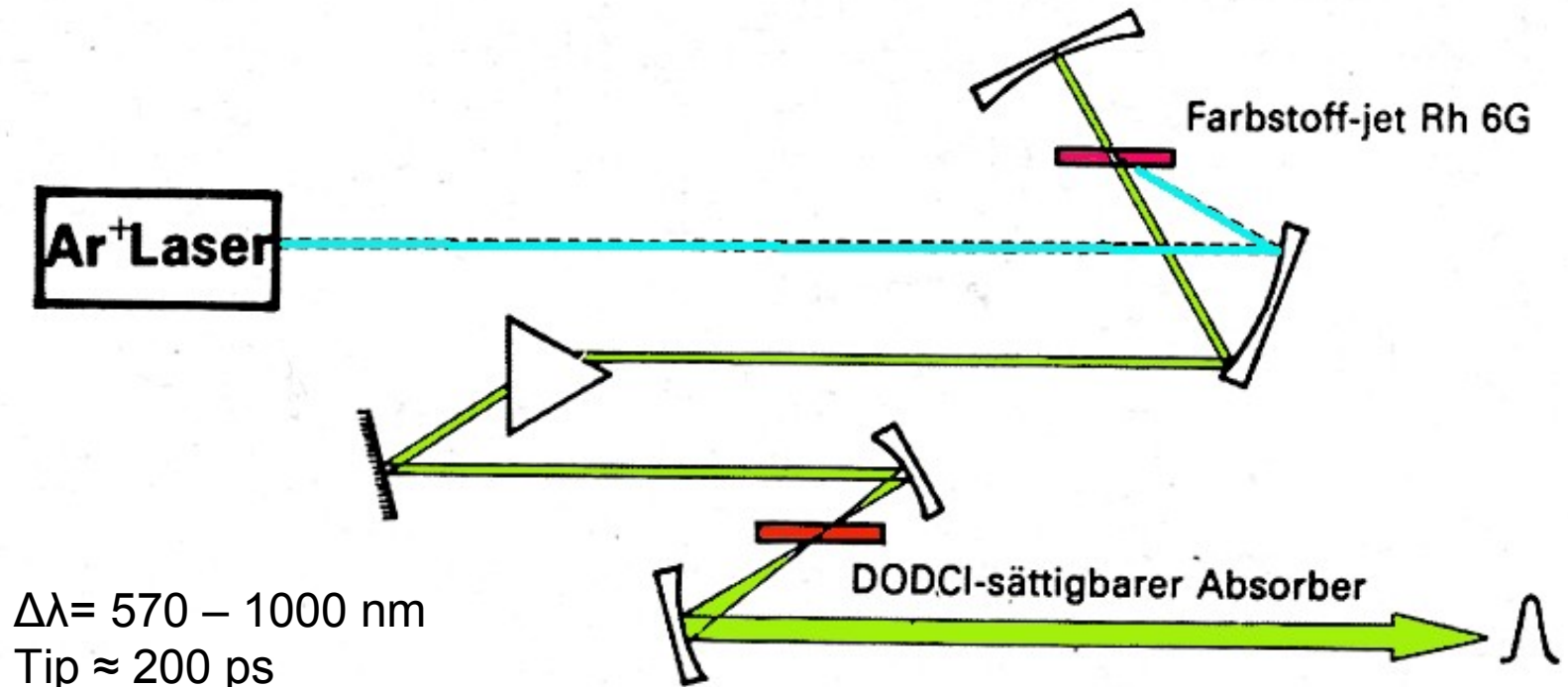


# Dye laser

1. Mirror (tuning)
2. Grating
3. Beam expander
4. cuvette with dye
5. Resonator mirror
6. Pumping (by laser)



# Picosecond laser

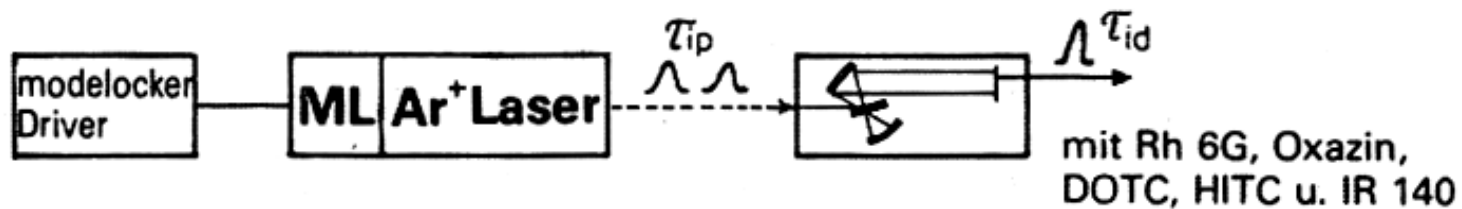


$\Delta\lambda = 570 - 1000 \text{ nm}$

Tip  $\approx 200 \text{ ps}$

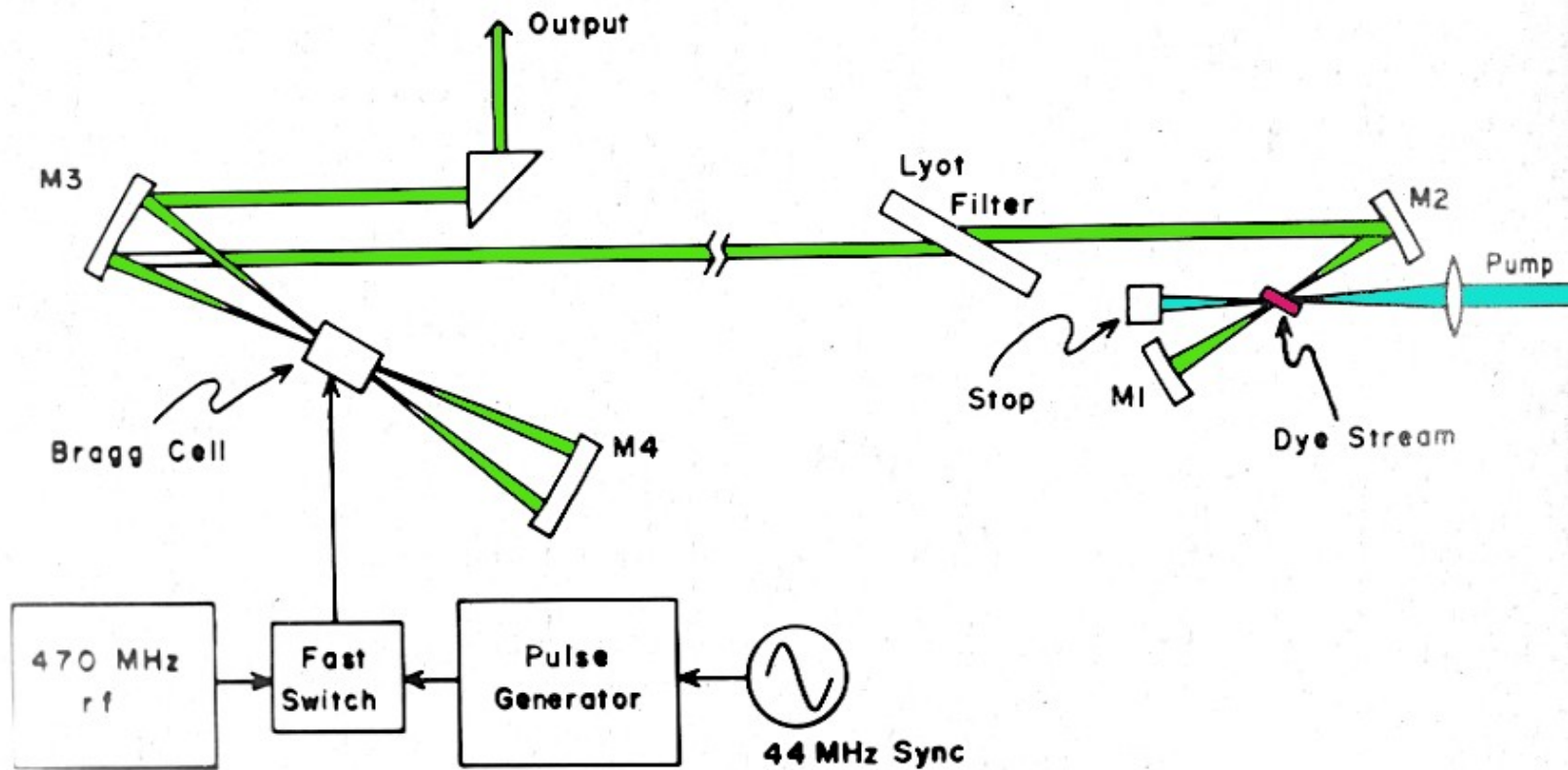
$\tau_{id} \approx 0,8 - 50 \text{ ps}$

$f = 50 - 150 \text{ MHz}$





# Mode-locked dye laser



# Types of dyes for lasers

## Emission range (nm)

340 - 430

360 - 480

410 - 440

440 - 520

460 - 540

510 - 700

540 - 1200

630 - 720

## Structural type

stilbenes

oxazoles

anthracenes

acridines

coumarins

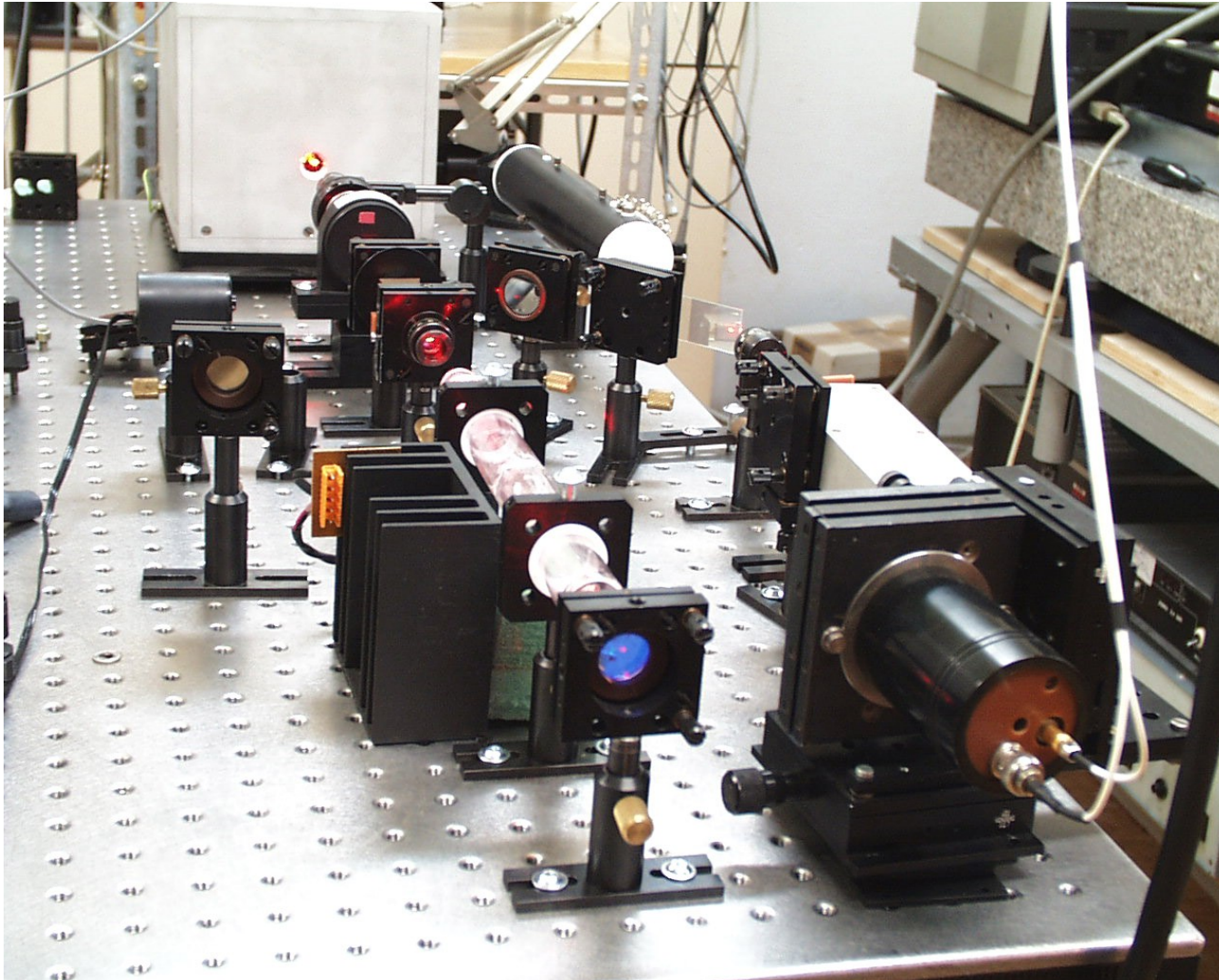
xanthenes

cyanines

oxazines



# Experimental laser workplace



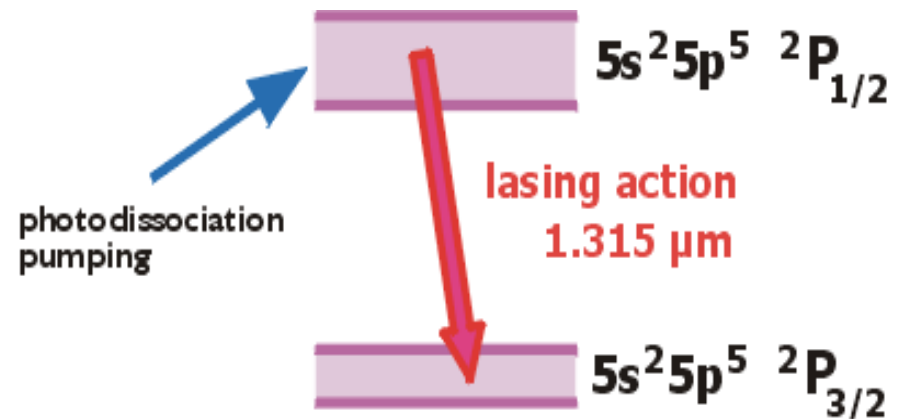
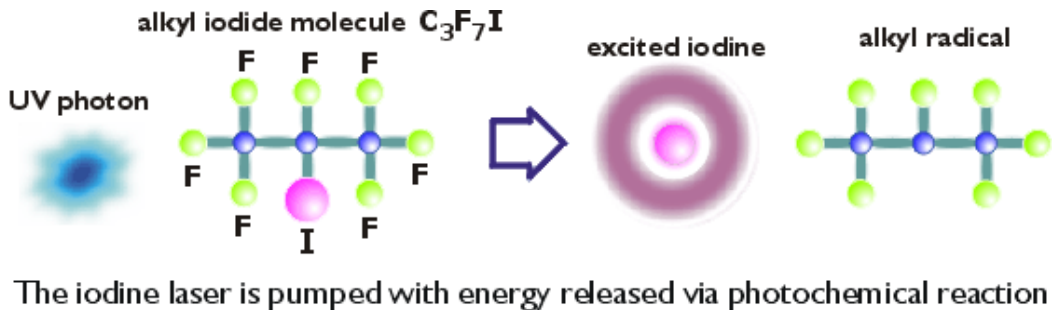
# Prague Asterix Laser System

- The backbone of the PALS Research Center is the giant iodine laser system. In its current configuration and at a base wavelength of 1315 nm, it is capable of providing pulses of up to 1 kJ in the main laser beam, in addition to 100 J in two smaller additional beams. The wavelength of the laser beams can be converted to a second (658 nm, red) or a third (438 nm, blue) harmonic base frequency. Due to the very short laser pulse length (approx. 350 ps), the peak laser pulse power is enormous - up to 3 TW, i.e. 3 million megawatts. The laser is able to deliver such a giant pulse about once every half hour. The PALS output beam is high quality: spatially homogeneous, and stable. Its energy practically does not change from shot to shot.



# Asterix iodine laser

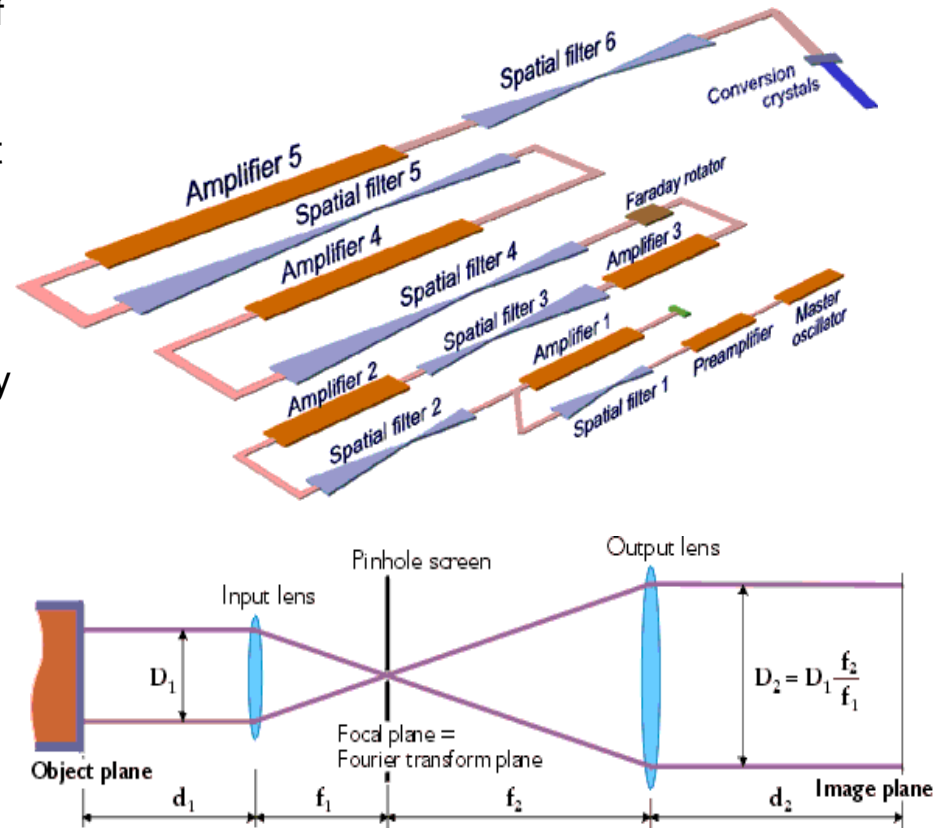
Asterix IV is a gas laser in which iodine atoms are used to generate near-infrared radiation at a wavelength of 1,315  $\mu\text{m}$ . The iodine atom is obtained from the parent molecule of the alkyl iodide  $\text{C}_3\text{F}_7\text{I}$  by photodissociation. The atom is released from the chemical bond by means of pulsed UV radiation delivered by the lamps. The electron envelope of iodine emerging from the photodissociation reaction is excited, thereby automatically forming a population inversion with respect to the underlying state. This creates the conditions for the laser action.





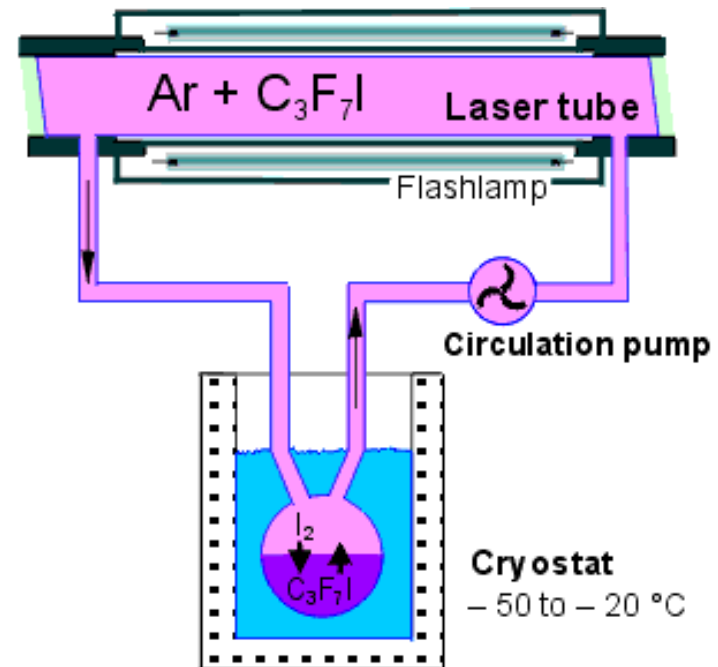
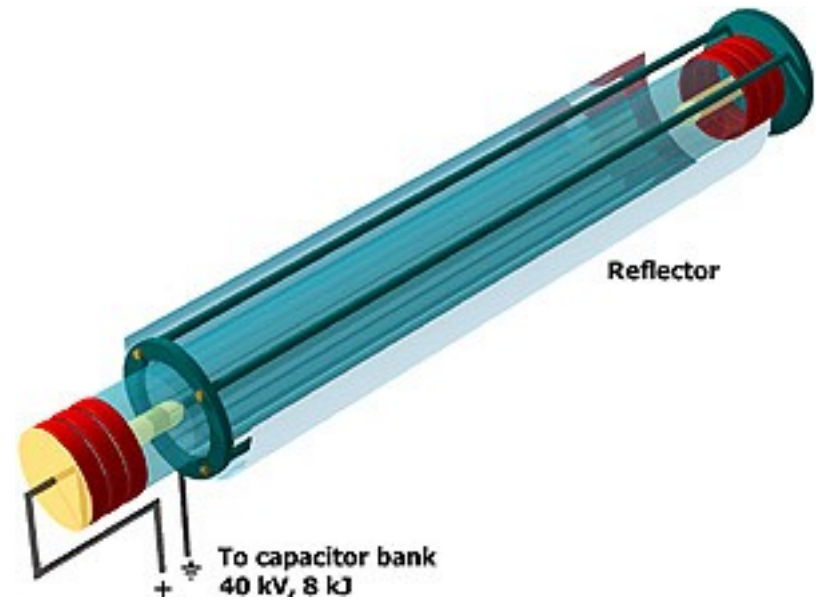
# The overall arrangement PALS

PALS is a single-beam laser system, consisting of an oscillator section generating an initial low light pulse and a string of five laser amplifiers that gradually amplify the pulse. Such an arrangement scheme is called "master oscillator - power amplifiers" (MOPA), or control oscillator - power amplifiers. The size of the amplifiers increases from one amplification stage to the next, so that the diameter of the amplified laser beam gradually increases, from the initial 8 mm to the final 290 mm. In this way, the power density of the laser beam is maintained at a value at which the surface of the individual optical elements cannot yet be damaged by excessive light load.



# Optical amplifier

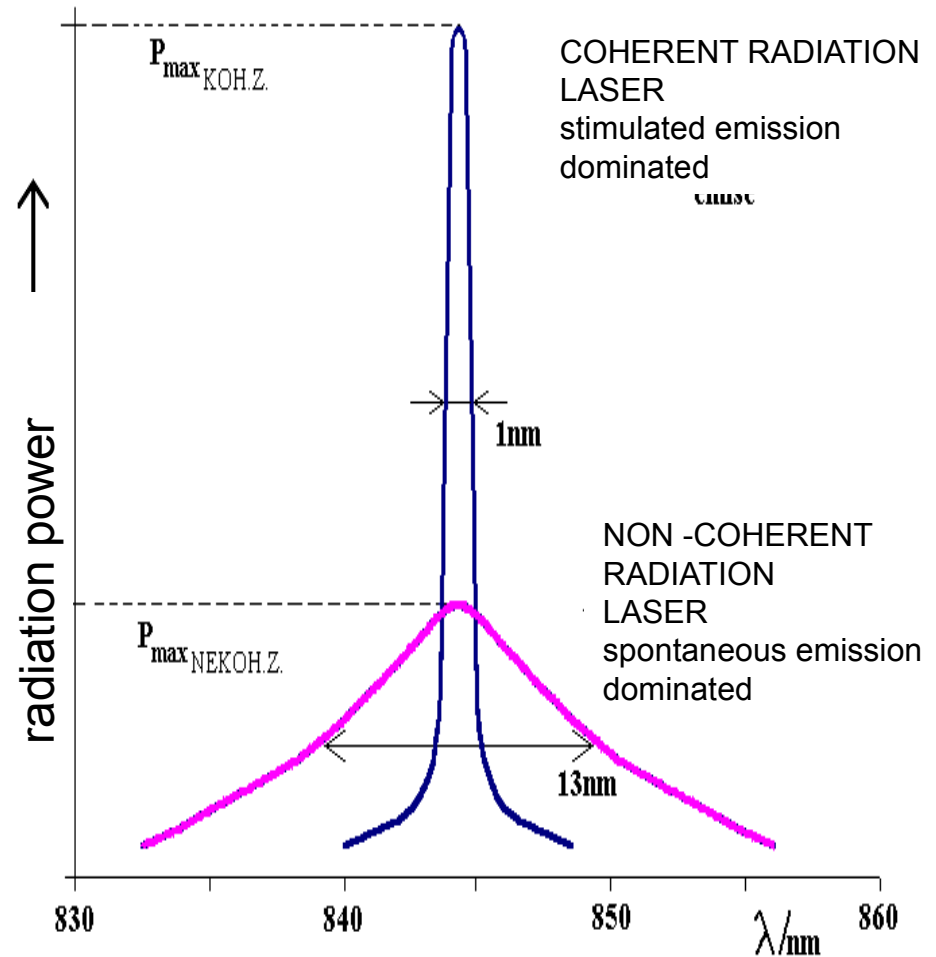
The Asterix IV / PALS laser chain includes a total of five power amplifiers. Their task is to amplify pulses coming from the oscillator to energy up to one kilojoule. The size of each amplifier gradually increases towards the end of the string - the final fifth amplifier is over 13 m long (see picture) and provides a laser beam with a diameter of 29 cm. In the split of second before the actual laser shot, the amplifiers are "activated" by discharging large capacitor batteries into the lamps surrounding the amplifier cells containing the gaseous working environment. The intense pulse of incoherent ultraviolet radiation produced by the lamps gives rise to a large number of excited iodine atoms in the cuvettes, which are "ready" to deliver their excess energy to the laser pulse coming from the oscillator.



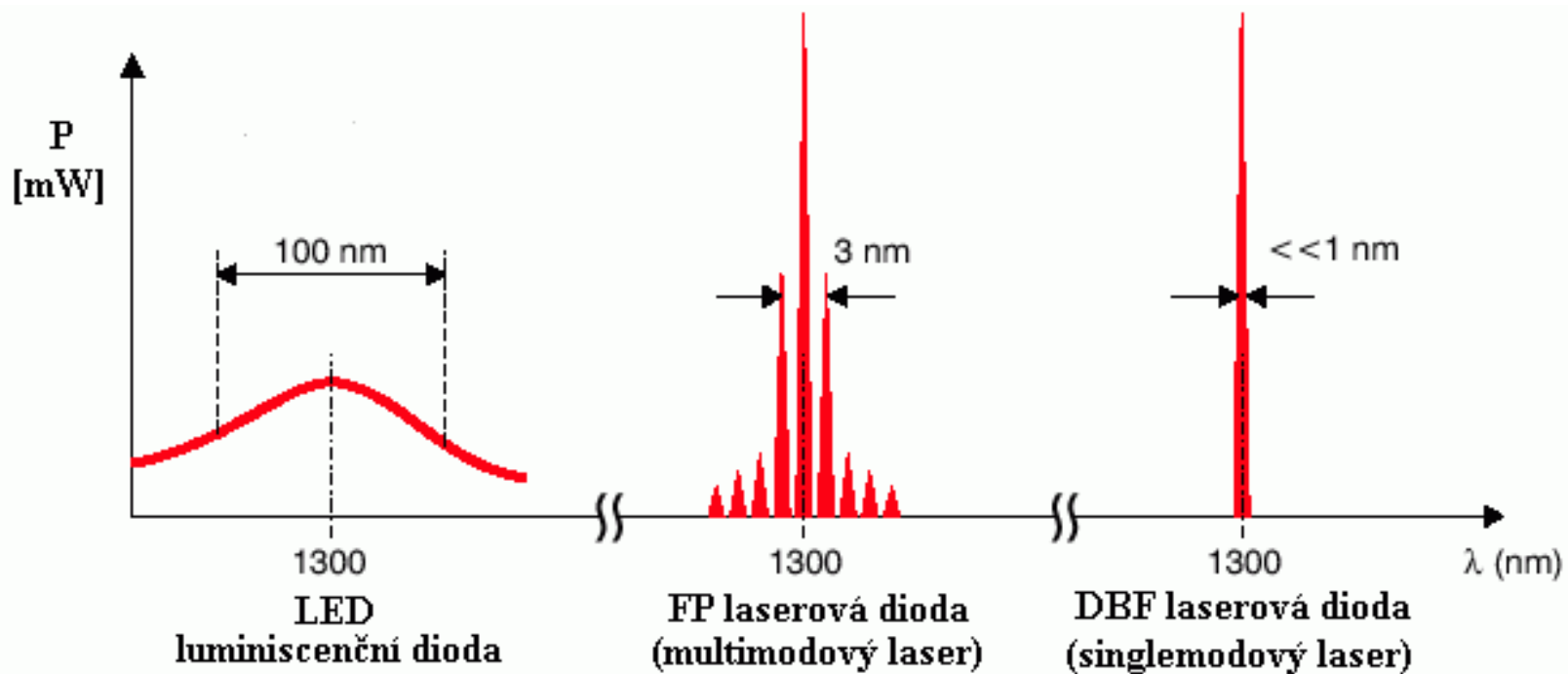


# Laser diodes

- For small currents, LED radiation is spontaneous and is a linear function of the excitation current. Upon reaching the threshold current, the power of the stimulated radiation increases sharply and coherent radiation is again emitted from the resonator mirrors, again linearly dependent on the magnitude of the excitation current. At the same time, there is a qualitative change in the shape of the radiation characteristic of the laser diode, expressed by decreasing the angle of radiation in a plane perpendicular and parallel to the plane of the PN transition, as well as reducing the bandwidth of the emitted radiation.

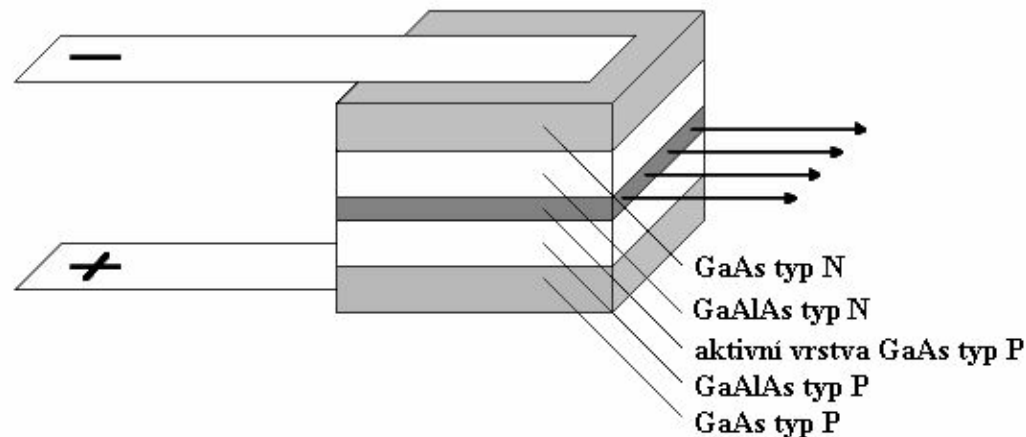


# Spectrum of LEDs and laser diodes



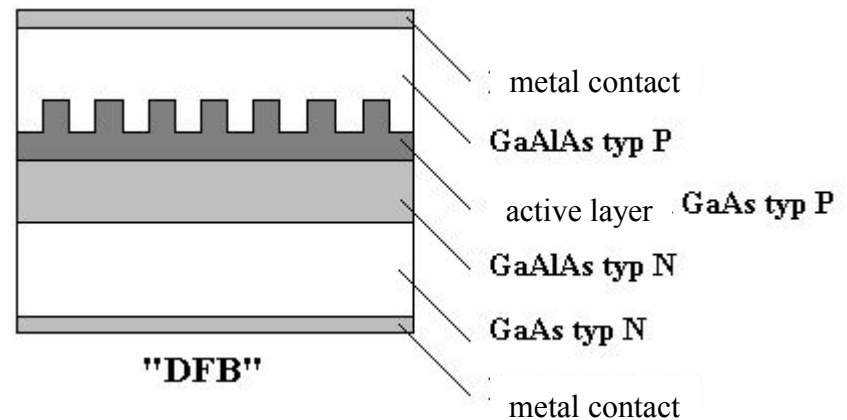
# Heterostructure lasers

- In these types of laser with *hetero-transitions*, the delimitation of the waveguide is determined by a sharp change of the refractive index in the region of *hetero-transition*. Along with efficient light guidance, *heterostructure* provides conditions for effective concentration of minority carriers. The effect of hetero-transition concentrates radiation and injection carriers in selected areas.



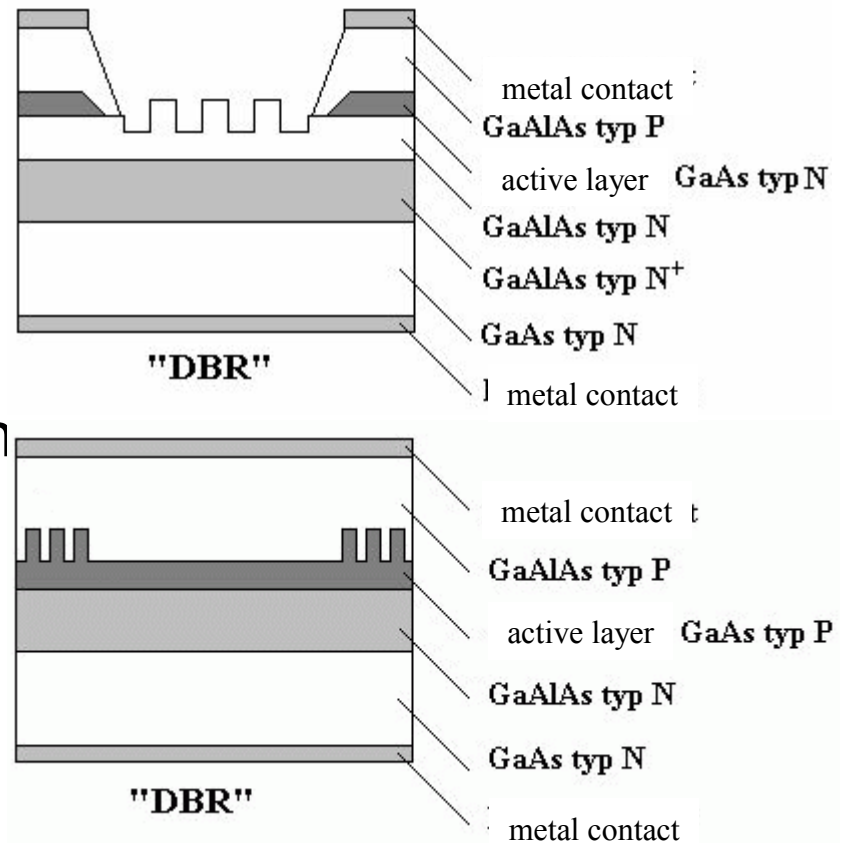
# Distributed Feed Back Lasers

- In this type of laser, the resonator is realized without mirrors using spatial periodic structures (diffraction gratings). The function is based on periodic change of refractive index in the direction of propagation. The feedback is created by the continuous binding of the propagating wave in the opposite direction by Bragg scattering. The grid is formed by etching directly on the surface of the active layer.
- These lasers are referred as DFB

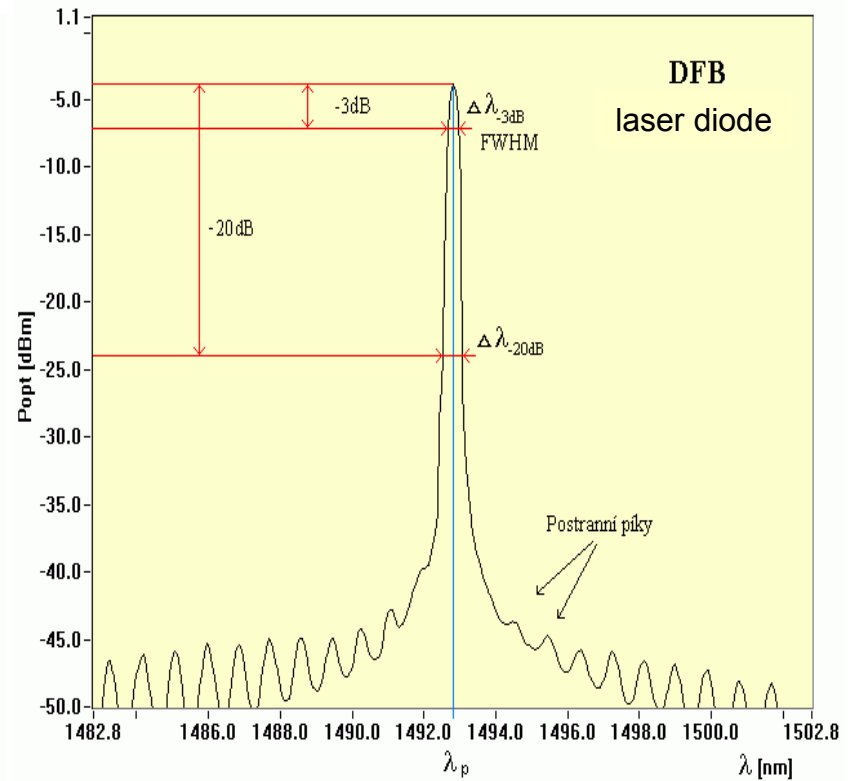
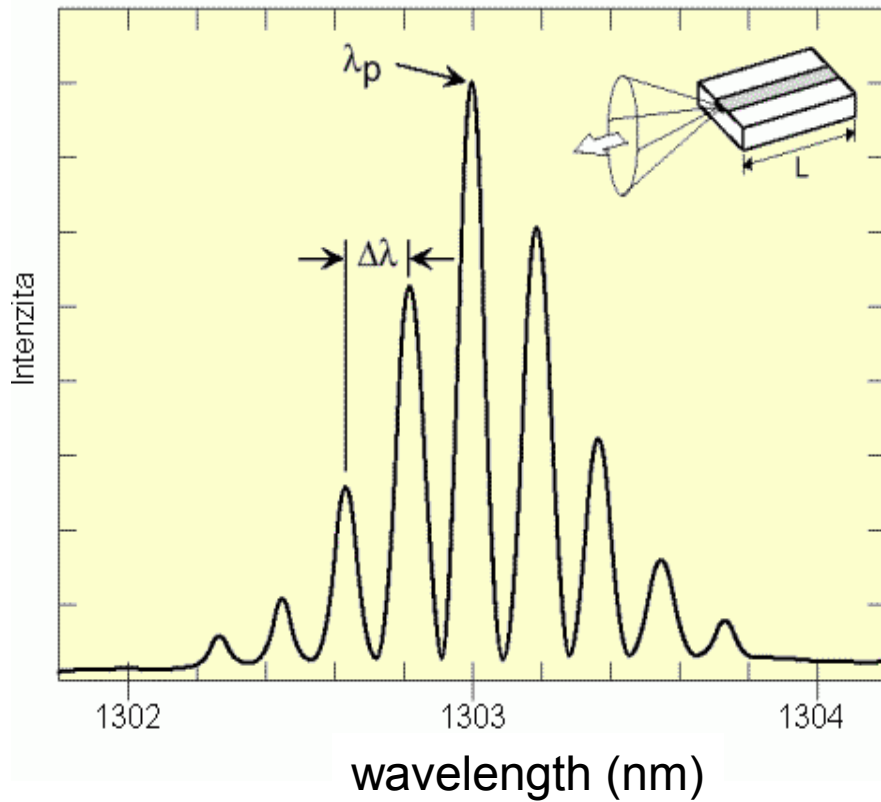


# Distributed Bragg Reflector Lasers

- Optical radiation generation and feedback (again using an optical gratings) take place in separate parts of the structure. Two types of construction are used, with one or two Bragg mirrors. In common practice, the type with two Bragg mirrors at the ends of the waveguide is more often used

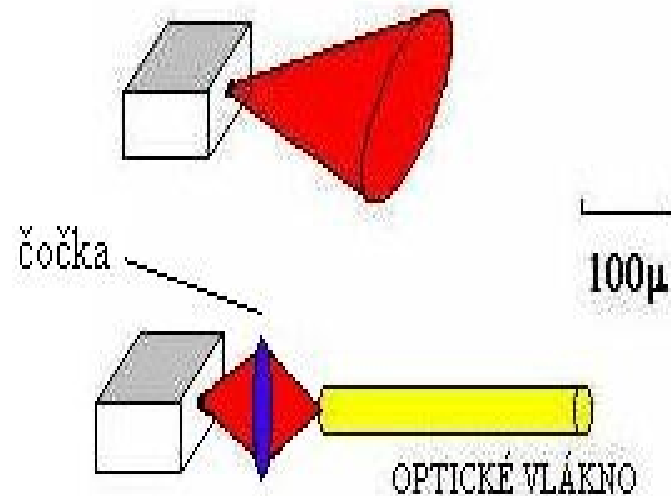


# Spectrum of laser diodes



# Edge-emitting lasers

- This type (Edge Emitting Lasers - EEL) emits radiation from the edge transition. Manufacturing and applications of laser diodes currently prevails



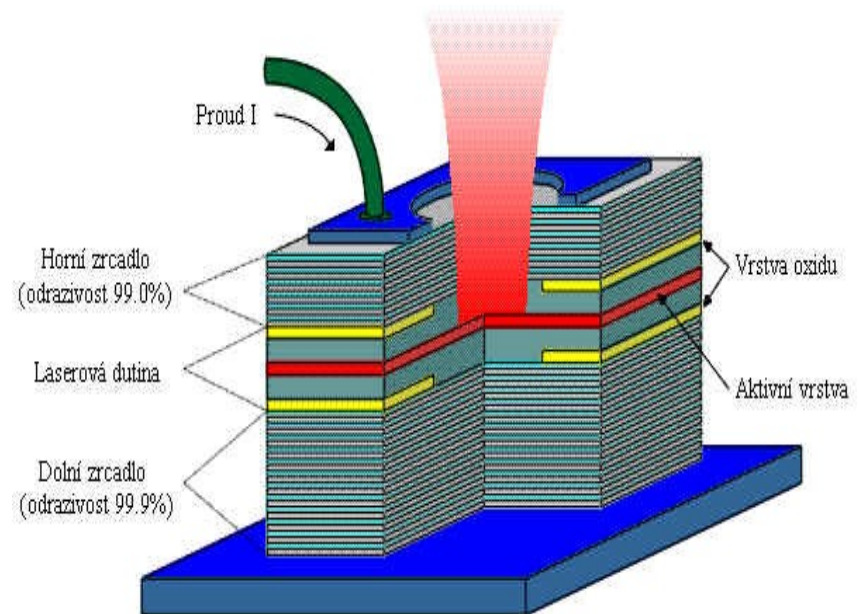
Edge Emitting LD





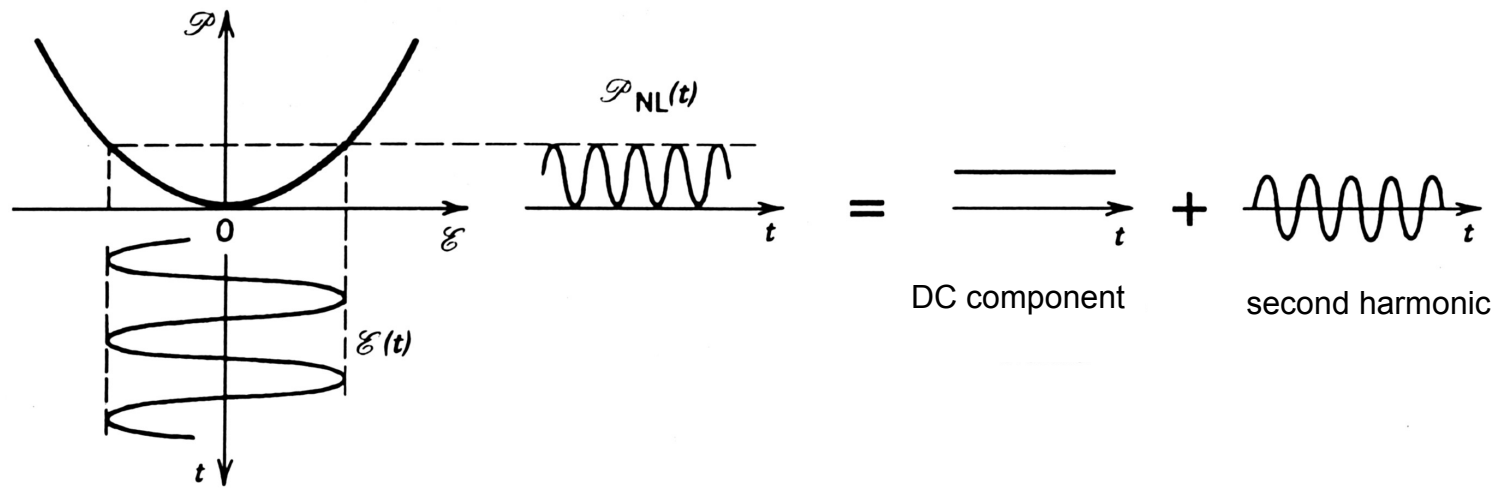
# Surface-emitting lasers

- VCSEL (Vertical Cavity Surface Emitting Lasers) emit radiation from the part surface parallel to the transition plane. Radiation emitted from the surface is absorbed by the substrate and lost or, more preferably, is reflected from the metal contact



# Frequency conversion

- The first option is to use nonlinear phenomena of the second (third) order. The intensity of the second harmonic radiation is proportional to the square of the coefficient of optical non-linearity and the intensity of the incident wave, inversely proportional to the square of the wavelength.

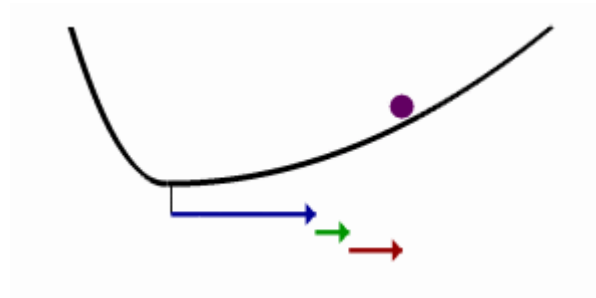
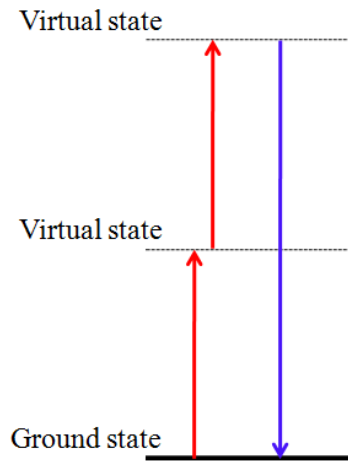
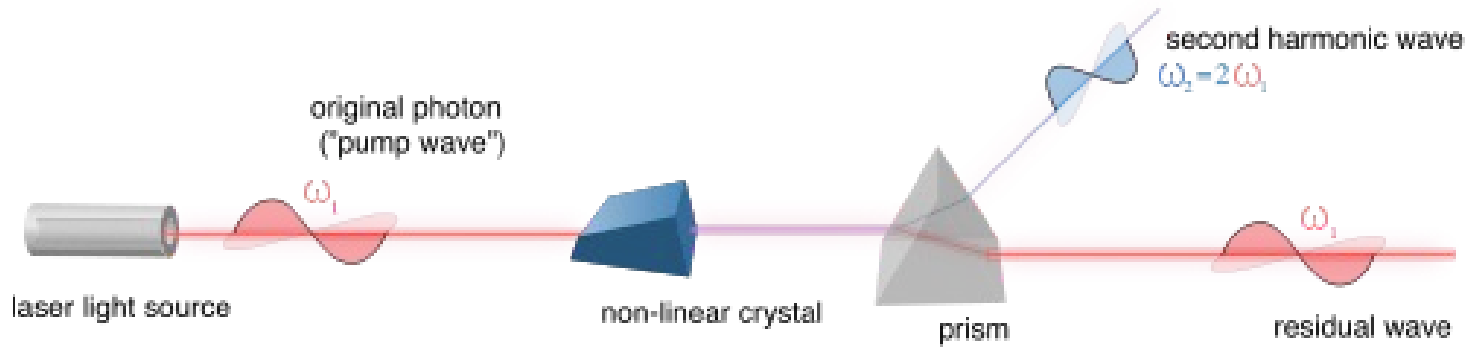


# Nonlinear optics - crystals

crystal	$\Delta\lambda$ ( $\mu\text{m}$ )	MW/cm <sup>2</sup>
KDP (potassium dihydrogen phosphate)	0,2-1,35	400
KDDP (deuterated KDP)	0,2-1,8	500
ADP (ammonium dihydrogen phosphate)	0,2-1,2	500
RDP (rubidium dihydrogen phosphate)	0,2-1,5	300
CDA (cesium dihydrogenarsenate)	0,26-1,6	500
LiIO <sub>3</sub>	0,3-4,5	60
LiNbO <sub>3</sub>	0,4-4,5	120
Ba <sub>2</sub> NaNb <sub>5</sub> O <sub>15</sub>	0,38-5	100
HIO <sub>3</sub>	0,4-1,3	100
BBO ( $\beta$ -BaB <sub>2</sub> O <sub>4</sub> )	0,2-1,5	400



# Frequency conversion



The blue arrow corresponds to ordinary (linear) susceptibility, the green arrow corresponds to second-harmonic generation, and the red arrow corresponds to optical rectification.



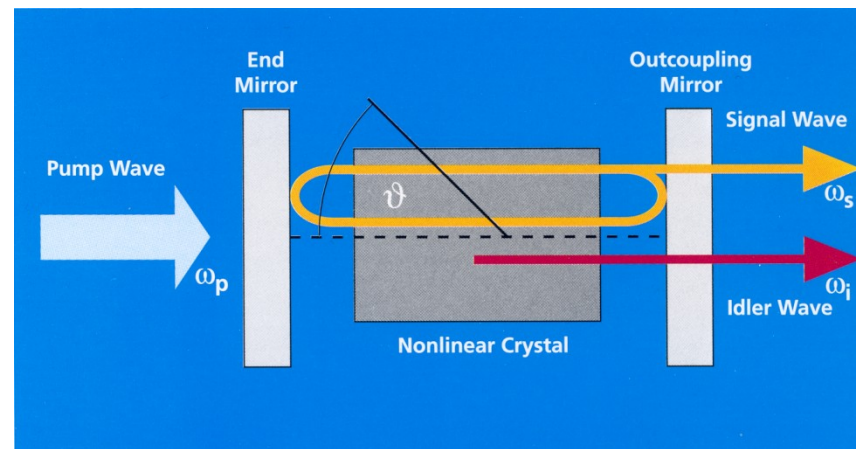
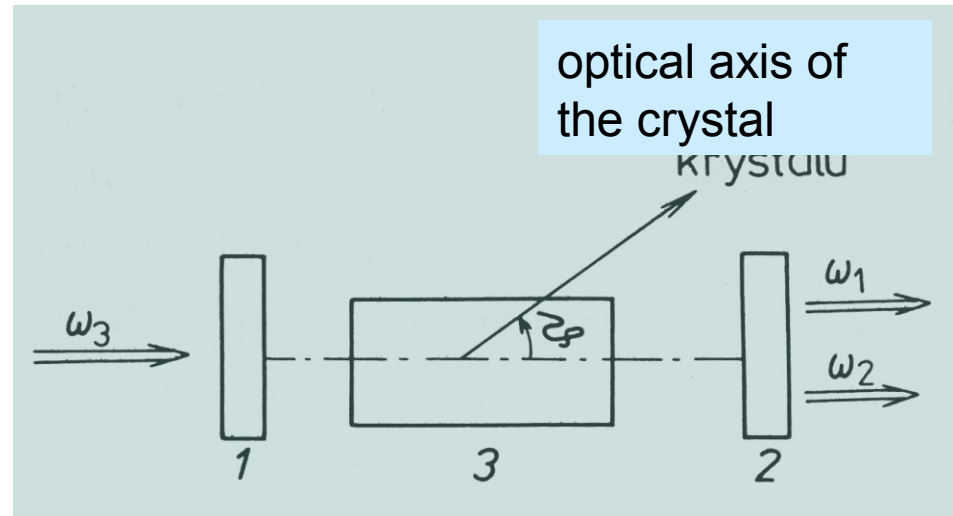
# Optical Parametric Oscillator (OPO)

- Based on a coherent decay photon of angular frequency  $\omega_3$  into two photons whose angular frequencies  $\omega_1$  and  $\omega_2$  (signal wave and idler wave), whereby the following applies:

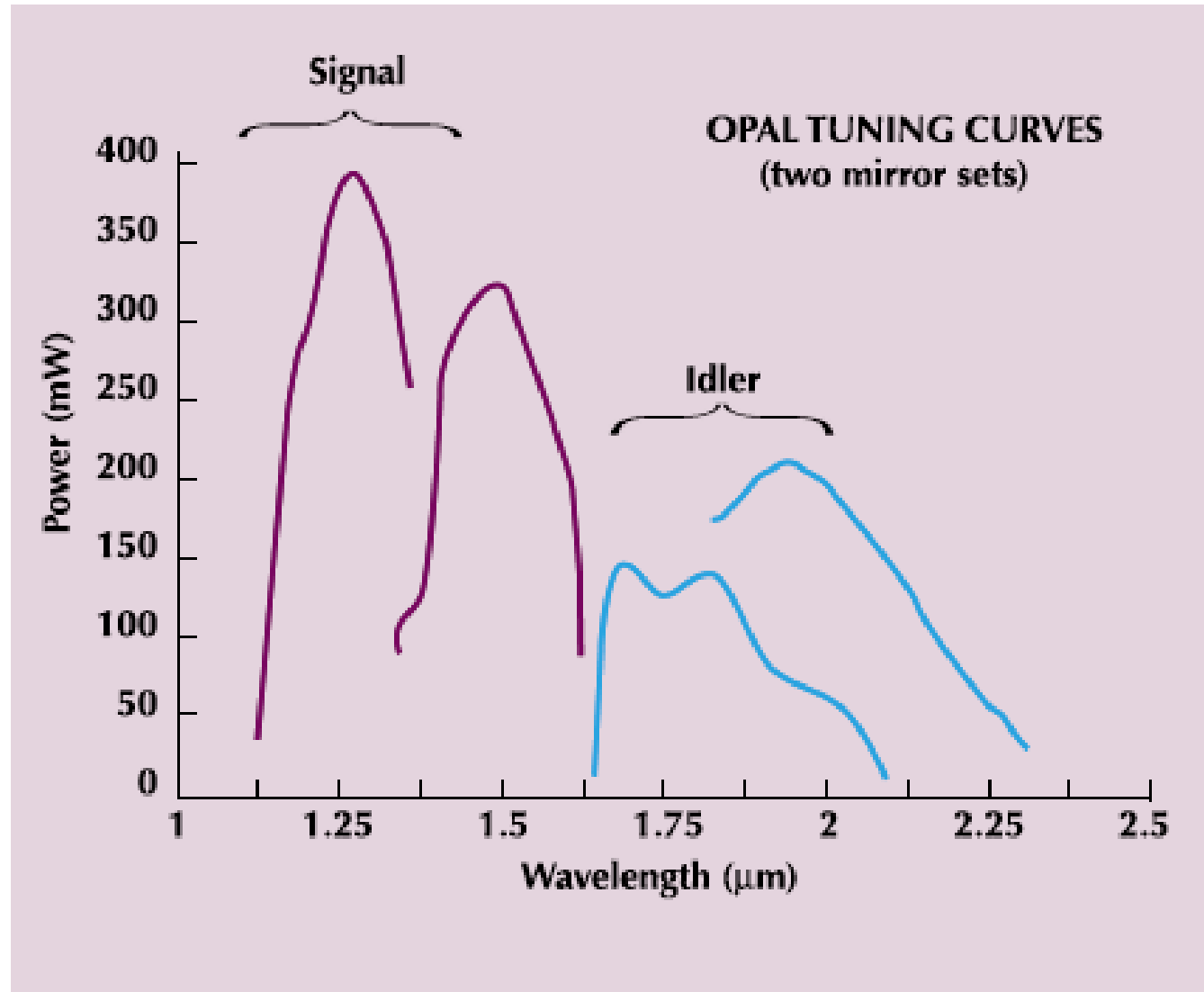
$$\omega_3 = \omega_1 + \omega_2$$

and ratio

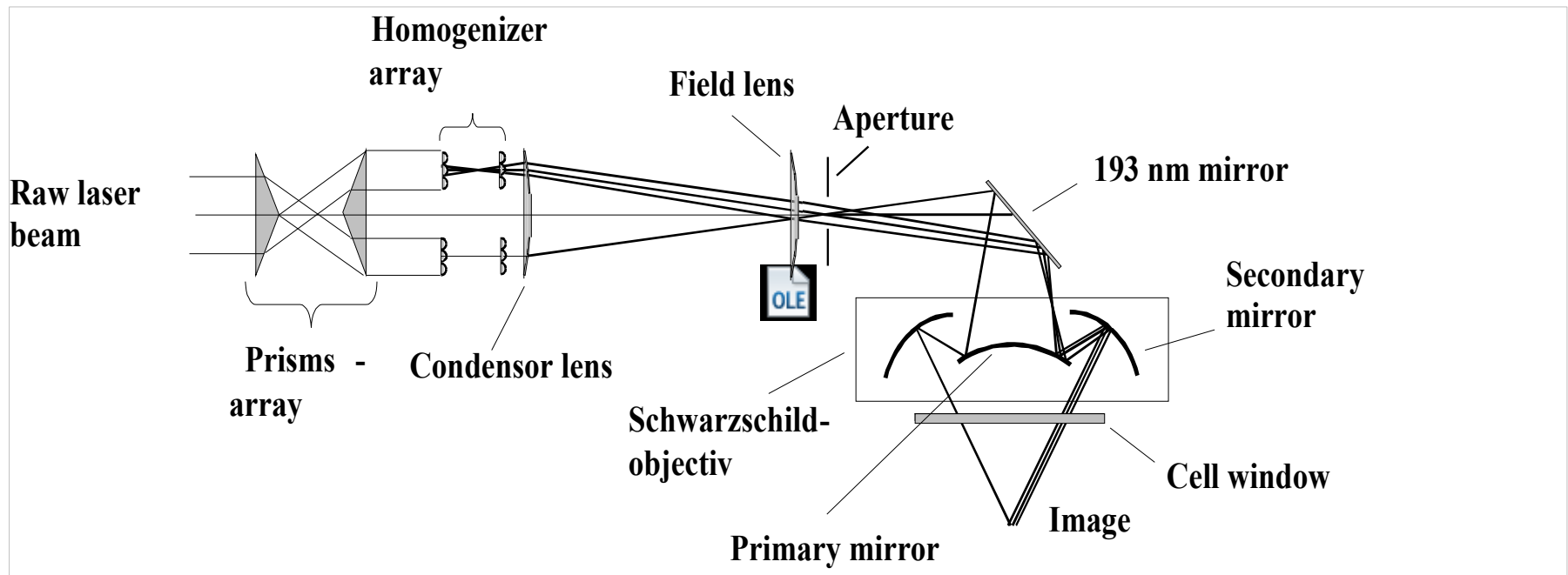
$$\omega_1 / \omega_2 = f(u)$$



# OPO Spectra Physics



# Laser beam homogenization

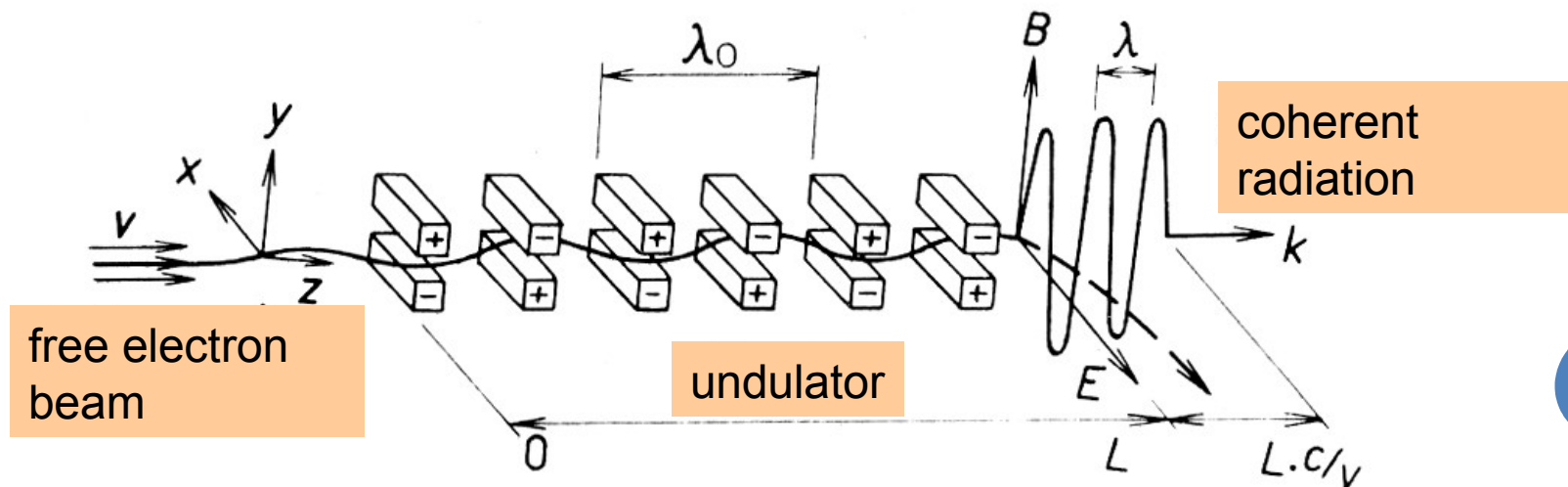


# Free electron laser FEL

The active environment is relativistic electrons passing through the periodic magnetic field. Electrons emit electromagnetic radiation as they move along curved orbits with wavelength:

( $\gamma \ll \lambda_0$ ,  $\gamma$  is the so-called relativistic factor):

$$\lambda = \lambda_0 \left( \frac{c}{v} - 1 \right) = \frac{\lambda_0}{2\gamma^2} \quad \gamma = \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}}$$



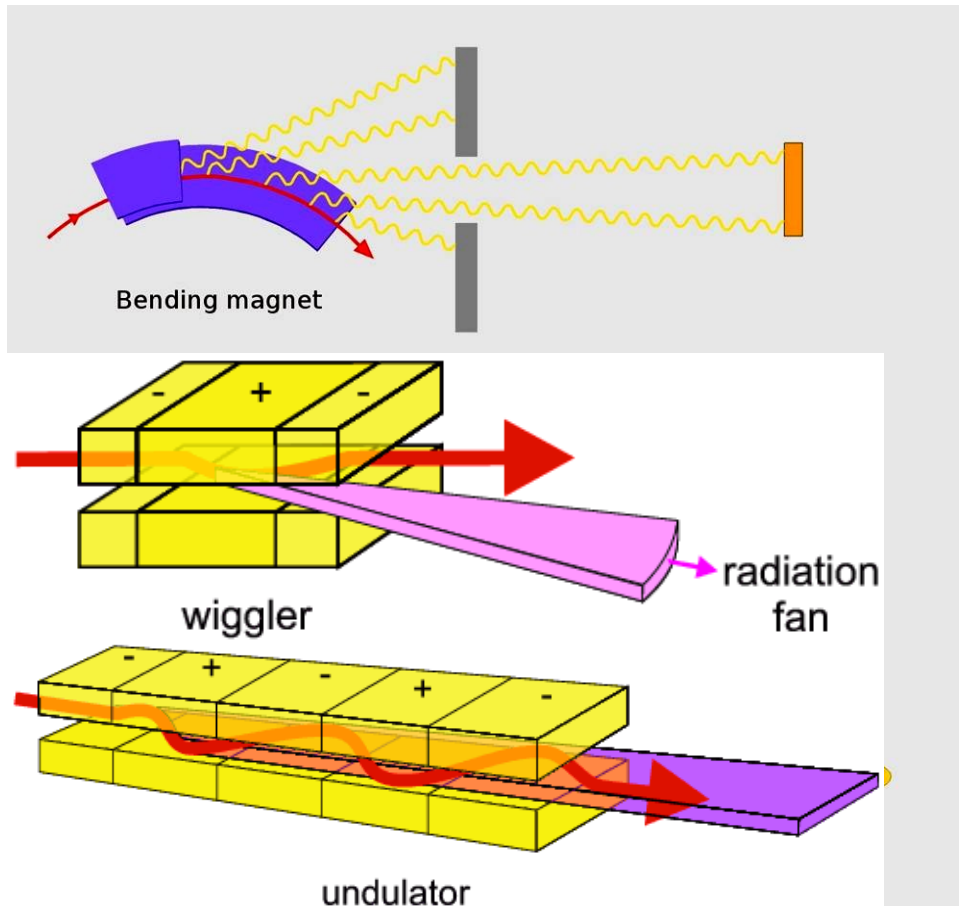


# Spectral brilliance

- For characterization of high-intensity sources (mainly synchrotron) radiation, the term **spectral brilliance** is introduced, indicates the number of radiated photons per second per 1 mm<sup>2</sup> of the radiation source area, at 1 mrad<sup>2</sup> divergence and at 10% of the wavelength range ( $\Delta\lambda / \lambda = 0.1$ ).
- The narrower and more parallel the beam is and the more the photons are concentrated in the narrowest wavelength range, the higher the spectral brilliance.
- This is inversely related to **emittance**, which is essentially the product of the dimension of the radiation source and the radiation divergence.



# Synchrotron radiation



## Sources of magnetic field

- bending magnets
- undulator
- wigglers

## Free electron lasers

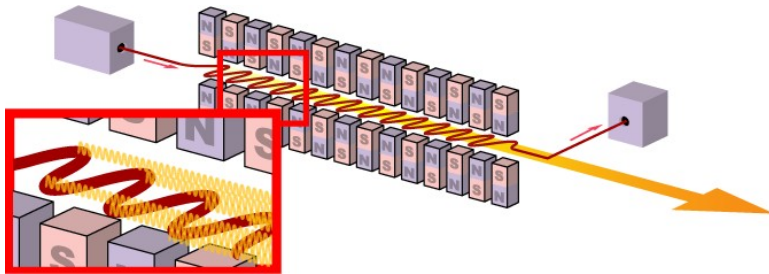


# Free Electron Laser

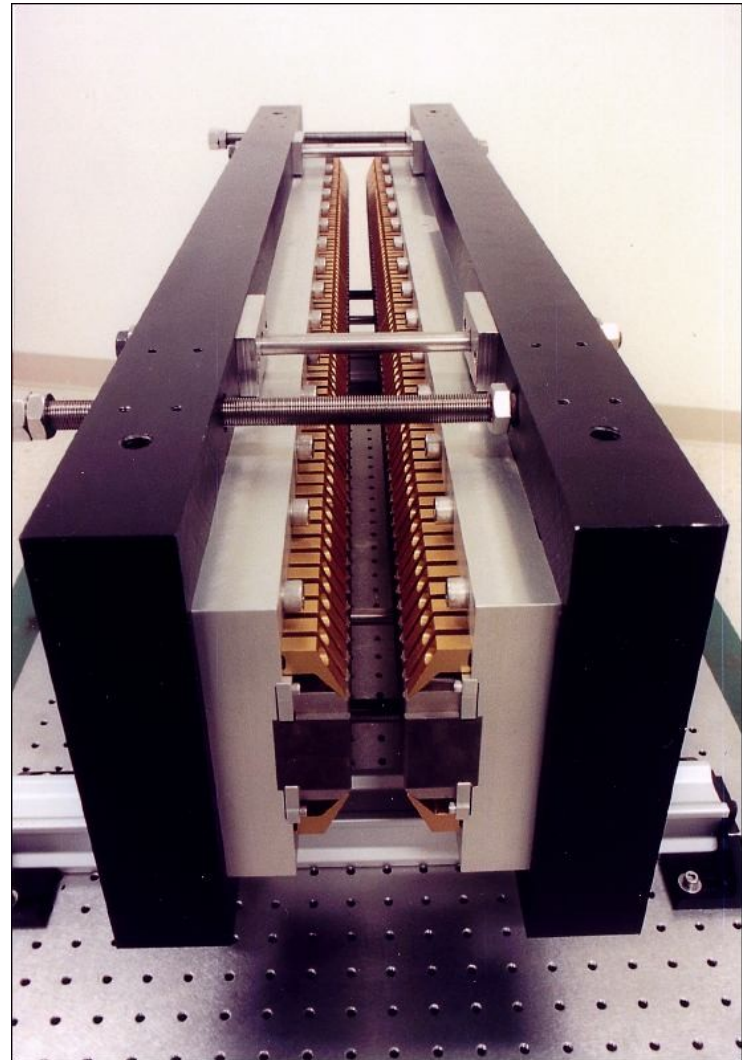
- The **fourth generation** of SZ sources is based on the use of linear accelerators, which allow to reduce emittance and shorten the pulse length. If a short electron burst runs through a sufficiently long undulator, then the electromagnetic wave generated at each location of the undulator proceeds together with the electron beam and interacts with it and results **free electron laser – FEL**. It is characterized by high brilliance, significantly higher than that of a classic undulator, coherence and shortness of pulses reaching tens of fs. A very long linear accelerator is required to accelerate electrons to GeV values.



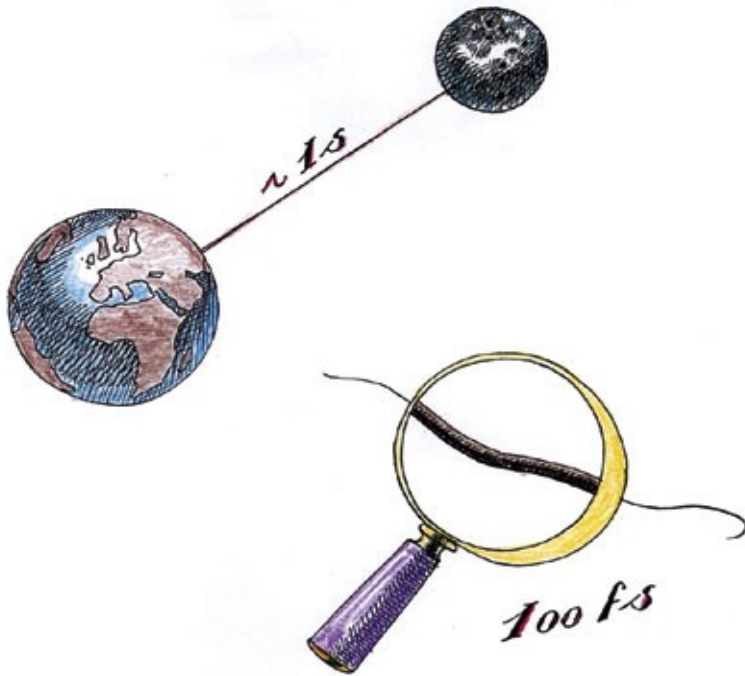
# Free Electron Laser



Electron clusters move along a wavy orbit. Charged particles that change their speed (direction enough) emit. For understanding, let's imagine that electrons in a cluster move along a sinusoid along an undulator. If we look at them from the end of this axis, we do not see that they are moving towards us, but we see oscillate cloud of charged particles. The clusters thus generate a coherent X-ray beam. Behind the undulator, the electrons are diverted by a strong magnetic field and the resulting X-ray laser beam continues into the experiment hall

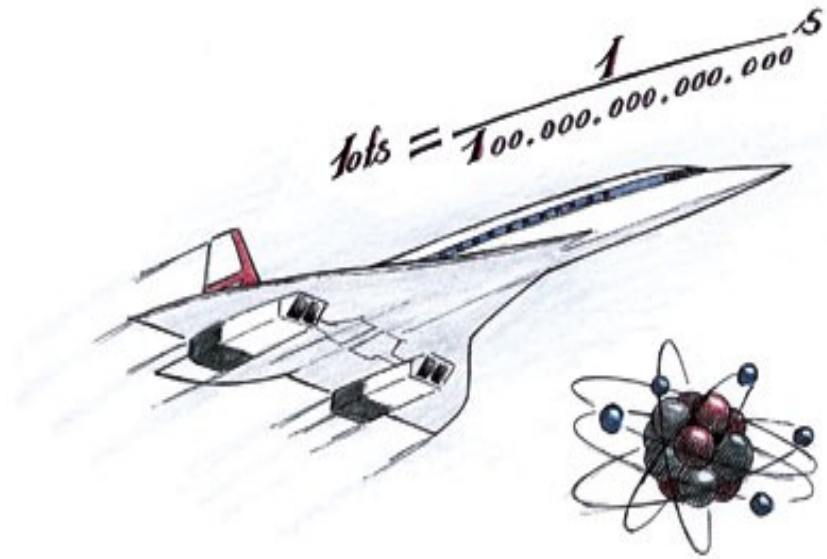


# Femtoseconds



Average Earth – Moon distance is around 380,000 km. Light moving at 300,000 km per second will travel this distance in a little more than 1 second...

Within 100 fs however light passes only 30  $\mu\text{m}$ , i.e. less than the thickness of the hair.



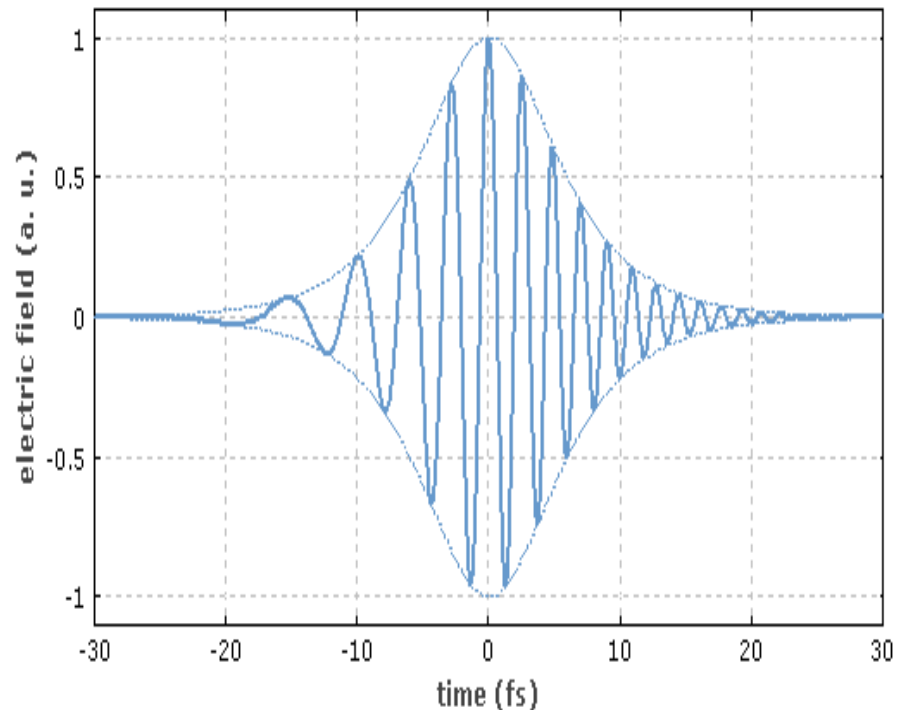
Concorde supersonic flies at approximately 2 Mach (twice the speed of sound in air), or 600 m / sec (2160 km / h). Over a period of 10 fs fly only 6 picometres ( $6 \cdot 10^{-12}$  m), it is 10 times less than the diameter of the carbon atom.



# Chirp

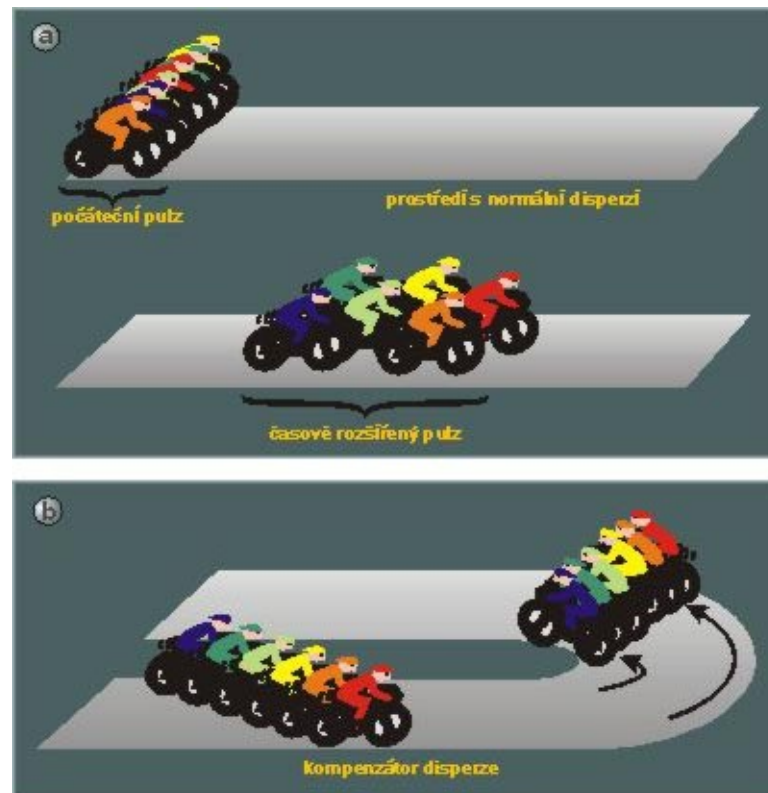
- In the case of pulses of optical radiation, a **Chirp** means a gradual change of frequency during the pulse (increase or decrease). This means that the frequency on the leading edge is different from the closing edge. If such a pulse propagates in a dispersion medium, the rate of propagation of the radiation in the leading edge is less (or greater) than in the closing portion, thereby shortening (or prolonging) the pulse.

Uncertainty relations make it impossible to have a short ( $\sim$  fs) light pulse in the visible region of the spectrum that is monochromatic



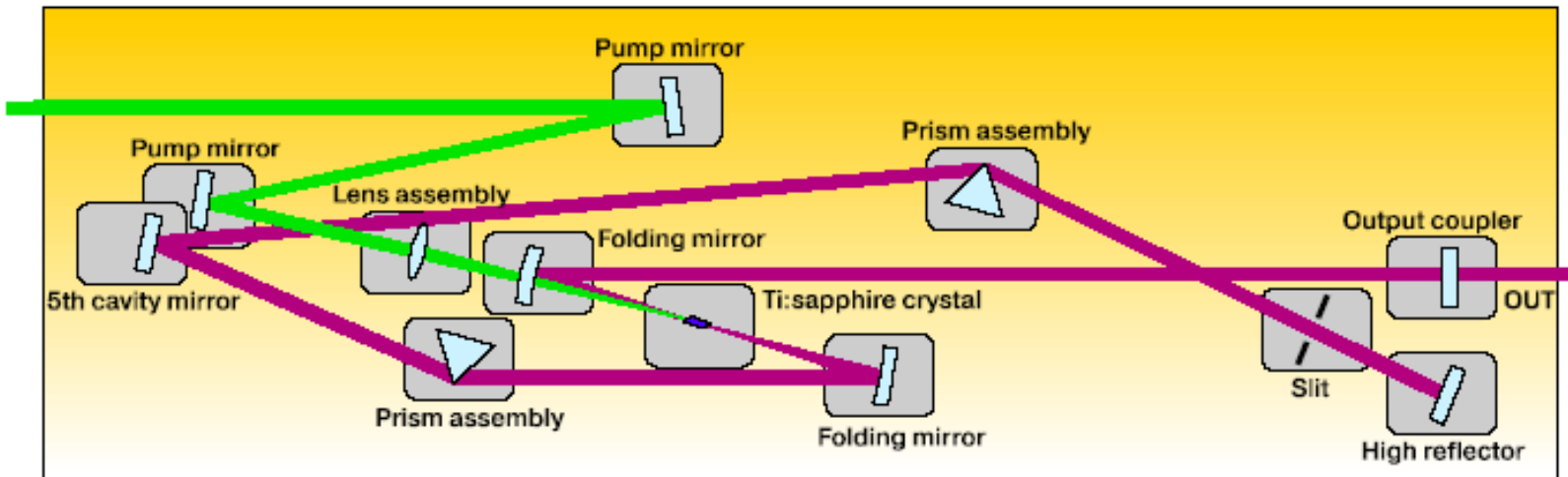
# Chirp of femtosecond pulse

- In a dispersed environment, the femtosecond pulse, which has a large frequency range (large wavelength range), is extended by the different dispersion (refractive index, light velocity) of different wavelengths. It is therefore necessary to make a correction with the dispersion compensator to maintain the pulse time profile.



# Model TISSA-20: $< 20$ fs

- Stable Kerr-lens mode-locking operation 5-mirror compact cavity design.
- Model TISSA-20: Seeding source of broadband femtosecond pulses for Ti:sapphire amplifiers

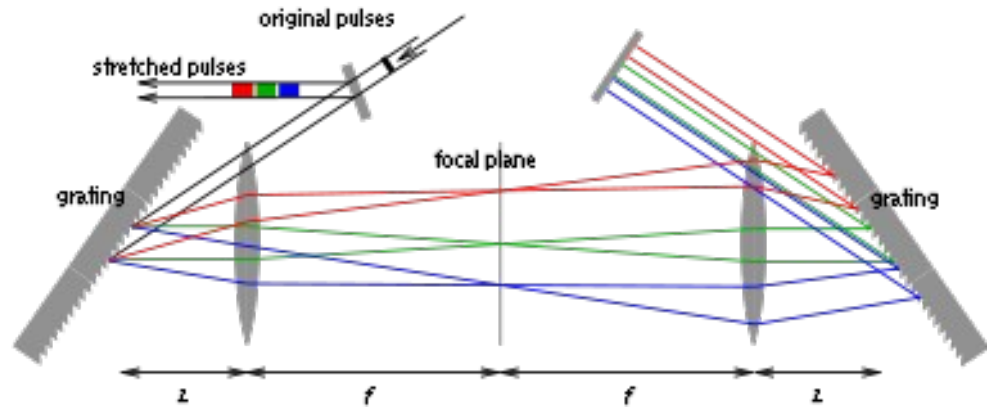
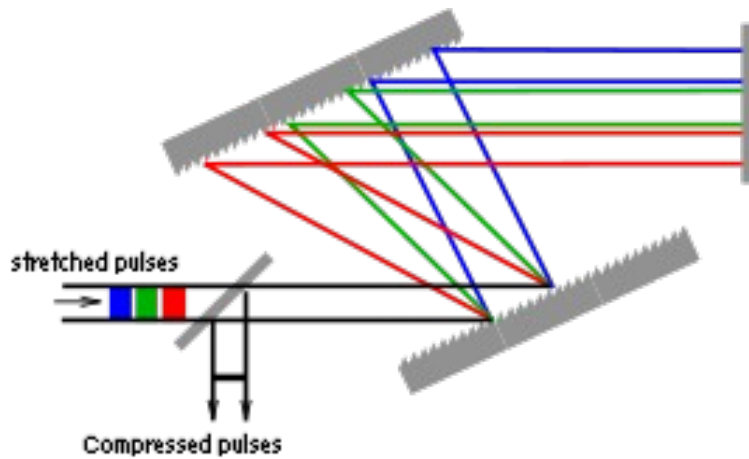




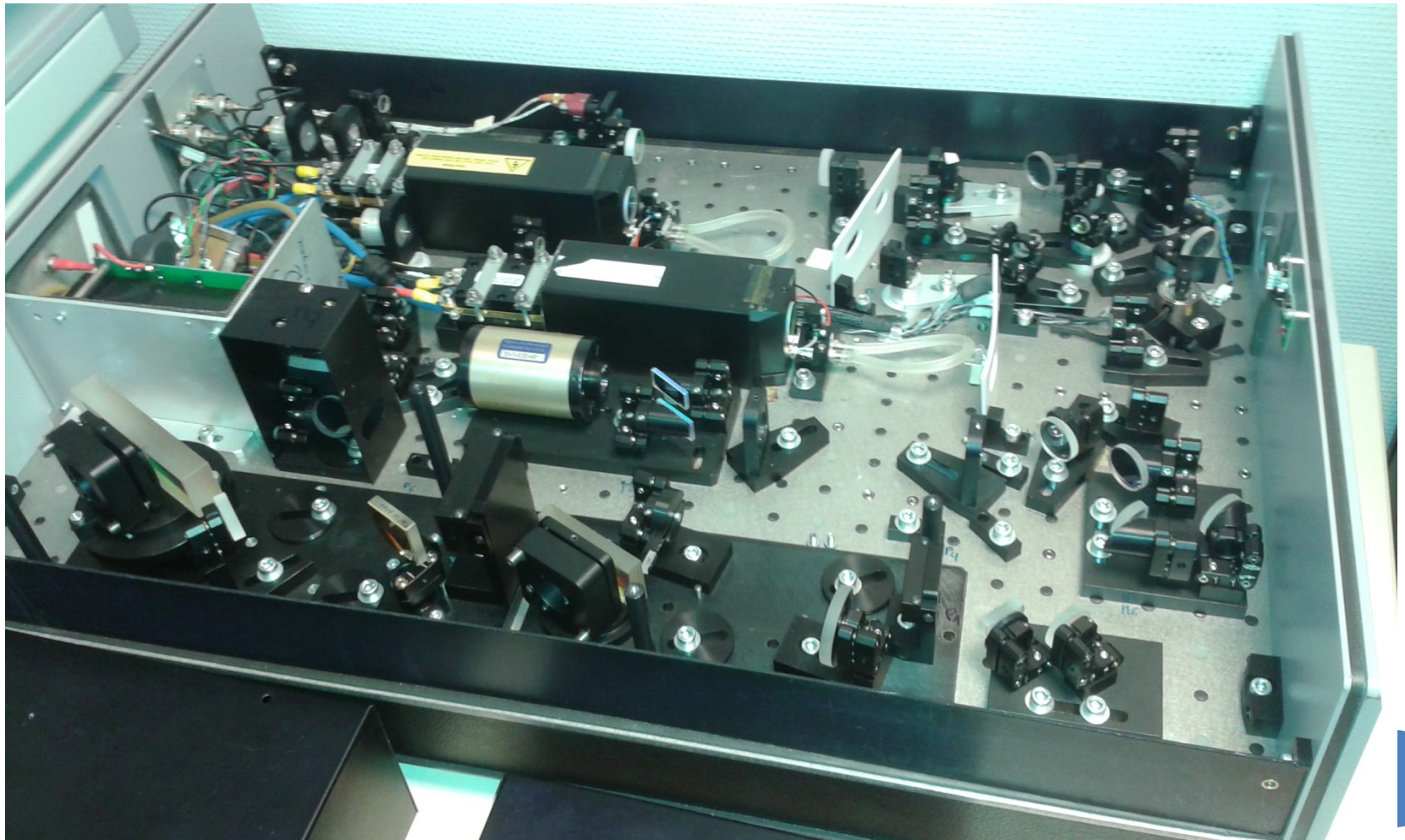
# Stretcher and compressor design

With gratings compressor

With gratings stretcher



# Stretcher and compressor design



# Amplification of femtosecond pulses with power PW

