

4 Electrons in a periodic potential (41 pts)

Due date: **Wednesday, 27th November 2024.**

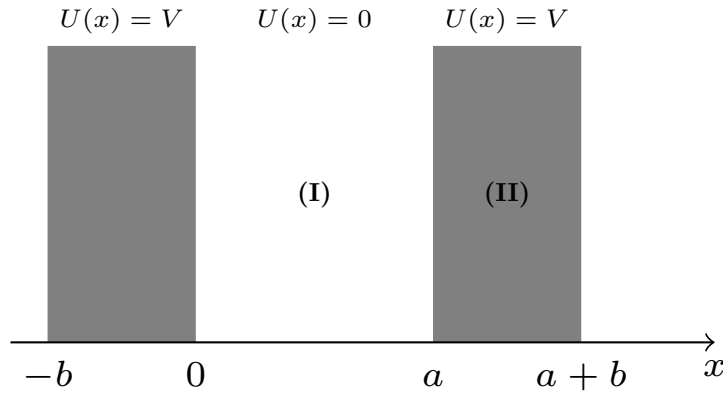
4.1 Concept of bands: 1D approximation of a crystal

Consider a one-dimensional crystal with the potential given by:

$$\forall x \in \dots \cup [0, a[\cup [a + b, (a + b) + a[\cup \dots, U(x) = 0$$

$$\forall x \in \dots \cup [-b, 0[\cup [a, a + b[\cup [(a + b) + a, 2(a + b)[\cup \dots, U(x) = V > 0$$

periodically repeated over $N \gg 1$ periods of length $d \equiv a + b$. We use periodic boundary conditions. The x axis is therefore divided in zones, where the potential is either 0 or V . The region $[0, a]$ ($U(x) = 0$) is labeled region I, the region $[a, a + b]$ is labeled region II ($U(x) = V$).



- (i) (1 pt) Give the expression for a wave function for a particle in this potential, using Bloch's theorem. The corresponding wave vector will be denoted k .
- (ii) (1 pt) Give the general expressions for the solutions of the Schrödinger equation $\psi_I(x)$ and $\psi_{II}(x)$ for a particle of mass m , in region I and in region II respectively. These involve plane waves with wave vector k_1 and k_2 respectively, as well as two unknowns each, which will be denoted A, B for region I, C, D for region II. For simplicity, assume $E > V$, where E is the energy of the particle. Please pay attention to the specific domain of definition for each of the functions $\psi_I(x)$ and $\psi_{II}(x)$!
- (iii) (4 pts) Use continuity constraints at points $x = 0$ and $x = a$ on the wave function, in order to obtain a set of four independent linear equations involving the unknowns introduced in the preceding question.
- (iv) (6 pts) Show that the dispersion relation is given by:

$$\cos(k_1 a) \cos(k_2 b) - \frac{k_1^2 + k_2^2}{2k_1 k_2} \sin(k_1 a) \sin(k_2 b) = \cos(kd) \quad (13)$$

$$k_1^2 - k_2^2 = \frac{2mV}{\hbar^2}$$

This is not a difficult calculation, but it is long; you need to keep things organized.

- (v) (2 pts) Deduce from the above the expression for the dispersion relation in the configuration where $E < V$. (Use the notation $\frac{\hbar^2 \chi^2}{2m} = V - E$, you do **not** need to repeat the derivation!).
- (vi) (5 pts) Use $a = \frac{10\hbar}{\sqrt{2m}}$, $b = \frac{2\hbar}{\sqrt{2m}}$, and $V = 1$. Plot the the profile of the left hand side of the dispersion relation, as a function of E , for $0 < E < 7$. Provide me with a transcript of your code, explaining how it works. Comment on your results, in relation to the concept of bands in solids.

4.2 Band structure in a periodic potential

Consider a two dimensional square lattice, with periodic boundary conditions, and with the periodic potential given by

$$U(x, y) = -4V \cos\left(\frac{2\pi x}{a}\right) \cos\left(\frac{2\pi y}{a}\right),$$

where a is the lattice parameter and $V > 0$. In this problem, we explore the band structure near the point $(\frac{\pi}{a}, \frac{\pi}{a})$, i.e., the corner of the 1st BZ. In all questions, we assume that $V \ll (\hbar^2/2m)(\pi/a)^2$. The arguments used when solving the problem can also be used when discussing band structures of simple metals.

- (i) (1 pt) Give a description of the reciprocal lattice.
- (ii) (2 pts) Give the expression of the Fourier expansion of $U(\mathbf{r})$.
- (iii) (1 pt) Explain in plain words what the statement $V \ll (\hbar^2/2m)(\pi/a)^2$ tells us about the system. In particular, explain why this allows us to look for solutions of the Schrödinger equation in the form

$$\psi(\mathbf{r}) = \sum_{\mathbf{k} \in \text{BVK conditions}} c_{\mathbf{k}} e^{i\mathbf{k} \cdot \mathbf{r}}. \quad (14)$$

- (iv) (3 pts) Derive the so-called *central equation*, that is, the Schrödinger equation in \mathbf{k} -space. This provides you with a relation between all the coefficients $c_{\mathbf{k}}$: for any vector \mathbf{K} from the reciprocal lattice,

$$\left[\frac{\hbar^2}{2m} (\mathbf{k} - \mathbf{K})^2 - E \right] c_{\mathbf{k} - \mathbf{K}} + \sum_{\mathbf{K}'} U_{\mathbf{K}' - \mathbf{K}} c_{\mathbf{k} - \mathbf{K}'} = 0 \quad (15)$$

Hint: Substitute the solution in the form given in (iii) to Schrödinger equation.

- (v) (2 pts) Show that this allows you to write the solutions of the Schrödinger equation in the form

$$\psi_{\mathbf{k}}(\mathbf{r}) = \sum_{\mathbf{K}} c_{\mathbf{k} - \mathbf{K}} e^{i(\mathbf{k} - \mathbf{K}) \cdot \mathbf{r}}, \quad (16)$$

where the sum runs over all reciprocal lattice vectors, and \mathbf{k} is any vector in the first Brillouin zone.

For \mathbf{k} close to $(\frac{\pi}{a}, \frac{\pi}{a})$, only the four following coefficients in the expansion of $\psi_{\mathbf{k}}(\mathbf{r})$, Eq. (16) are non-negligible:

$$c_{\mathbf{k} - \mathbf{K}}, \quad \forall \mathbf{K} \in \left\{ \begin{pmatrix} 1 \\ 1 \end{pmatrix}, \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \begin{pmatrix} 0 \\ 1 \end{pmatrix}, \begin{pmatrix} 0 \\ 0 \end{pmatrix} \right\},$$

in units of $\frac{2\pi}{a}$.

- (vi) (4 pts) Solve the central equation for \mathbf{k} close to $(\frac{\pi}{a}, \frac{\pi}{a})$: calculate the perturbed energy levels $E(\mathbf{k})$.

4.3 Density of states from the tight-binding model

Consider a tight-binding model on a square lattice, whose band structure is given by (we assume the lattice parameter is $a = 1$ for simplicity)

$$\epsilon(\vec{\mathbf{k}}) = -2t \sum_i \cos(k_i), \quad (17)$$

where i runs over all dimensions. We consider lattices in dimensions $d = 1, 2, 3$, so that the density of states can be written in a general manner

$$D(\epsilon) = \frac{2}{L^d} \sum_{\vec{\mathbf{k}}} \delta(\epsilon - \epsilon_{\vec{\mathbf{k}}}) \quad (18)$$

- (i) (3 pts) Compute analytically the expression for $D(\epsilon)$ for $d = 1$.
- (ii) (3 pts) For $d = 2, d = 3$, use numerics. Plot the profiles of $D(\epsilon)$ for $d = 1$ (analytic and numerical results together), $d = 2, d = 3$. Provide me with a transcript of your code, which should implement the elementary summation in some form (do not use the “histogram” feature of some numerical libraries).
- (iii) (3 pts) Discuss and comment your results.