

Laser-induced fluorescence

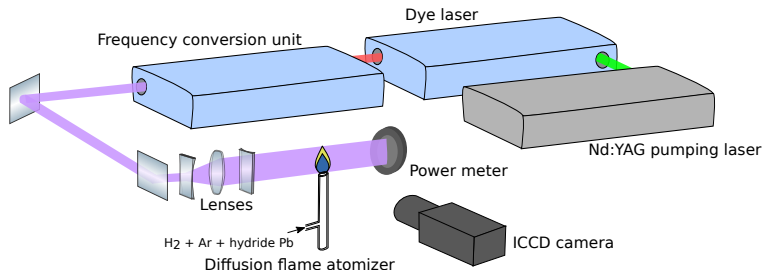
Martina Mrkvičková

3. 12. 2020

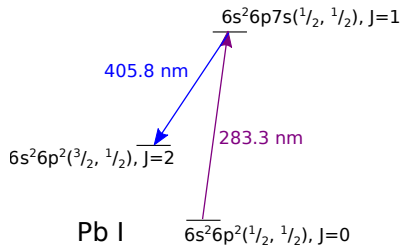
Table of contents

- 1 Introduction
- 2 Experimental

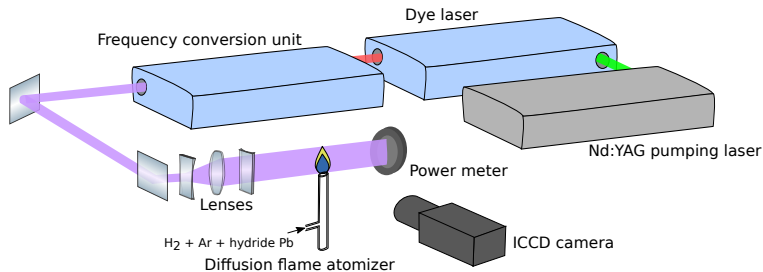
Laser-induced fluorescence



- method for determination of concentrations of various particles in plasma (atoms, molecules, radicals) also in the ground state
- high temporal and spatial resolution
- high sensitivity



Laser-induced fluorescence

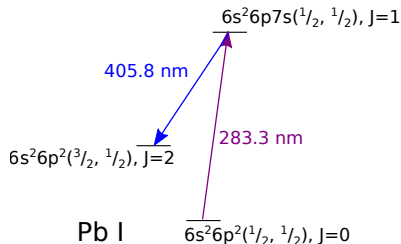


fluorescence signal:

$$f(\vec{r}) = \int_0^{\infty} A_{32} n_3(\vec{r}, t) dt$$

concentration of measured species:

$$n = n_1(0) = n_1(t) + n_2(t) + n_3(t)$$



Examples

(elsewhere)

Laser-induced fluorescence

Connection between the fluorescence signal and the concentration of excited state is easy:

$$f(\vec{r}) = \int_0^{\infty} A_{32} n_3(\vec{r}, t) dt$$

But what is the connection between n_3 and Σn ?

Laser-particle interactions (not only)

- absorption of laser photon
- two-photon absorption
- spontaneous emission
- stimulated emission
- collisional quenching
- amplified spontaneous emission
- atomisation by laser
- ionisation by laser
- ...

Spontaneous emission

$$\left(\frac{dn_3}{dt}\right)_{\text{sp.emission}} = -(A_{31} + A_{32}) n_3$$

solution:

$$n_3 = n_3(0) e^{-\frac{t}{\tau_0}}$$
$$\tau_0 = \frac{1}{\sum_i A_{3i}}$$

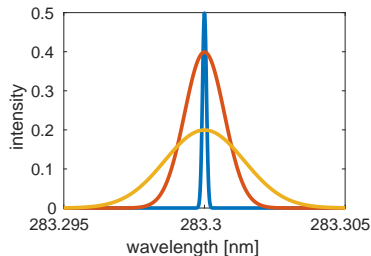
- A – Einstein coefficient of emission (you can find on NIST)
- τ_0 – radiative lifetime

Absorption of laser photon

$$\left(\frac{dn_3}{dt}\right)_{\text{absorption}} = \kappa \frac{B_{13}}{c} I n_1$$

- B_{13} - Einstein coefficient of absorption ($B_{13} \sim A_{31}$)
- I – laser intensity
- κ – spectral overlap of laser $I(\nu)$ and absorption line $g(\nu)$

$$\kappa = \int_{\nu} g(\nu) I(\nu) d\nu$$



Stimulated emission

$$\left(\frac{dn_3}{dt}\right)_{\text{stim.emission}} = -\kappa \frac{B_{31}}{c} I n_3$$

Collisional quenching

$$\left(\frac{dn_3}{dt}\right)_{\text{quenching}} = -(Q_{31} + Q_{32}) n_3$$
$$Q = \sum_m q_m n_m$$

- Q – effective quenching rate
- n_m – concentration of collisional partner m
- q_m – quenching coefficient for collisions with m

Two-photon absorption

$$\left(\frac{dn_3}{dt}\right)_{\text{TPabsorption}} = \kappa^{(2)} G^{(2)} \sigma^{(2)} \left(\frac{I}{h\nu}\right)^2 n_1$$

- used in method Two-photon absorption laser-induced fluorescence (TALIF)
- $\sigma^{(2)}$ – two-photon absorption cross section
- $G^{(2)}$ – two-photon statistical factor

Atomisation, ionisation, ... by laser

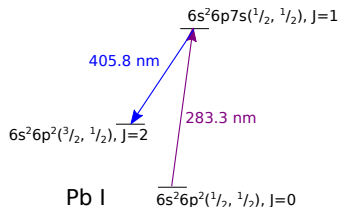
- processes leading to gain or loose of particles in the fluorescence process due to laser
- example for gain: laser causes dissociation of O_3 to $O_2 + O$
- example for loose: laser causes ionization of O to O^+

Rate equations – simple three-level model

$$\begin{aligned}\frac{dn_1}{dt} &= -\kappa \frac{B_{13}}{c} I n_1 + \kappa \frac{B_{31}}{c} I n_3 + \\ &\quad (A_{31} + Q_{31}) n_3 + (A_{21} + Q_{21}) n_2 \\ \frac{dn_2}{dt} &= (A_{32} + Q_{32}) n_3 - (A_{21} + Q_{21}) n_2 \\ \frac{dn_3}{dt} &= \kappa \frac{B_{13}}{c} I n_1 - \kappa \frac{B_{31}}{c} I n_3 - \\ &\quad (A_{31} + A_{32} + Q_{31} + Q_{32}) n_3\end{aligned}$$

Measured fluorescence signal:

$$f(\vec{r}) = \int_0^{\infty} A_{32} n_3(\vec{r}, t) dt$$



We want to find relation between f and Σn .

Solution: low-intensity approximation – "Linear regime"

assumptions:

- $n_1 \gg n_2, n_3$
- $n_1 \approx \text{konst.}$
- no stimulated emission

solution:

$$n_3(t) = \left(\frac{\kappa}{c} B_{13} I n_1 - \frac{dn_3}{dt} \right) \tau$$
$$f = A_{32} \int_0^{\infty} \left(\frac{\kappa}{c} B_{13} I n_1 - \frac{dn_3}{dt} \right) \tau dt$$
$$f(\vec{r}) = A_{32} \tau \kappa \frac{B_{13}}{c} n_1 L(\vec{r})$$

- $L(\vec{r})$ – temporal integral of laser intensity, $L(\vec{r}) = \int_0^{\infty} I(\vec{r}, t) dt$
- τ – lifetime of the excited state, $\tau = \frac{1}{A_{31} + A_{32} + Q_{31} + Q_{32}}$

Note for the lifetime

What is happening to the excited state after the laser pulse?

$$\frac{dn_3}{dt} = -(A_{31} + A_{32} + Q_{31} + Q_{32}) n_3$$

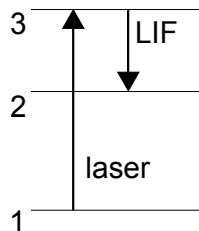
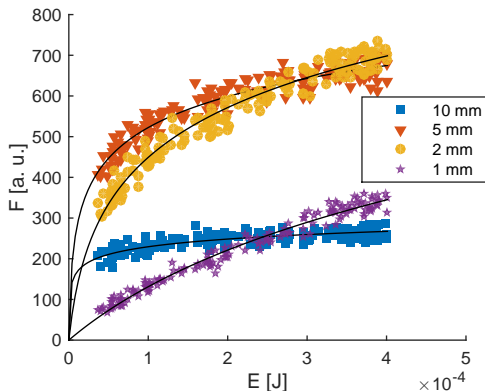
solution:

$$n_3 = n_3(0) e^{-\frac{t}{\tau}}$$
$$\tau = \frac{1}{A_{31} + A_{32} + Q_{31} + Q_{32}}$$

- τ – "real" lifetime of the excited state

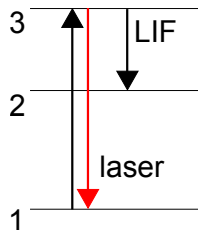
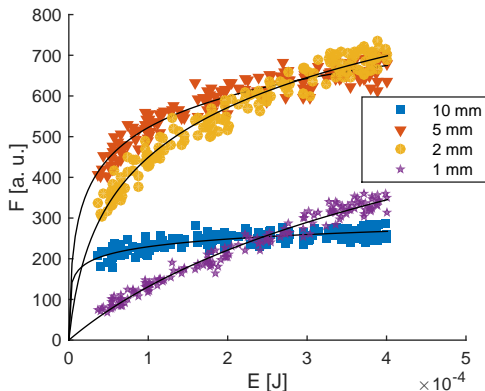
A bit more laser intensity?

- the dependence of signal on energy stops to be linear
- reason: stimulated emission and depletion of the ground state



A bit more laser intensity?

- the dependence of signal on energy stops to be linear
- reason: stimulated emission and depletion of the ground state



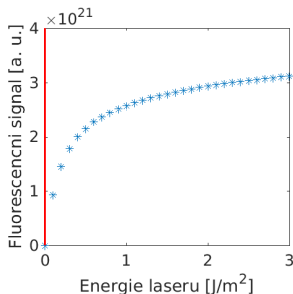
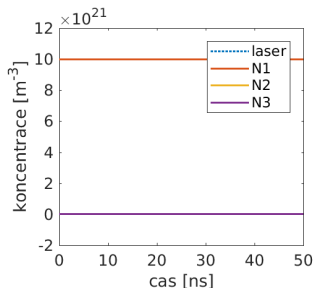
Saturation of LIF

$$\frac{dn_1(t)}{dt} = -\frac{I_{laser}(t)}{c} B_{13} n_1(t) + \frac{I_{laser}(t)}{c} B_{31} n_3(t) + (A_{31} + Q_3) n_3(t) + (A_{21} + Q_2) n_2(t)$$

$$\frac{dn_2(t)}{dt} = A_{32} n_3(t) - (A_{21} + Q_2) n_2(t)$$

$$\frac{dn_3(t)}{dt} = \frac{I_{laser}(t)}{c} B_{13} n_1(t) - \frac{I_{laser}(t)}{c} B_{31} n_3(t) - (A_{31} + A_{32} + Q_3) n_3(t)$$

$$F = \int_0^\infty A_{32} n_3(t) dt$$



Saturation of LIF

$$\frac{dn_1(t)}{dt} = -\frac{I_{laser}(t)}{c} B_{13} n_1(t) + \frac{I_{laser}(t)}{c} B_{31} n_3(t) + (A_{31} + Q_3) n_3(t) + (A_{21} + Q_2) n_2(t)$$

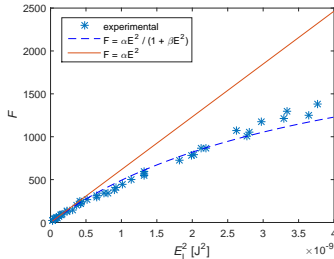
$$\frac{dn_2(t)}{dt} = A_{32} n_3(t) - (A_{21} + Q_2) n_2(t)$$

$$\frac{dn_3(t)}{dt} = \frac{I_{laser}(t)}{c} B_{13} n_1(t) - \frac{I_{laser}(t)}{c} B_{31} n_3(t) - (A_{31} + A_{32} + Q_3) n_3(t)$$

$$F = \int_0^{\infty} A_{32} n_3(t) dt$$

Saturation: conclusion

Thanks to some modelling, we are able to find "correction factors" which quantificate how many excited species were lost in comparison with ideal linear regime.



How much fluorescence is really detected?

- signal M is integrated from some definite volume V , not only point in space \vec{r}
- fluorescence is emitted to all angles, only small part $\frac{\Omega}{4\pi}$ goes to the camera
- wavelength filters have some transmittivity T , camera has some sensitivity C , ...

$$M = T \cdot C \cdot \frac{\Omega}{4\pi} \cdot \int f(\vec{r}) \, dV$$

How to obtain $\frac{\Omega}{4\pi}$?

- Let's measure some another signal caused by the exactly same shape of the laser beam, but we know exactly how much signal we should get from it.
 - 1 Rayleigh scattering of the laser
 - 2 fluorescence of noble gases with known concentration
- Than compare the theoretical signal and really obtained signal.

Formula for absolute concentration

$$n = \frac{M}{E_L} \zeta \frac{1}{T C} \frac{c}{\tau A_{32} B_{13} \kappa} \frac{M_{C,theoretical}}{M_C}$$

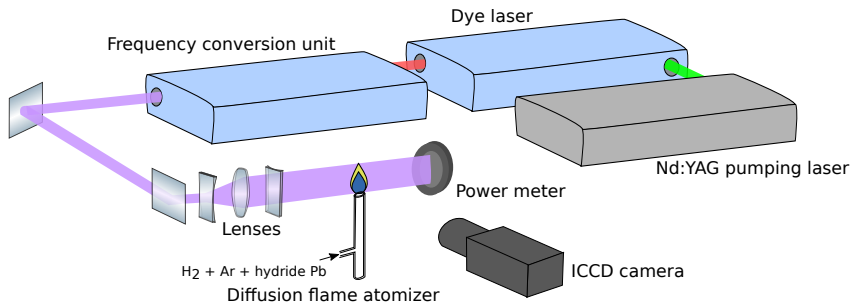
- n – concentration of measured species
- M – measured fluorescence signal
- E_L – laser pulse energy , $E_L = \int \int_{S,t} I(\vec{r}, t) dS dt$
- ζ – correction factor describing the saturation of signal
- T – wavelength filter transmissivity
- C – camera sensitivity (depends on wavelength)
- c – speed of light
- τ – lifetime of the excited state
- A_{32} – Einstein coefficient of the fluorescence transition
- B_{13} – Einstein coefficient of absorption of laser photon
- κ – spectral overlap of laser and absorption line
- $M_{C,theoretical}$ – theoretical amount of calibration signal
- M_C – really detected calibration signal

Obsah přednášky

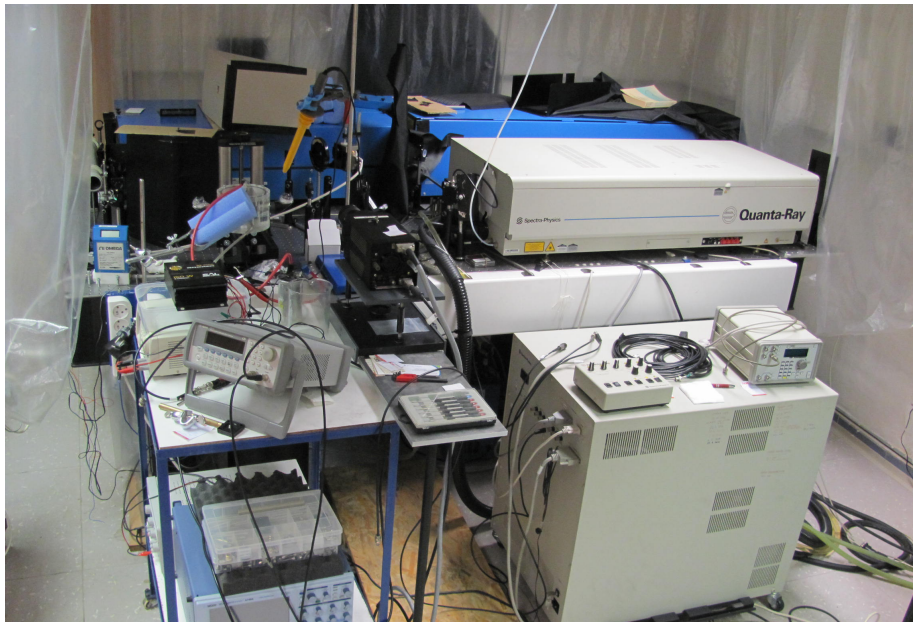
1 Introduction

2 Experimental

LIF equipment

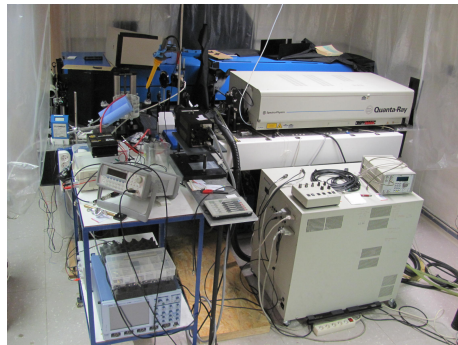


Our LIF laboratory



Our laser

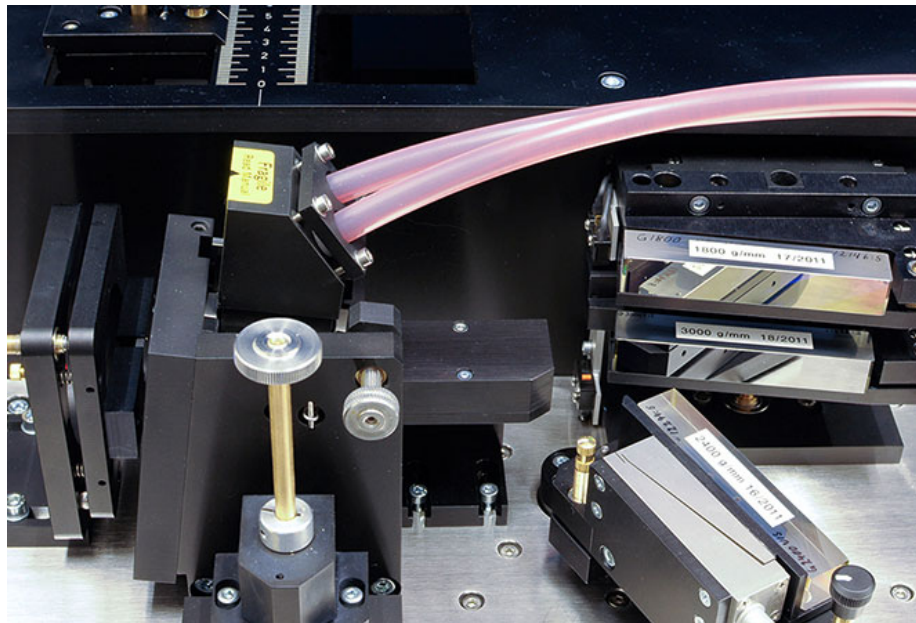
- Nd:YAG pumping laser + dye laser + frequency conversion unit
- pulsed – 10 ns, 30 Hz
- tunable wavelength (1064 nm fundamental; higher harmonics; cca 200 - 900 nm from dye laser and FCU)
- peak power: kW (in UV) - GW (in IR)
 - ▶ comparison: laser pointer 20 mW, ELI 10 PW



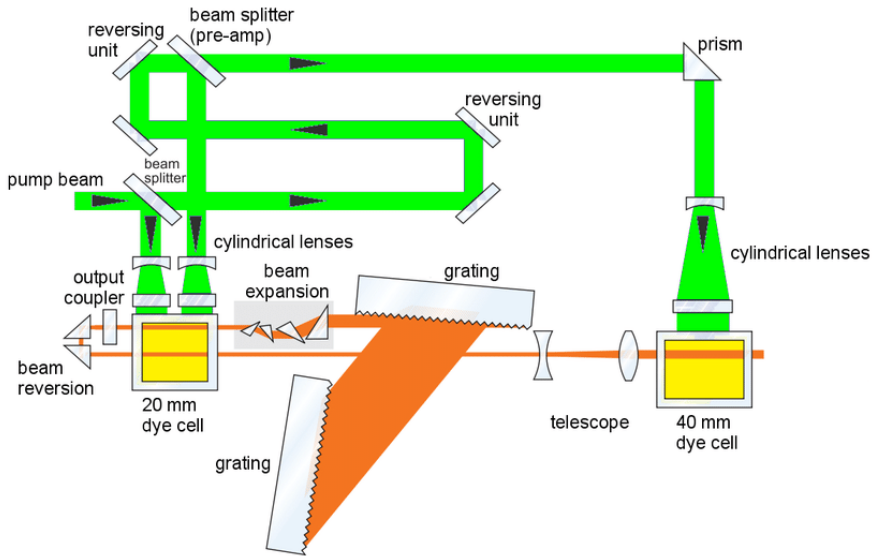
Inside the pump laser



Dye laser



Dye laser



Detector: ICCD camera

- intensified CCD kamera (with microchannel plate)
- gate down to 2 ns
- high sensitivity, both in UV and near IR
- external trigger
- optics: quartz lens, wavelength filter

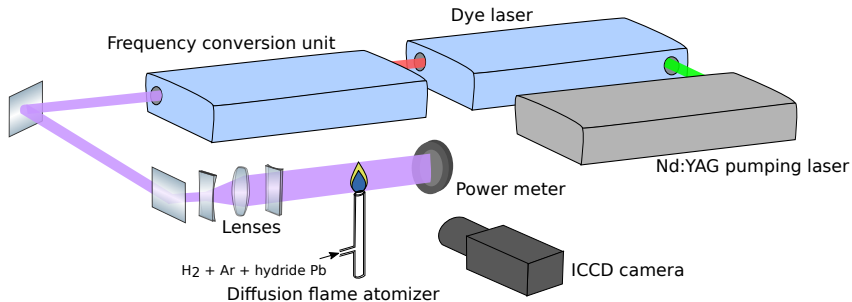


Measurement

- prepare optical path (...)
- set the correct wavelength and find the first fluorescence signal
- check parasitic effects, signal saturation
 - ▶ does the signal disappear when the laser wavelength is detuned?
 - ▶ does the signal have linear dependence on the laser energy?
- measure the shape of the absorption line ($\rightarrow \kappa$)
- measure the timeshape of the fluorescence ($\rightarrow \tau$)
- nice pictures of the whole fluorescence signal ($\rightarrow M$)
- calibration ($\rightarrow M_c$)
- do not forget to get darkframes and monitor the energy E_L for all the measurements

$$n = \frac{M}{E_L} \zeta \frac{1}{T C} \frac{c}{\tau A_{32} B_{13} \kappa} \frac{M_{c,theoretical}}{M_c}$$

Practise: measurement of lead atoms in a diffuse flame atomizer



Practise: measurement of lead atoms in a diffuse flame atomizer

tasks to do in Octave/Matlab:

- 1 load the measurement 'citlivost1.SPE' and see how the fluorescence signal look like
- 2 evaluate the fluorescence lifetime from the file 'dobazivota1.SPE'
- 3 estimate the absolute concentrations of atomic lead in the measurement 'citlivost1.SPE'